

952. Nonlinear vibrations of beams with spring and damping delayed feedback control

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Abstract. The primary, subharmonic, and superharmonic resonances of an Euler–Bernoulli beam subjected to harmonic excitations are studied with damping and spring delayed-feedback controllers. By method of multiple scales, the non-linear governing partial differential equation is transformed into linear differential equations directly. Effects of the feedback gains and time-delays on the steady state responses are investigated. The velocity and displacement delayed-feedback controllers are employed to suppress the primary and superharmonic resonances of the forced nonlinear oscillator. The stable vibration regions of the feedback gains and time-delays are worked out based on stability conditions of the resonances. It is found that proper selection of feedback gains and time-delays can enhance the control performance of beam's nonlinear vibration. Position of the bifurcation point can be changed or the bifurcation can be eliminated.

Keywords: time-delays, nonlinear vibration, nonlinear control, resonance, suppression.

1. Introduction

In the past decades, a considerable amount of research has helped us better understand the effect of time delays on the behavior of nonlinear dynamical systems, which are controlled by the linear or nonlinear feedback controllers to mitigate vibrations. The intrinsic time delays, which inevitably exist in processing of the actuation mechanisms, filters, circuits, and controllers, can lead to an unacceptable detrimental delay period between the controller input and real-time system actuation in the control engineering [1-2]. However, the time-delays feedback controllers can be utilized to change the dynamic behavior of nonlinear systems [2]. Recently, the application of time delays in the suspension controlling has become a hot research topic.

The studies on delayed dynamic systems are basically parallel to those on the dynamic systems without any delay. The major approaches include those on stability of systems, such as eigenvalue and Lyapunov method [3], and those on the Hopf-bifurcation, the center manifold reduction [4], the Lyapunov-Schmidt [5], perturbation [6] and frequency domain [7] method. The stability criteria of linear systems, the Hopf bifurcation and chaotic behavior of non-linear systems have attracted many researchers' attention.

The dynamic behavior of delayed systems of discrete systems has received considerable concerns. Sato et al. [8] discussed the free and forced vibration of non-linear system with time-delays, using the center current theorem and average method to study its stable periodic solution and stability. Hu et al. presented analytical and numerical studies of primary resonance and subharmonic resonance of a harmonically forced Duffing oscillator under state feedback control with a time delay [9]. Ji and Leung studied the primary, superharmonic, and subharmonic resonances of a harmonically excited non-linear s.d.o.f system with two distinct time-delays in the linear state feedback [10]. Nima and Nader [11] found that the experimental results and theoretical findings were in good agreement, which demonstrated that the non-linear modeling framework could provide a better dynamic representation of the micro cantilever than the previous linear models.

Recently, stabilization of beam with delayed feedback has raised many researchers' interests. Mohammed et al. presented a comprehensive investigation of the effect of feedback delays on non-linear vibrations of a piezoelectric actuated cantilever beam [12]. Khaled et al. investigated the effect of time delays on stability, amplitude, and frequency-response behavior of a beam and found that even the minute amount of delays can completely alter the behavior and stability of the parametrically excited beam, leading to unexpected behavior and responses [13]. Mustapha and Mohamed [14] examined the control of self-excited vibration of a simply-supported beam subjected to axially high-frequency excitation. The primary resonance of a cantilever beam under state feedback control with a time delay was investigated [15]. Vibration control and high-amplitude response suppression could be performed with appropriate time-delays and feedback gains. Wang and Hu [16] used the Lambert W function to calculate the rightmost characteristic root of retarded time-delay systems and analyzed the stability of high order mode or the systems with multiple delays. Alhazza [17] presented a single-input and single-output multimode delayed-feedback control methodology to mitigate the free vibrations of a flexible cantilever beam. A comprehensive investigation of the effect of feedback delays on the non-linear vibrations of a piezo-electrically actuated cantilever beam is presented. Stability boundaries of mechanical controlled system with time delay was studied [18]. Qian and Tang [19] discussed the primary resonance and the subharmonic resonances of a non-linear beam under moving load by using time-delay feedback controller. Gohary et al. [20] studied the vibration suppression of a dynamical system to multi-parametric excitations via time-delay absorber.

Methods of treating weakly non-linear continuous systems can be broadly divided into three groups. In the first approach, harmonic balance method is applied to obtain a nonlinear boundary-value problem. In the second approach, the Galerkin procedure is always used to discretize the non-linear governing partial differential equation and boundary conditions. The two approximate methods are easy and straightforward for application. However, the calculation of higher-order approximate solution may be complex. In the third approach, Nayfeh et al. [21] calculated an approximate solution of the nonlinear differential equation by directly solving the continuous problem. In the study, the non-linear governing partial differential equation and boundary conditions were transformed into linear differential equations by using multiple scales directly. The obtained mode shapes were in full agreement with those obtained by using discretization because the latter was performed by using a complete set of basis functions that satisfied the boundary conditions. This method is beneficial to the elimination of coupling terms and removal of solving the multi-dimensional equations and it can easily be used to analyze the dynamic response of higher modes.

Liao et al. [22] presented a feasible methodology by which could achieve good control performance of a dynamic beam structure system with time delay effect. A spring time-delay controller was designed to achieve good control performance. In this paper, the delayed damping and spring feedback controllers are designed to change the beam's nonlinear dynamical behavior. The feedback gains and time-delays of damping controllers are easy to set and change by using active control facilities.

In this study, the spring and damping controllers are designed to control the dynamic behavior of an inextensible Euler-Bernoulli beam in the primary, subharmonic, and superharmonic resonances. The multiple scales method is applied to obtain the linear equations, which can easily analyze the high order mode dynamic behavior of beam. The stable vibration regions of the feedback gains and time-delays are calculated by using the stability conditions of resonances. The unstable behavior of the non-linear dynamic system can be changed by the action of the spring and damping controllers. The effects of the feedback gains and time-delays on the steady state responses of the dynamic beam structure system are investigated.

2. Formulation and analysis

2.1. Primary resonance

The non-linear resonance of bridge can take place under the moving load and wind, which can lead to serious accidents. The bridge is assumed to be a simply-simply supported beam and inextensible. It is also assumed to follow the Euler–Bernoulli beam theory and possess uniform cross-sectional area, where rotary inertia terms and shear deformations are negligible (see Fig. 1). The non-dimensional form of equations of motion and boundary conditions of the beam can be obtained from [23-25]:

$$\ddot{w} + \varepsilon\eta\dot{w} + w^{(4)} = \varepsilon\alpha w'' \int_0^1 w'^2 dx + \varepsilon f(x)\cos(\Omega t) + \varepsilon M[H''(x - x_1) - H''(x - x_2)], \quad (1)$$

$$w = w'' = 0 \text{ at } x = 0, 1, \quad (2)$$

where the primes and overdots indicate the derivatives with respect to the position x and time t , respectively. w is the nondimensional bending vibration, η the damping coefficient, ε the small parameter, Ω the excitation frequency, and ε the perturbation parameter, H the Heaviside unit doublet function. $\alpha = WA/I$, where W is a characteristic transverse displacement, A the area of cross-section, and I the inertia momentum. l is the length of beam. $x = X/l$, $x_1 = X_1/l$, $x_2 = (X_1 + X_2)/l$. The active control torque M is expressed as:

$$M = ka[-u(t - \tau_1) + aw'(x_1, t) - aw'(x_2, t)] + ca[-\dot{u}(t - \tau_2) + a\dot{w}'(x_1, t) - a\dot{w}'(x_2, t)], \quad (3)$$

where $k = k_0 l^4/EIW$, k_0 is the stiffness of the spring. $c = c_0 l^4/EIW$, c_0 is the coefficient of the damping controller. $a = \bar{a}/l$. $u(t)$ is the displacement of spring caused by time-delayed active controller. τ_1 and τ_2 are the time-delays, which are produced by the displacement time-delayed and velocity time-delayed controllers. That $u(t) = 0$ in Eq. (3) is corresponding to a passive structure control system.

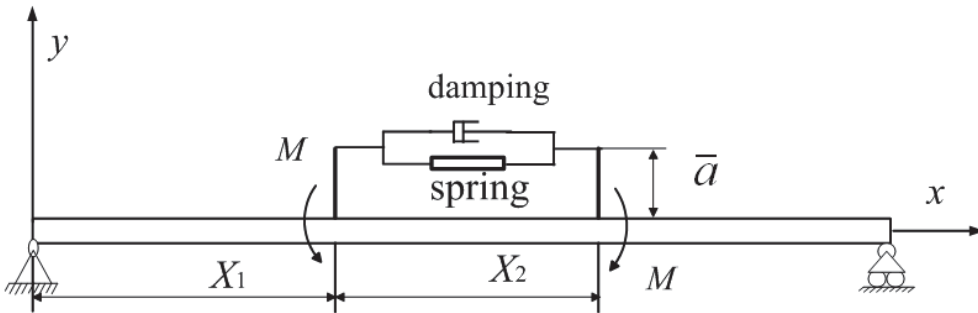


Fig. 1. Dynamic beam structure

The beam bending vibration w can be expanded by order of ε as [22-24]:

$$w(x, t; \varepsilon) = w_0(x, T_0, T_1) + \varepsilon w_1(x, T_0, T_1) + \dots, \quad (4)$$

where T_0 and T_1 are the time scales. Substituting expression (4) into the partial differential Eq. (1) and boundary conditions (2), and separating terms at orders of ε , one yields:

$$D_0^2 w_0 + w_0^{(4)} = 0, \quad (5)$$

$$w_0 = w_0'' = 0 \text{ at } x = 0, 1, \quad (6)$$

$$D_0^2 w_1 + w_1^{(4)} = -2D_0 D_1 w_0 - \eta D_0 w_0 + \alpha w_0'' \int_0^1 w_0'^2 dx + f \cos(\Omega T_0) + M[H''(x - x_1) - H''(x - x_2)], \quad (7)$$

$$w_1 = w_1'' = 0 \text{ at } x = 0, 1, \quad (8)$$

where $D_0 = \frac{\partial}{\partial T_0}$, $D_1 = \frac{\partial}{\partial T_1}$. The Galerkin approximation can be used to represent w as a series of products of spatial functions of x and time-dependent functions as:

$$w(x, t) = \sum_{n=1}^N w_n(x, t) = \sum_{n=1}^N \varphi_n(x) q_n(t), \quad (9)$$

where φ_n is the n th eigenfunctions of a linear uniform beam, $\varphi_n = \sqrt{2} \sin(n\pi x)$. q_n is the n th generalized time-dependent coordinates.

$u(t)$ is the displacement of spring and damping caused by delayed feedback active controllers. When the functions are designed as displacement and velocity time-delayed feedback control, the simplest form with time-delays of $u(t - \tau_1)$ and $\dot{u}(t - \tau_2)$ can be written as the linear functions of $q_k(t)$ and $\dot{q}_k(t)$. By introducing the time-delays, $u(t - \tau_1)$ and $\dot{u}(t - \tau_2)$ can be expressed as:

$$u(t - \tau_1) = g q_k(t - \tau_1), \quad \dot{u}(t - \tau_2) = g \dot{q}_k(t - \tau_2), \quad (10)$$

where g is feedback gain parameter. $g > 0$, $g < 0$, and $g = 0$ point to positive, negative and no feedback, respectively.

The solution of linear Eq. (5) is:

$$w_{0k} = \varphi_k(x) [A_k(T_1) e^{i\omega_k T_0} + \bar{A}_k(T_1) e^{-i\omega_k T_0}], \quad (11)$$

where φ_k is the k th eigenfunction of a linear uniform beam. A_k is the k th approximate amplitude of Eq. (5), which is the function of time T_1 , \bar{A}_k is the conjugate item of A_k . The solvability condition demands that the eigenfunctions are orthogonal, i.e.:

$$\langle \varphi_n, \varphi_m \rangle = \int_0^1 \varphi_n(s) \varphi_m(s) ds = \delta_{nm}, \quad (12)$$

δ_{nm} is the Kronecker delta. For the primary resonance, the frequency of force is close to the natural frequency of system. One can get:

$$\Omega = \omega_k + \varepsilon \sigma, \quad (13)$$

where σ is the detuning parameter. ω_k is the circular frequency of k th order mode. By substituting Eqs. (11), (12) and (13) into (7), applying the conditions of Eq. (8), and letting $A_k = \frac{1}{2} a_k e^{i\beta_k}$, the secular terms which should be equal to zero turn into:

$$2i\omega_k \left(\frac{1}{2} D_1 a_k + \frac{1}{2} i a_k D_1 \beta_k \right) + \frac{1}{2} i \eta \omega_k a_k - \frac{3}{8} g_{kk} a_k^3 - F_k e^{i(\sigma T_1 - \beta_k)} + M_{1k} a_k + i\omega_k M_{2k} a_k - g_{1k} a_k e^{-i\omega_k \tau_1} - i g_{2k} a_k e^{-i\omega_k \tau_2} = 0, \quad (14)$$

where $F_k = \frac{1}{2} \int_0^1 f \varphi_k(x) dx$, $g_{kk}(\omega_k) = \alpha \langle \varphi_k(x), [\varphi_k'(x) \int_0^1 \varphi_k'^2 dx] \rangle$,

$$g_{1k} = \frac{1}{2}kag [\varphi'_k(x_2) - \varphi'_k(x_1)], \quad g_{2k} = \frac{1}{2}\omega_k cag [\varphi'_k(x_2) - \varphi'_k(x_1)],$$

$$M_{1k} = \frac{1}{2}ka^2[\varphi'_k(x_1) - \varphi'_k(x_2)]^2, \quad M_{2k} = \frac{1}{2}ca^2[\varphi'_k(x_1) - \varphi'_k(x_2)]^2.$$

Let $\gamma_k = \sigma T_1 - \beta_k$. Separating the real and imaginary parts of Eq. (14), we can get:

$$D_1 a_k = -\mu_k a_k + \frac{F_k}{\omega_k} \sin(\gamma_k), \tag{15}$$

$$a_k D_1 \beta_k = \sigma_k a_k + v_k a_k^3 + \frac{F_k}{\omega_k} \cos(\gamma_k), \tag{16}$$

where $\mu_k = \frac{1}{2}\eta + \frac{\omega_k M_{2k} + g_{1k} \sin \omega_k \tau_1 - g_{2k} \cos \omega_k \tau_2}{\omega_k}$, $\sigma_k = \sigma - \frac{M_{1k} - g_{1k} \cos \omega_k \tau_1 - g_{2k} \sin \omega_k \tau_2}{\omega_k}$, $v_k = \frac{3g_{kk}}{8\omega_k}$.

The nominal damping and detuning coefficients are expressed as the function of time-delays and feedback gains.

By setting $D_1 a_k = D_1 \beta_k = 0$, the steady-state response in the system can be expressed as:

$$-\mu_k a_k + \frac{F_k}{\omega_k} \sin(\gamma_k) = 0, \tag{17}$$

$$\sigma_k a_k + v_k a_k^3 + \frac{F_k}{\omega_k} \cos(\gamma_k) = 0. \tag{18}$$

From Eqs. (17) and (18), the frequency response equation of the system can be obtained by:

$$(\mu_k a_k)^2 + [\sigma_k a_k + v_k a_k^3]^2 = \left(\frac{F_k}{\omega_k}\right)^2. \tag{19}$$

The peak amplitude at primary resonance, obtained from Eq. (19), can be expressed as:

$$a_{k\max} = \frac{F_k}{|\mu_k| \omega_k}. \tag{20}$$

For the purpose of comparison, the equation of motion for the nonlinear primary oscillator without attached controllers can be written as:

$$\ddot{w} + \varepsilon \eta \dot{w} + w^{(4)} = \varepsilon a w'' \int_0^1 w'^2 dx + \varepsilon f(x) \cos(\Omega t). \tag{21}$$

The peak amplitude at primary resonance without the controllers is written as:

$$\bar{a}_{k\max} = \frac{F_k}{|\bar{\mu}_k| \omega_k}. \tag{22}$$

As the response amplitude cannot be found analytically for a nonlinear system, a different method has to be developed here to study the performance of vibration controllers. In suppressing the primary resonance vibrations of the nonlinear oscillator, the performance of the vibration controller will be examined in the paper by defining a ratio. An attenuation ratio of nonlinear response is defined by the ratio of the amplitude of nonlinear vibrations with and without the controllers. The attenuation ratio, denoted by r_k , can be expressed as:

$$r_k = \frac{a_k}{\bar{a}_k}. \tag{23}$$

The attenuation ratio of the peak amplitude of primary resonance response is also defined by the ratio of the peak amplitude of primary resonance vibrations of the nonlinear system with and without the controllers. The attenuation ratio R_k is [25]:

$$R_k = \frac{a_{k\max}}{\bar{a}_{k\max}} = \frac{|\bar{\mu}_k|}{|\mu_k|}. \quad (24)$$

As can be seen from the definition given by Eq. (24), under a fixed value of the amplitude of excitation, a small value of the attenuation ratio R_k indicates a large reduction in the nonlinear vibrations of the nonlinear system.

For simplicity, one expresses the time-delays in the form $\tau_1 = \tau$ and $\tau_2 = \phi + \omega_k \tau$. The parameters μ_k become [10]:

$$\mu_k = \frac{1}{2}\eta + \frac{\omega_k M_{2k} + g_{1k} \sin \omega_k \tau - g_{2k} \cos(\omega_k \tau + \phi)}{\omega_k}. \quad (25)$$

As the phase of velocity is ahead of displacement $\pi/2$, the phase difference ϕ can be assumed as $\pi/2$. Eq. (25) can be expressed as:

$$\mu_k = \frac{1}{2}\eta + \frac{\omega_k M_{2k} + G_k \sin \omega_k \tau}{\omega_k}, \quad (26)$$

where $G_k = g_{1k} + g_{2k}$. The attenuation ratio R_k , can be expressed as:

$$R_k = \frac{a_{k\max}}{\bar{a}_{k\max}} = \frac{1}{1 + \left| \frac{2\omega_k M_{2k} + 2G_k \sin \omega_k \tau}{\eta \omega_k} \right|}. \quad (27)$$

2.2. Determining of control parameters for the nonlinear vibration system

The primary resonance response amplitude is determined by the real solution a_k^2 of Eq. (19). One or three real solutions can be obtained. Three real solutions exist between two points of vertical tangents (saddle-node bifurcation), which are determined by differentiation of Eq. (19) implicitly with respect to a_k^2 . This leads to the condition:

$$\sigma_k^2 + 4v_k a_k^2 \sigma_k + \mu_k^2 + 3v_k^2 a_k^4 = 0. \quad (28)$$

Its solutions are:

$$\sigma_k^\pm = -2v_k a_k^2 \pm (v_k^2 a_k^4 - \mu_k^2)^{1/2}. \quad (29)$$

For the case of $v_k^2 a_k^4 - \mu_k^2 > 0$, three real and positive solutions a_k^2 of Eq. (29) exist in an interval $\sigma < \sigma_k < \sigma^+$. In the limit $v_k^2 a_k^4 \rightarrow \mu_k^2$, this interval shrinks to the point $\sigma_k \rightarrow 2v_k a_k^2$. The critical force amplitude obtained from Eq. (19) can be written as:

$$F_{kc} = \omega_k \mu_k (2\mu_k / v_k)^{1/2}. \quad (30)$$

For $F > F_{kc}$, there are three solutions. While $F < F_{kc}$, there is only one.

The stability of the solutions can be determined by the eigenvalues of the corresponding Jacobian matrix of Eq. (17) and (18). The corresponding eigenvalues are the roots of [10]:

$$\lambda^2 + 2\mu_k \lambda + \mu_k^2 + (\sigma_k + v_k a_k^2)(\sigma_k + 3v_k a_k^2) = 0. \quad (31)$$

As can be seen from Eq. (31), the sum of the two eigenvalues is $-2\mu_k$. $\mu_k < 0$, which means that at least one of the eigenvalues will always have a positive real part. The system will be unstable. $\mu_k = 0$, there is a pair of purely imaginary eigenvalues for the system. The system is also unstable and a Hopf bifurcation may occur. If $\mu_k = 0$, the sum of two eigenvalues is always negative, at least one of the two eigenvalues will always have a negative real part. The other eigenvalues is zero when:

$$f(\sigma_k) = \mu_k^2 + (\sigma_k + v_k a_k^2)(\sigma_k + 3v_k a_k^2) = 0. \quad (32)$$

From Eq. (26), we can find that μ_k is the function of the feedback gains, damping parameters and time-delays. As a larger positive value, μ_k is prone to have the negative real parts of characteristic roots, M_{2k} should have a larger value in order to improve the control performance by selecting proper coefficients of damping coefficient c , arm of force a and the coordinates of installation of the controller.

Based on the analyses mentioned above, the sufficient conditions to guarantee the system stability are:

$$f(\sigma_k) > 0, \mu_k > 0. \quad (33)$$

If there is no real solution to the Eq. (32), the sufficient conditions to guarantee the system stability are:

$$v_k^2 a_k^4 - \mu_k^2 < 0, \mu_k > 0. \quad (34)$$

Letting $\mu_k^2 \geq v_k^2 a_{k\max}^4 \geq v_k^2 a_k^4$, we can find that the inequalities satisfy the Eq. (34) and get:

$$G_k \sin \omega_k \tau \geq \omega_k \left[\sqrt[3]{\frac{|v_k| F_k^2}{\omega_k^2}} - M_{2k} - \frac{1}{2} \eta \right], \mu_k > 0. \quad (35)$$

If there are two real solutions to the Eq. (32), the solutions are the same as Eq. (29). Let:

$$\sigma_{k-} \leq -2v_k a_{k\max}^2 - (v_k^2 a_{k\max}^4 - \mu_k^2)^{1/2} \leq \sigma_{k-}^- \text{ or } \sigma_{k+} \geq (v_k^2 a_{k\max}^4 - \mu_k^2)^{1/2} \geq \sigma_{k+}^+. \quad (36)$$

We can find that $f(\sigma_{k\pm}) > 0$. The sufficient conditions of guaranteeing the system stability can be written as:

$$G_k \cos \omega_k \tau \leq \omega_k [M_{1k} - \sigma \omega_k - 2v_k a_{k\max}^2 - (v_k^2 a_{k\max}^4 - \mu_k^2)^{1/2}] \text{ and } \mu_k > 0, \quad (37)$$

$$G_k \cos \omega_k \tau \geq \omega_k [M_{1k} - \sigma \omega_k + (v_k^2 a_{k\max}^4 - \mu_k^2)^{1/2}] \text{ and } \mu_k > 0.$$

For the fixed time-delay, the range of the feedback gains can be obtained based on the stable conditions of the primary response resonances. For the given feedback gains of the controllers, the time delay can be designed for the requirement of the suppression of the nonlinear systems.

2.3. Subharmonic resonance

In the case of subharmonic resonance, the excitation is:

$$Q(x) = f(x). \quad (38)$$

Substituting Eqs. (4) and (38) into Eq. (1) and equating coefficients of like powers of ϵ , we have:

$$D_0^2 w_0 + w_0^{(4)} = Q(x) \cos(\Omega T_0), \quad (39)$$

$$w_0 = w_0'' = 0 \text{ at } x = 0, 1, \quad (40)$$

$$D_0^2 w_1 + w_1^{(4)} = -2D_0 D_1 w_0 - \eta D_0 w_0 + \alpha w_0'' \int_0^1 w_0'^2 dx + M[H''(x - x_1) - H''(x - x_2)], \quad (41)$$

$$w_1 = w_1'' = 0 \text{ at } x = 0, 1. \quad (42)$$

The solution to Eq. (39) can be written as:

$$w_0 = \varphi_k(x)[A_k(T_1)e^{i\omega_k T_0} + B_k e^{i3(\omega_k T_0 + \sigma T_1)} + cc], \quad (43)$$

where $B_k = \frac{F_k}{2(\omega_k^2 - \Omega^2)}$, $F_k = \int_0^1 \varphi_k(x)Q(x)dx$. cc are the conjugate items of the solution. Substituting (43) into Eq. (41), we can obtain the equation of secular terms:

$$i2\omega_k D_1 A_k + i\eta \omega_k A_k - 3g_{kk} A_k^2 \bar{A}_k - 6g_{kk} A_k B_k \bar{B}_k - 3g_{kk} \bar{A}^2 B_k e^{i(3\sigma T_1)} + 2M_{1k} A_k + i2M_{2k} A_k - 2g_{1k} A_k e^{-i\omega_k \tau_1} - i2g_{2k} A_k e^{-i\omega_k \tau_2} = 0. \quad (44)$$

Let $A_k = \frac{1}{2} a_k e^{i\beta_k}$, $\gamma_k = 3\sigma T_1 - 3\beta_k$. After substituting them into Eq. (44) and separating the real and imaginary parts, we have:

$$D_1 a_k = -\mu_k a_k + \frac{3g_{kk} B_k a_k^2}{4\omega_k} \sin(\gamma_k), \quad (45)$$

$$a_k D_1 \gamma_k = 3 \left[\sigma_k a_k + \Gamma_k a_k + v_k a_k^3 + \frac{3g_{kk} B_k a_k^2}{4\omega_k} \cos(\gamma_k) \right], \quad (46)$$

where $\Gamma_k = \frac{3g_{kk} B_k \bar{B}_k}{\omega_k}$. In the case of $D_1 a_k = a_k D_1 \gamma_k = 0$, the steady state response corresponds to the solutions of:

$$-\mu_k a_k + \frac{3g_{kk} B_k a_k^2}{4\omega_k} \sin(\gamma_k) = 0, \quad (47)$$

$$\sigma_k a_k + \Gamma_k a_k + v_k a_k^3 + \frac{3g_{kk} B_k a_k^2}{4\omega_k} \cos(\gamma_k) = 0. \quad (48)$$

By eliminating γ_k from Eqs. (47) and (48), the frequency response equation is obtained as:

$$(u_k a_k)^2 + (\sigma_k a_k + \Gamma_k a_k + v_k a_k^3)^2 = \left(\frac{3g_{kk} B_k a_k^2}{4\omega_k} \right)^2. \quad (49)$$

There are two possibilities: either a trivial solution a_k , or non-trivial solutions, which can be acquired by:

$$(\mu_k)^2 + (\sigma_k + \Gamma_k + v_k a_k^2)^2 = \left(\frac{3g_{kk} B_k}{4\omega_k} \right)^2 a_k^2. \quad (50)$$

Solving Eq. (50), we can get:

$$a_k^2 = p_k \pm \sqrt{p_k^2 - q_k}, \quad (51)$$

where $p_k = \left(\frac{3g_{kk}B_k}{16v_k\omega_k}\right)^2 - \frac{2(\sigma_k + \Gamma_k)}{4v_k^2}$, $q_k = \frac{u_k^2 + (\sigma_k + \Gamma_k)^2}{v_k^2}$. Because $q_k > 0$, the condition of the positive solution is $p_k > 0$ and $p_k^2 \geq q_k$. The necessary condition of 1/3 subharmonic resonance can be written as:

$$\frac{1}{2} \left[\frac{3g_{kk}B_k}{4\omega_k} - 2(\sigma_k + \Gamma_k) \right] \geq \sqrt{u_k^2 + (\sigma_k + \Gamma_k)^2}. \quad (52)$$

The steady-state solutions of subharmonic resonance response are determined by the eigenvalues of the characteristic equation, which are the roots of:

$$\lambda^2 + 2\mu_k\lambda + \mu_k^2 + (\sigma_k + \Gamma_k + v_ka_k^2)(\sigma_k + \Gamma_k + 3v_ka_k^2) = 0. \quad (53)$$

2.4. Superharmonic resonance

For the case of superharmonic resonance, the solution of Eq. (39) is:

$$w_0 = \phi_k(x) \left[A_k(T_1)e^{i\omega_k T_0} + B_k e^{i\frac{\omega_k(T_0 + \sigma T_1)}{3}} + cc \right]. \quad (54)$$

By substituting Eq. (54) into Eq. (41), the equation of secular terms is written as:

$$i2\omega_k D_1 A_k + i\eta\omega_k A_k - 3g_{kk}A_k^2 \bar{A}_k - 6g_{kk}A_k B_k \bar{B}_k - g_{kk}B_k^3 e^{i(\sigma T_1)} + 2M_{1k}A_k + i2M_{2k}A_k - 2g_{k1}A_k e^{-i\omega_k T_1} - i2g_{2k}A_k e^{-i\omega_k T_2} = 0. \quad (55)$$

Let $A_k = \frac{1}{2}a_k e^{i\beta_k}$, $\gamma_k = \sigma T_1 - \beta_k$. Substituting them into Eq. (55) and separating the real and imaginary parts, we have:

$$D_1 a_k = -\mu_k a_k + \frac{g_{kk}B_k^3}{\omega_k} \sin(\gamma_k), \quad (56)$$

$$a_k D_1 \gamma_k = \sigma_k a_k + \Gamma_k a_k + v_k a_k^3 + \frac{g_{kk}B_k^3}{\omega_k} \cos(\gamma_k). \quad (57)$$

In the case of $D_1 a_k = a_k D_1 \gamma_k = 0$, the fixed points of this system are expressed as:

$$-\mu_k a_k + \frac{g_{kk}B_k^3}{\omega_k} \sin(\gamma_k) = 0, \quad (58)$$

$$\sigma_k a_k + \Gamma_k a_k + v_k a_k^3 + \frac{g_{kk}B_k^3}{\omega_k} \cos(\gamma_k) = 0. \quad (59)$$

Eliminating γ_k from Eqs. (58) and (59), we obtain the frequency response equation:

$$(\mu_k a_k)^2 + (\sigma_k a_k + \Gamma_k a_k + v_k a_k^3)^2 = \left(\frac{g_{kk}B_k^3}{\omega_k} \right)^2. \quad (60)$$

The peak amplitude of the superharmonic resonance, obtained from Eq. (60), is given by:

$$a_{k\max} = \frac{|g_{kk}B_k^3|}{\omega_k |\mu_k|}. \quad (61)$$

For the purpose of comparison, the peak amplitude of the superharmonic resonance without control can be written as:

$$\bar{a}_{k\max} = \frac{|g_{kk}B_k^3|}{\omega_k|\bar{\mu}_k|}. \quad (62)$$

The attenuation ratio of the peak amplitude of superharmonic resonance response is defined by the ratio of the peak amplitude of superharmonic resonance vibrations of the nonlinear system with and without the controllers. The attenuation ratio R_k is [26]:

$$R_k = \frac{a_{k\max}}{\bar{a}_{k\max}} = \frac{|\bar{\mu}_k|}{|\mu_k|}. \quad (63)$$

The characteristic equation of superharmonic resonance response is:

$$\lambda^2 + 2\mu_k\lambda + \mu_k^2 + (\sigma_k + \Gamma_k + v_k a_k^2)(\sigma_k + \Gamma_k + 3v_k a_k^2) = 0. \quad (64)$$

The sufficient conditions of guaranteeing the system stability are:

$$f(\sigma_k) = \sigma_k^2 + (2\Gamma_k + 4v_k a_k^2)\sigma_k + \mu_k^2 + (\Gamma_k + v_k a_k^2)(\Gamma_k + 3v_k a_k^2) > 0, \mu_k > 0. \quad (65)$$

If there is no real solution for the Eq. (65), the sufficient conditions of guaranteeing the system stability are:

$$v_k^2 a_k^4 - \mu_k^2 < 0, \mu_k > 0. \quad (66)$$

Let $\mu_k^2 \geq v_k^2 a_{k\max}^4 \geq v_k^2 a_k^4$. We can find that the inequalities satisfy the Eq. (66). We can get:

$$G_k \sin \omega_k \tau \geq \omega_k \left[\sqrt[3]{\frac{|v_k| F_s^2}{\omega_k^2}} - M_{2k} - \frac{1}{2} \eta \right], \mu_k > 0, \quad (67)$$

where $F_s = g_{kk}B_k^3$.

If there are two real solutions for the $f(\sigma_k) = 0$, the solutions can be written as:

$$\sigma_k^\pm = -\Gamma_k - 2v_k a_k^2 \pm (v_k^2 a_k^4 - \mu_k^2)^{1/2}. \quad (68)$$

Let:

$$\begin{aligned} \rho_{k-} &\leq -\Gamma_k - 2v_k a_{k\max}^2 - (v_k^2 a_{k\max}^4 - \mu_k^2)^{1/2} \leq \sigma_k^- \text{ or} \\ \sigma_{k+} &\geq -\Gamma_k + (v_k^2 a_{k\max}^4 - \mu_k^2)^{1/2} \geq \sigma_k^+. \end{aligned} \quad (69)$$

We can find that $f(\sigma_{k\pm}) > 0$. The sufficient conditions of guaranteeing the system stability can be written as:

$$\begin{aligned} G_k \cos(\omega_k \tau) &\leq \omega_k \left[M_{1k} - \sigma \omega_k - \Gamma_k - 2v_k a_{k\max}^2 - (v_k^2 a_{k\max}^4 - \mu_k^2)^{\frac{1}{2}} \right] \text{ or} \\ G_k \cos(\omega_k \tau) &\geq \omega_k \left[M_{1k} - \sigma \omega_k - \Gamma_k + (v_k^2 a_{k\max}^4 - \mu_k^2)^{\frac{1}{2}} \right] \text{ and } \mu_k > 0. \end{aligned} \quad (70)$$

For the fixed time-delay, the range of the feedback gains can be obtained based on stable conditions of the superharmonic response resonances. For the given feedback gains of the

controllers, the time delay can be designed for the requirement of the suppression of the nonlinear systems.

3. Case study

This section illustrates the effect of the feedback gains and time-delays on the non-linear dynamical behavior of the controlled system. The variation of attenuation ratio with time-delays and feedback gains is also discussed. The results are shown in a set of figures.

3.1. The response analysis of first mode

3.1.1. The primary resonance of first mode

The beam's non-dimensional geometrical parameters l , a , x_1 and x_2 are 1, 0.05, 0.2, and 0.4, respectively. $\alpha = 0.2$, $\eta = 0.25$. The coefficient of elasticity of the spring is 80. The feedback gain of the spring is 5. The coefficient of damping controllers is 0.3.

Figs. 2 and 3 show the variation of μ_1 and μ_2 with the time-delay τ for $\Phi = \pi/2$ and $M_{21} = 0, 0.3$ and 0.6 , respectively. We can find that μ_k varies with value of time-delays and M_{2k} . As larger μ_k is prone to have the negative real parts of characteristic roots, M_{2k} should have a larger value in order to improve the control performance. Fig. 2 shows that the value of μ_1 of first order mode is larger than $\bar{\mu}_1$ in the range $0.32 < \tau < 0.63$, which means that the attenuation ratio R is smaller than 1. Fig. 3 shows that the value of μ_2 of second order mode is larger than $\bar{\mu}_2$ in the range $0.08 < \tau < 0.159$. From Figs. 2 and 3, we also find that the value of time-delay for large μ_k of lower mode is larger than the higher one. For the fixed feedback gains and M_{2k} , different control purposes can be easily achieved by the right choice of the time-delays. μ_k should have a larger value in order to improve the suppression control performance.

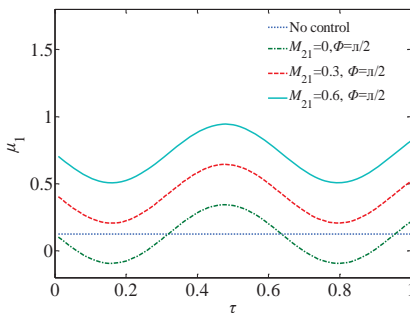


Fig. 2. Variation of μ_1 with the time-delays of the first mode

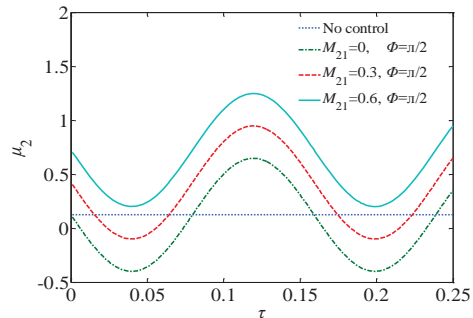


Fig. 3. Variation of μ_2 with the time-delays of the second mode

Figs. 4 and 5 show the variation of attenuation ratio R_k with time-delays and feedback gains. As can be seen from the figures, under a fixed value of the amplitude of excitation, a small value of the attenuation ratio R_k indicates a large reduction in the nonlinear vibrations of the nonlinear primary system. A selected value of time-delay can get a large positive value of μ_k and small attenuation ratio R_k . From Eq. (35), we can find appropriate value of feedback controllers, which can satisfy the stability conditions. The proper selection of the feedback gains and the time-delays can guarantee the stability of the non-linear system. For the fixed feedback gains and M_{2k} , the good suppression control for the nonlinear vibration of beams can obtain from the right choice of the time-delays easily.

Fig. 6 shows the primary response curves of first mode for three different sets of the time-delays. There is no jump and hysteresis phenomenon when $\tau = \pi/3\omega$, $\Phi = \pi/2$ and no control.

This suggests that saddle node bifurcation and jump phenomenon can be eliminated by certain values of the time-delays. Three solutions exist in a region of coexistence of $\tau = \pi/2\omega$, $\Phi = \pi/2$. The bending of the frequency response curves is responsible for a jump phenomenon. Moreover, the peak amplitude of the primary resonance response at $\tau = \pi/3\omega$, $\Phi = \pi/2$ is smaller than that in the other two cases.

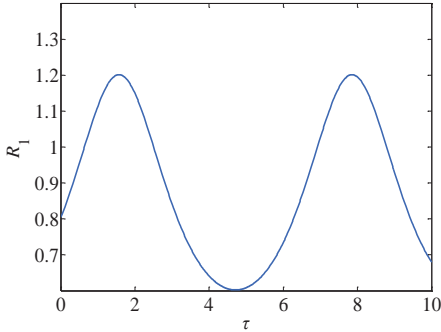


Fig. 4. Variation of attenuation ratio R_1 with the time-delays

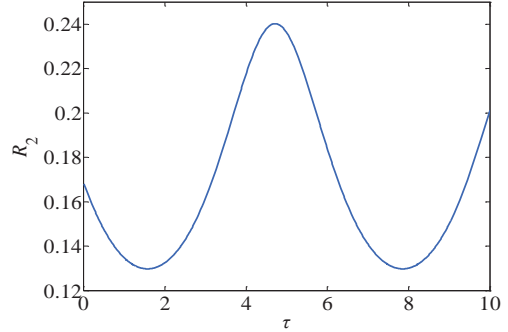


Fig. 5. Variation of attenuation ratio R_2 with the time-delays

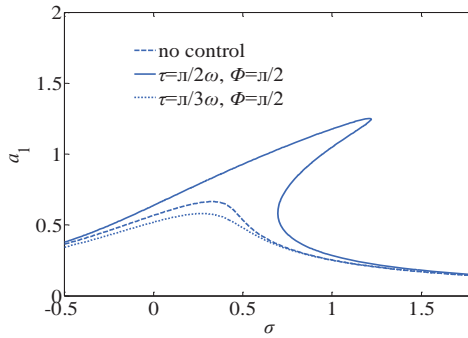


Fig. 6. First mode frequency-response curves of primary resonance for three sets of the time-delays

3.1.2. Subharmonic resonance

The excitation amplitude of the subharmonic resonance is $B_k = 0.30$. The curve of the subharmonic resonance with time-delays is shown as Fig. 7.

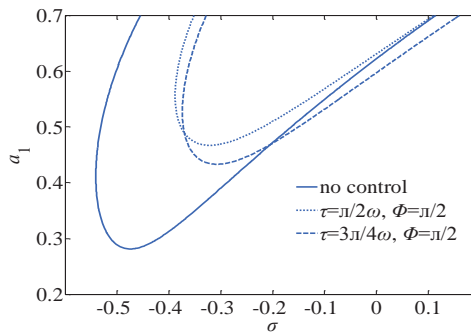


Fig. 7. First mode frequency-response curves of subharmonic resonance for three sets of the time-delays

Fig. 7 shows the regions where subharmonic response exists for three different sets of time-

delays. The time-delays can change the range of the occurrence of subharmonic resonance. It is noted that the regions for the existence of subharmonic responses are different. This shows that there always exist certain regions of the time-delays where subharmonic response does not exist.

3.1.3. Superharmonic resonance

The excitation amplitude of the superharmonic resonance of beam is $B_k = 0.38$. The curve of the superharmonic resonance with time-delays is shown in Fig. 8. For the superharmonic resonance response, the suitable choice of the time-delays and feedback gains can also improve the control performance. Moreover, the occurrence of saddle-node bifurcation, jump and hysteresis phenomena can be delayed or eliminated.

Fig. 8 shows the first mode frequency-response curves of superharmonic resonance response for three sets of the time-delays. $\tau = \pi/2\omega$, and $\Phi = 0$ mean that the system is stable. When $\tau = \pi/\omega$, $\Phi = \pi/2$ and no control, jumps exist. This suggests that saddle node bifurcation and jump phenomenon can be eliminated by certain values of the time-delays. Thus, the control performance can be enhanced by the optimal selection of the feedback gains and time-delays.

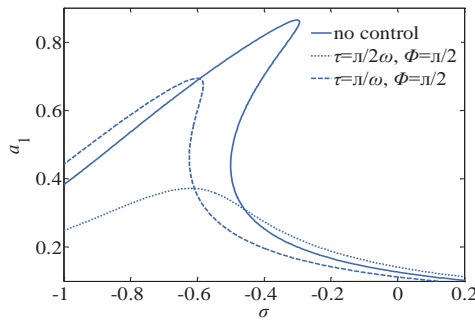


Fig. 8. Frequency-response curves of superharmonic resonance for three sets of the time-delays

3.2. Second mode response analysis

3.2.1. Primary resonance of second mode

The beam's geometry non-dimensional parameters l , a , x_1 and x_2 are 1, 0.05, 0.2, and 0.4, respectively. $\alpha = 0.2$. The coefficient of elasticity is 80; the damping ratio is 0.3; the feedback gain is 10. The parameters η and ω are $\eta = 0.25$, $\omega = 4\pi^2$. The curves of the primary resonance with time-delays are shown in Figs. 9-10.

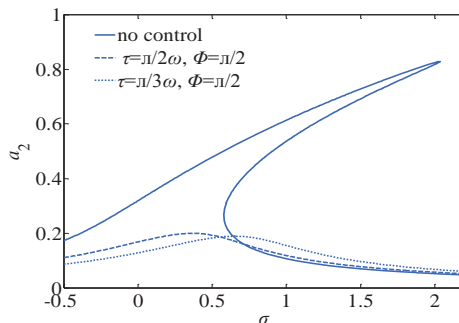


Fig. 9. Second mode frequency-response curves of primary resonance for three sets of the time-delays

Fig. 9 shows the second mode frequency-response curves of primary response for three sets of the time-delays. $\tau = \pi/3\omega$, $\Phi = \pi/2$ and $\tau = \pi/2\omega$, $\Phi = \pi/2$ mean that the system is stable while jumps occur in the frequency response without control. This indicates that saddle-node bifurcation and jump phenomenon can be eliminated by selecting suitable time-delays.

3.2.2. Subharmonic resonance

The excitation amplitude of the subharmonic resonance of beam is $B_k = 0.10$. The curve of the subharmonic resonance with respect to time-delays is shown in Fig. 10.

Fig. 10 shows the regions where subharmonic response exists for the three different sets of time-delays. The time-delays can change the range of the occurrence of subharmonic resonance. It is noted that the regions for the existence of subharmonic responses are different. This shows that there always exist certain ranges of the time-delays where subharmonic response does not exist.

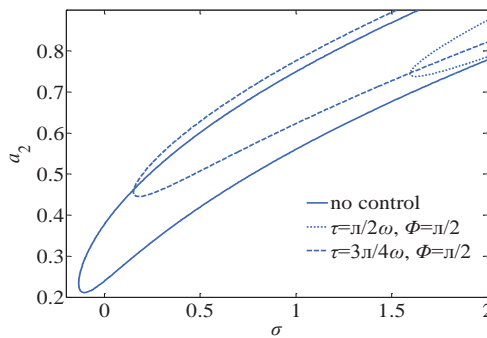


Fig. 10. Second mode frequency-response curves of subharmonic resonance for three sets of the time-delays

4. Conclusions

The primary, subharmonic, and superharmonic resonances of the first and second modes of the Bernoulli-Euler beam are studied with spring and damping delayed feedback controllers. The results show that the proper selection of the feedback gains and time-delays can enhance the control performance, but the inappropriate feedback gains and time-delays may lead to instability of the system. By using the joint function of spring and damping delayed feedback control strategy, the bifurcation can be eliminated or the bifurcation point's position can be changed.

Time-delays feedback controllers can be utilized to suppress the dynamic behaviors of nonlinear systems. Most of the vibrational energy of the nonlinear oscillator is transferred to the controllers through coupling spring and damping. The vibration controllers can effectively suppress the amplitude of oscillations of the nonlinear oscillator. Hence, by properly choosing the feedback gains and time-delays of the linked spring and damping controllers, the primary and superharmonic resonance response of the nonlinear oscillator can be reduced to relatively small amplitude. The control performance of beam's nonlinear vibration can be improved with these delayed controllers by selecting appropriate time delays and feedback gains.

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