## 958. Nonlinear vibration characteristics of a rotor system with pedestal looseness fault under different loading conditions

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**Abstract.** Taking a single-span rotor system with two discs as the research object, the pedestal looseness fault is simulated by a piecewise linear spring-damper model, and this model is combined with the FE (finite element) model of the rotor system. Two different loading conditions are determined based on API Standard 617 and the spectrum cascades, rotor orbits and Poincaré maps are used to analyze the influences of the stiffness of non-loosened bolts, looseness clearance and rotating speed on the dynamic characteristics of the system. The results show that different bifurcation forms, multiple periodic, quasi-periodic and chaotic motions can be observed under two loading conditions, and the system motion is more complicated under the second loading condition. The results will provide theoretical references for fault diagnosis, dynamic design, and safe operation of the rotor-bearing system.

Keywords: nonlinear vibration, rotor system, pedestal looseness, finite element method.

#### 1. Introduction

In the rotating machinery, the pedestal looseness is a common fault. The poor installation quality and the long-term vibration will lead to the looseness between the mechanical components. Especially, the serious looseness fault may induce other faults such as rub-impact fault between the rotor and the stator, even may lead to disastrous accidents [1]. Therefore, the research on pedestal looseness fault is significant in engineering practice for the safe operation of rotating machinery, the extension of service life and the improvement of its work efficiency.

Many researchers have studied the pedestal looseness by using piecewise linear models [2-7]. Muszynska et al. [2, 3] proposed a bi-linear model of a rotating machine with one loose pedestal. Their numerical results shows the synchronous and subsynchronous fractional components of the response, which were verified by the experiments. Chu et al. [4] analyzed the vibration characteristics of a rotor-bearing system with pedestal looseness by building a non-linear mathematical model. The stability of these periodic solutions was discussed by using the shooting method and the Floquet theory. In some cases, the pedestal looseness could result in 1/2 fractional harmonic and multiple harmonic components such as 1X/2, 2X, 3X, etc., and could occur intermittently under some special conditions [5]. Duan et al. [6] proposed a non-linear mathematical model of a rotor system with pedestal looseness, which considers both the variation of the stiffness and the periodic impact action caused by the pedestal looseness. Chen [7] introduced a dynamic model of a coupled rotor-rolling bearing-casing system considering of the coupling faults of the rotor unbalance and the pedestal looseness.

The piecewise linear models have also been proved to well simulate a discontinuous stiffness affected by a bearing clearance, rotor-stator contact, etc. Aiming at a Jeffcott rotor with a piecewise-linear non-linearity at the bearing support, Kim et al. [8] presented a multiple harmonic balance method and analyzed the internal resonant steady state vibration of the rotor system. Hossain et al. [9] studied preloading effects on the bilinear system with clearance and analyzed the nonlinear responses by the changing of the equilibrium positions for a wide range of the excitation frequency, the results show that equilibrium position is an important parameter in the bilinear system with clearance problem, which can show the chaotic and the multiple periodic motion. Considering the nonlinearity of a discontinuous stiffness affected by a radial

clearance between the elastically supported rotor and an elastically supported outer ring, Gonsalves et al. [10] and Karpenko et al. [11] investigated the nonlinear vibration responses of a two degree of freedom rotor dynamic system. Aiming at a rotor system with bearing clearance and stabilizing rods, Karlberg et al. [12] presented a piecewise linear two degrees-of-freedom rotor system model and analyzed the nonlinear vibration responses of the system. Considering the effects of the piecewise stiffness and damping caused by the bearing clearance, Gao et al. [13] established a six-degree-of-freedom (DOF) model of a machine-tool spindle-bearing system and investigated the effect of bearing clearance, bifurcations and routes to chaos of this nonsmooth system.

Recently, many researchers have studied the rotor looseness fault by using the finite element method. Based on the non-linear finite element theory, Ren et al. [14] analyzed the influence of the looseness fault on the non-linear dynamic characteristics of a single-span rotor system with two discs by the simulation. Ma et al. [15] analyzed the effects of the looseness parameters on the system dynamic characteristics by using the finite element method based on the piecewise linear stiffness model. Behzad [16] developed a finite-element code for studying the effects of loose rotating discs on the vibration responses of the rotor-bearing system. The developed finite-element model can numerically give the response of rotors with any number of loose discs at any location with isotropic or orthotropic supports.

Most of previous researches on looseness fault focused on the simple Jeffcott rotor systems using the lumped mass model and only a few works were on the complex rotor system with looseness fault using finite element method (FEM). In this paper, taking a single-span rotor system with two discs as the object, the pedestal looseness is simulated by piecewise linear stiffness and damping models. And then the models are coupled with the finite element model of the rotor system. Based on the two different loading conditions, the effects of the stiffness of non-loosened bolts, looseness clearance, and rotational speed on the dynamic characteristics of the system are investigated by using the spectrum cascades, rotor orbits and Poincaré maps. The results can provide theoretic basis for the fault diagnosis of the pedestal looseness.

#### 2. Dynamic model of a rotor system with pedestal looseness fault

#### 2.1. Finite element model of a rotor system

In order to simulate the dynamic characteristics of the rotor system efficiently, the system is simplified according to the following assumptions:

(a) The shaft is simulated by a Timoshenko beam and the element model is shown in Fig. 1.

(b) The two oil film bearings are identical, which are simulated ideally by linear stiffness and damping models and the effects of the pedestal looseness on the stiffness and damping of the bearing are ignored.

(c) The pedestal looseness locates in the right bearing position, and the stiffness and damping between the pedestal and the base are only considered in the looseness direction.



Fig. 1. Finite element model of shaft section element

Neglecting axial and torsional deformations, the general displacement vector of a beam element for the shaft  $\mathbf{u}_s$  is given as:

(1)

 $\mathbf{u}_{s} = [y_{A} \quad z_{A} \quad \theta_{yA} \quad \theta_{zA} \quad y_{B} \quad z_{B} \quad \theta_{yB} \quad \theta_{zB}]^{\mathrm{T}}.$ 

#### 2.2. Equivalent stiffness model of the loosened pedestal

When one or a few bolts loosen and vibration aggravates, the pedestal and the base may be separated partially. Assuming that the right pedestal looses in vertical direction (y direction), as is shown in Fig. 2. In the figure,  $y_p$  is pedestal displacement and  $\delta_1$  is the looseness clearance.



Fig. 2. Bolt looseness schematic

When  $y_p < 0$ , the pedestal contacts with the base and the equivalent stiffness of the right pedestal  $k_{br}$  is the base stiffness  $k_b$ . When  $0 \le y_p \le \delta_1$ , non-loosened bolts will be in a state of elastic deformation due to the pulling force and  $k_{br}$  is the stiffness of non-loosened bolts  $k_{b1}$ . When  $y_p > \delta_1$ , the loosened bolts and non-loosened bolts will be pulled together, here  $k_{br}$  is  $k_{b1} + k_{b2} - k_{b2}\delta_1/y_p$  ( $k_{b2}$  is the stiffness of loosened bolts). Based on the above analysis, it is clear that  $k_{br}$  is a piecewise function related to  $y_p$  and its expression can be written as:

$$k_{br} = \begin{cases} k_{b}, & y_{p} < 0, \\ k_{b1}, & 0 \le y_{p} \le \delta_{1}, \\ k_{b1} + k_{b2} - k_{b2}\delta_{1}/y_{p}, & y_{p} > \delta_{1}. \end{cases}$$
(2)

Assuming that base stiffness is approximately equal to tensile stiffness of bolts, namely,  $k_b \approx k_{b1} + k_{b2} - k_{b2}\delta_1/y_p$ , then Eq. (2) can be simplified as:

$$k_{br} \approx \begin{cases} k_{b1}, & 0 \le y_p \le \delta_1, \\ k_{b}, & \text{others.} \end{cases}$$
(3)

The equivalent damping of the right pedestal  $c_{br}$  is similar to  $k_{br}$ , and its expression is:

$$c_{br} \approx \begin{cases} c_{b1}, & 0 \le y_p \le \delta_1, \\ c_b, & \text{others.} \end{cases}$$
(4)

#### 2.3. Finite element model of a rotor system with pedestal looseness

Considering the piecewise linear model of the right pedestal looseness [15], the non-linear equation of motion of the rotor-bearing system can be written as:

$$\mathbf{M}\ddot{\mathbf{u}} + \mathbf{D}\dot{\mathbf{u}} + \mathbf{K}\mathbf{u} = \mathbf{F}_{u},\tag{5}$$

where M, D and K denote the mass matrix, damping matrix and stiffness matrix of the whole

system, respectively. The mass matrix **M** includes the rotor mass and pedestal mass; the damping matrix **D** includes the viscous damping ( $C_1$ ), bearing damping and gyroscopic damping; the stiffness matrix **K** includes the rotor stiffness, bearing stiffness and equivalent pedestal stiffness. **u** is the displacement vector and  $F_u$  is the external load vector. In this paper, Rayleigh damping form is adopted to determine the viscous part ( $C_1$ ) and it can be obtained by the following formula:

$$\mathbf{C}_1 = \alpha_1 \mathbf{M} + \beta_1 \mathbf{K},\tag{6}$$

where:

$$\alpha_1 = \frac{60(\omega_{n2}\xi_1 - \omega_{n1}\xi_2)\omega_{n1}\omega_{n2}}{\pi(\omega_{n2}^2 - \omega_{n1}^2)}, \ \beta_1 = \frac{\pi(\omega_{n2}\xi_2 - \omega_{n1}\xi_1)}{15(\omega_{n2}^2 - \omega_{n1}^2)}.$$
(7)

Here  $\omega_{n1}$  and  $\omega_{n2}$  are the first and second critical speeds (r/min),  $\xi_1$  and  $\xi_2$  the corresponding modal damping ratios, respectively.

The FE model of the rotor system with pedestal looseness is shown in Fig. 3. In the figure,  $k_{zl}$ ,  $k_{yl}$ ,  $k_{zr}$ ,  $k_{yr}$  and  $c_{zl}$ ,  $c_{yl}$ ,  $c_{zr}$ ,  $c_{yr}$  denote the stiffnesses and damping in z and y directions, respectively;  $m_{br}$  is the right pedestal mass;  $k_{br}$ ,  $c_{br}$  are the equivalent stiffness and equivalent damping of the pedestal, respectively.



# 3. Nonlinear characteristics of the rotor system with pedestal looseness under two loading conditions

Based on the API Standard 617 [17], two unbalance loading conditions are determined for the rotor system. The material parameters, bearing parameters and loads which are used to simulate the looseness fault are shown in Table 1. The modal damping ratios are assumed as  $\xi_1 = \xi_2 = 0.04$  and the geometric parameters of the rotor system can be found in Ref. [18]. Without regard to pedestal looseness, the unbalance responses of disc 2 (node 19) at vertical direction (y direction) under two loading conditions are shown in Fig. 4. It can be observed from the figure that the first order resonant response is predominant at 1680 r/min under the first loading condition and the second order resonant response at 6450 r/min under the second loading condition.

		,		
Material	Elastic modulus <i>E</i> (GPa)	Poisson's ratio v	Density $\rho$ (kg/m <sup>3</sup> )	
parameters	207	0.3	7850	
	Horizontal	Vertical	Horizontal	Vertical
Bearing	stiffness	stiffness	damping	damping
parameters	$k_{zl} = k_{zr} (MN/m)$	$k_{yl} = k_{yr} (MN/m)$	$c_{zl} = c_{zr} (kN \cdot s/m)$	$c_{yl} = c_{yr} (kN \cdot s/m)$
-	200	500	2	2
Dedestal	Base stiffness	Base damping cb	Pedestal mass	
parameters	$k_b (MN/m)$	(kN·s/m)	$m_{br}(\mathrm{kg})$	
	200	0.2	1.256	
	Unbalance	Dhasa angla of	Unbalance	Phase angle of
First unbalance	moment of disc 1	unbalance at dise 1 (°)	moment of disc 2	unbalance at disc
loading condition	$mr(g \cdot m)$	undalance at disc 1 ()	<i>mr</i> (g·m)	2 (°)
	0.156	0	0.156	0
Second unbalance loading condition	Unbalance	Dhasa angla of	Unbalance	Phase angle of
	moment of disc 1	unbalance at dise 1 (°)	moment of disc 2	unbalance at disc
	$mr(g \cdot m)$	unbanance at disc 1 ()	<i>mr</i> (g·m)	2 (°)
	0.156	0	0.156	180

Table 1. Parameters of the rotor, bearing and pedestal



Fig. 4. Unbalance response of the rotor system in y direction under two loading conditions

#### 3.1. Dynamic characteristics under the first loading condition

Under the first loading condition, the influences of the stiffness of non-loosened bolts  $(k_{b1})$ , looseness clearance  $(\delta_1)$  and rotating speed  $(\omega)$  on the nonlinear dynamic characteristics of the right journal are analyzed. Only vibration responses of the right journal at looseness position (node 26) are presented in this paper. Since the vibration of the rotor system is violent under the resonance conditions, the influences of  $k_{b1}$  and  $\delta_1$  are analyzed under the resonance speed of 1680 r/min and the influences of  $\omega$  are discussed finally.

#### 3.1.1. The influence of the stiffness of non-loosened bolts

Assuming the pedestal displacement is less than looseness clearance  $(y_p \le \delta_1)$  at  $\omega_{n1} = 1680$  r/min, namely  $k_{br}$  changes only when the pedestal contacts the base, and the damping of non-loosened bolts  $c_{b1}$  changes with different  $k_{b1}$ , as is shown in Table 2.

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<i>k</i> <sub>b1</sub> / (N/m)	$\lambda = \log(k_{b1} / 2)$	$c_{b1}$ / (N·s/m)
2×10 <sup>8</sup>	8	2000
2×107	7	1000
2×10 <sup>6</sup>	6	500
2×10 <sup>5</sup>	5	250
2×10 <sup>4</sup>	4	125
2×10 <sup>3</sup>	3	62.5
2×10 <sup>2</sup>	2	31.25
2×10 <sup>1</sup>	1	15.63
2	0	7.81

Table 2. Stiffness and damping coefficients of non-loosened bolts

The spectrum cascades and rotor orbits of the system are shown in Fig. 5 based on the parameters in Table 2. In the spectrum cascade, left-hand abscissa is the pedestal equivalent stiffness ( $\lambda = log(k_{b1}/2)$ ) and 1X denotes the rotating frequency. The rotor orbits under different  $\lambda$  are shown in the right side. From Fig. 5, it is clear that the amplitude of the right journal increases with the decrease of  $\lambda$ . The spectral components display only 1X at  $\lambda = 8$ , integer-order super-harmonics (nX, n = 1, 2, 3...) at  $\lambda = 7, 6, 5, 4$ ; 1/2 fractional harmonic components, such as 1X/2, 3X/2, etc., at  $\lambda = 3$ ; complicated lower frequency components at  $\lambda = 2$ ; and only 1X at  $\lambda = 1, 0$ . The rotor orbit shows a circle at  $\lambda = 8$ ; an ellipse with limited bottom at  $\lambda = 7, 6, 5, 4$ ; two ellipses at  $\lambda = 3$ , which shows the system motion is P2 (period-2); an upward spiral at  $\lambda = 2$ , 1 and a translational ellipse along pedestal looseness direction at  $\lambda = 0$  because the pedestal motion breaks away from the base. The system motion is also determined by Poincaré maps, as is shown in Fig. 6. From the figure, it can be seen that P2 motion appears at  $\lambda = 3$  and chaotic motion at  $\lambda = 2$ .







#### 3.1.2. The influence of the looseness clearance

Assuming  $\omega = \omega_{n1} = 1680$  r/min and  $\delta_1 = 1$  mm,  $k_{br}$  will change twice only when  $y_p > \delta_1$ . The loosened bolts under this condition, which is similar to double constraint problem. In this paper,  $y_p > \delta_1$  starts from  $\lambda = 3$ , and occurs at  $\lambda = 3$ , 2, 1, 0. The vibration responses of the right journal are shown in Fig. 7 and Poincaré maps at  $\lambda = 3$ , 2, 1 in Fig. 8. It can be seen from the spectrum cascade (see Fig. 7) that 1/3 fractional harmonic components, such as 1X/3, 2X/3, 4X/3, etc. and some continous spectra in the high frequency range appear at  $\lambda = 3$ ; the low combination frequency components besides high combination frequency components and contious spectra at  $\lambda = 2$ , 1, 0. The rotor orbit shows the limits of the bottom and the top. Poincaré maps (see Fig. 8) show that choatic motions appear at  $\lambda = 3$ , 2, 1. The appearances show that the contact and rebound times of the pedestal between the base and the loosened bolts increase with the decrease of  $\lambda$ .



**Fig. 7.** Vibration responses of rotor system at  $\lambda = 3, 2, 1, 0$  under condition 1



#### 3.1.3. The influence of rotating speeds

The stiffness and damping of non-loosened bolts are assumed as  $k_{b1} = 20$  kN/m and  $c_{b1} = 125$  N·s/m, respectively. Assuming  $y_p < \delta_1$ , namely  $k_{br}$  changes when the pedestal contacts the base. The rotating speed parameters are listed in Table 3. The vibration responses of the rotor system with pedestal looseness under different  $\omega$  are shown in Fig. 9 and Poincaré maps at  $\gamma = 3, 3.5, 4, 4.5, 5, 6$  in Fig. 10.

In the spectrum cascade, left-hand abscissa is the ratio of the rotating speed to the first critical speed ( $\gamma = \omega/\omega_{n1}$ ). From the figure it can be seen that the system response shows several integer-order super-harmonics (nX, n = 1, 2, 3...) at  $\gamma = 0.5, 1, 1.5, 2.1/\gamma$  fractional harmonic components can be observed at  $\gamma = 3, 4, 5, 6$ , which shows that the system motions

are period- $\gamma$ . Poincaré map at  $\gamma = 3.5$  shows that the chaotic motion appears; P10 at  $\gamma = 4.5$ , which is also proved by 1X/10 with small amplitude in spectrum cascade; P6 at  $\gamma = 5.5$ . From the rotor orbits, it is obvious that the right journal vibrates violently and its orbit is similar to a spiral spring along looseness direction at high rotating speeds.

	able et itotatin	S speed parameters	
Rotating speed $\omega$ / (r/min)	$\gamma = \omega / \omega_{n1}$	Rotating speed $\omega$ (r/min)	$\gamma = \omega / \omega_{n1}$
840	0.5	5880	3.5
1680	1.0	6720	4.0
2520	1.5	7560	4.5
3360	2.0	8400	5.0
4200	2.5	9240	5.5
5040	3.0	10080	6.0

Table 3. Rotating speed parameters



### (a) $\gamma = 0, 0.5, 1, 1.5, 2, 2.5, 3$

(b)  $\gamma = 3.5, 4, 4.5, 5, 5.5, 6$ 

Fig. 9. Vibration responses of the rotor system at different  $\gamma$  under condition 1



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#### 3.2. Dynamic characteristics under the second loading condition

Under the second loading condition, the parameters of the rotor, bearing and pedestal are the same as those under the first loading condition and the rotating speed  $\omega = \omega_{n2} = 6450$  r/min is selected to analyze the influences of  $k_{b1}$  and  $\delta_1$  on the vibration responses of the right journal.

#### 3.2.1. The influence of the stiffness of non-loosened bolts

The system vibration responses at different  $\lambda$  are shown in Fig. 11 and Poincaré maps at  $\lambda = 4$ , 3, 2 in Fig. 12. It can be seen from Figs. 11 and 12 that the system motion is P1 at  $\lambda = 8$ , 7, 6, 5 and rich multiple frequency components appear with the decrease of  $\lambda$ ; P4 appears at  $\lambda = 4$ , chaotic motion at  $\lambda = 3$ , 2 and P1 at  $\lambda = 1$ , 0. The rotor orbits are stable at  $\lambda = 8$ , 7, 6, 5, fluctuate along looseness direction at  $\lambda = 4$ , 3, 2, 1. An ellipse orbit at  $\lambda = 0$  indicates that the non-loosened bolts lose their binding force for the pedestal and the pedestal is suspended, which shows that there is no contact between the pedestal and the base by analyzing the position of the right pedestal in y direction.









#### 3.2.2. The influence of the looseness clearance

The vibration responses of the right journal under  $y_p > \delta_1$  and  $\delta_1 = 1$  mm are shown in Fig. 13 and Poincaré maps at  $\lambda = 3, 2, 1$  in Fig. 14. Multiple frequency components, such as 2*X*, 3*X*, etc., and some combination frequencies with small amplitude appears and rotor orbits are limited at the bottom and the top at  $\lambda = 4, 3, 2, 1$ , 0. The Poincaré maps are similar at  $\lambda = 3, 2, 1$ , which shows that the system motions are quasi-periodic. These features are different

than these under condition 1 (see Figs. 7 and 8).

#### 3.2.3. The influence of rotating speeds

The parameters are the same as those under the condition 1, and the vibration responses of the rotor system at different  $\gamma$  are shown in Fig. 15 and Poincaré maps at  $\gamma = 3, 3.5, 4, 4.5, 5, 6$  in Fig. 16. It can be seen from the Figs. 15 and 16 that the integer-order super-harmonics (nX, n = 1, 2, 3, ...) at  $\gamma = 0.5, 1, 1.5, 2$ . P8 appears at  $\gamma = 2.5$ ; P4 at  $\gamma = 3, 4$ ; chaotic motion at  $\gamma = 3.5, 6$ ; P10 at  $\gamma = 4.5$ ; P16 at  $\gamma = 5$ ; and P6 at  $\gamma = 5.5$ . The rotor orbits are similar to upward spirals at large  $\gamma$ , which are similar to those under condition 1.

The system motion forms under two loading conditions are listed in Table 4.



**Fig. 13.** Vibration responses of rotor system at  $\lambda = 4, 3, 2, 1, 0$  under condition 2



**Fig. 14.** Poincaré maps at  $\lambda = 3, 2, 1$  under condition 2



**Fig. 15.** Vibration responses of the rotor system at different  $\gamma$  under condition 2

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Fig. 16. Poincar	é maps at $\gamma =$	3, 3.5, 4, 4	4.5, 5, 6 un	der condition 2
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Domanatana of		Condition 1		Condition 2		
rarameters of	Malaaa	System	Frequency	Valesa	System	Frequency
Influence	values	motion	features	values	motion	features
	$\lambda = 8, 7, 6,$	D1	1X, 2X,	$\lambda = 8, 7,$	D1	1X, 2X,
$\lambda = \log(k_{b1} / 2)$ under $y_p < \delta_1$	5,4	PI	3 <i>X</i> , etc.	6, 5	PI	3 <i>X</i> , etc.
	$\lambda = 3$	P2	1X/2, 1X,	$\lambda = 4$	P4	1 <i>X</i> /4, 1 <i>X</i> /2,
			3X/2, etc.			3X/4, etc.
	1 0	Chaotic	Continuous	$\lambda = 3, 2$	Chaotic	Continuous
	$\lambda - \lambda$	motion	spectra		motion	spectra
	$\lambda = 1, 0$	P1	1X	$\lambda = 1, 0$	P1	1X
						1X, 2X, 3X and
$\lambda = \log(k_{b1}/2)$	$\lambda = 3, 2,$	Chaotic	Continuous	$\lambda = 4, 3,$	Quasi-	some combination
under $y_p > \delta_1$	1, 0	motion	spectra	2, 1, 0	periodic	frequency
<i></i>						components
	$\gamma = 0.5, 1,$	P1	1 <i>X</i> , 2 <i>X</i> ,	$\gamma = 0.5, 1,$	D1	1 <i>X</i> , 2 <i>X</i> ,
	$\gamma = 0.5, 1, 1.5, 2$	P1	1 <i>X</i> , 2 <i>X</i> , 3 <i>X</i> , etc.	$\gamma = 0.5, 1, 1.5, 2$	P1	1 <i>X</i> , 2 <i>X</i> , 3 <i>X</i> , etc.
	$\gamma = 0.5, 1,$ 1.5, 2 $\gamma = 3, 4,$	P1	$1X, 2X, 3X, etc. 1X/\gamma, 2X/\gamma,$	$\gamma = 0.5, 1,$ 1.5, 2	P1	1X, 2X, 3X, etc. 1X/8, 1X/4,
	$\begin{array}{c} \gamma = 0.5,  1, \\ 1.5,  2 \\ \gamma = 3,  4, \\ 5,  6 \end{array}$	P1 Period-γ	$1X, 2X, 3X, etc. 1X/\gamma, 2X/\gamma, etc.$	$\gamma = 0.5, 1,$ 1.5, 2 $\gamma = 2.5$	P1 P8	1X, 2X, 3X, etc. 1X/8, 1X/4, 3X/8, etc.
	$\begin{array}{l} \gamma = 0.5, 1, \\ 1.5, 2 \\ \gamma = 3, 4, \\ 5, 6 \end{array}$	P1 Period-γ Chaotic	$\begin{array}{c} 1X, 2X, \\ 3X, \text{etc.} \\ 1X/\gamma, 2X/\gamma, \\ \text{etc.} \\ \end{array}$ Continuous	$\gamma = 0.5, 1,$ 1.5, 2 $\gamma = 2.5$ $\gamma = 3$	P1 P8	1X, 2X, 3X, etc. 1X/8, 1X/4, 3X/8, etc. 1X/4, 1X/2,
	$ \begin{array}{l} \gamma = 0.5, 1, \\ 1.5, 2 \\ \gamma = 3, 4, \\ 5, 6 \\ \gamma = 3.5 \end{array} $	P1 Period-γ Chaotic motion	$\begin{array}{c} 1X, 2X, \\ 3X, \text{ etc.} \\ 1X/\gamma, 2X/\gamma, \\ \text{ etc.} \\ \hline \text{Continuous} \\ \text{ spectra} \end{array}$	$\gamma = 0.5, 1,$ 1.5, 2 $\gamma = 2.5$ $\gamma = 3$	P1 P8 P4	1X, 2X, 3X, etc. 1X/8, 1X/4, 3X/8, etc. 1X/4, 1X/2, 3X/4, etc.
y = (y)/(y)	$   \begin{array}{l} \gamma = 0.5, 1, \\ 1.5, 2 \\ \gamma = 3, 4, \\ 5, 6 \\ \gamma = 3.5 \\ \gamma = 4.5 \end{array} $	P1 Period-γ Chaotic motion P10	1 <i>X</i> , 2 <i>X</i> , 3 <i>X</i> , etc. 1 <i>X</i> /γ, 2 <i>X</i> /γ, etc. Continuous spectra 1 <i>X</i> /10, 1 <i>X</i> /5,	$\gamma = 0.5, 1,$ 1.5, 2 $\gamma = 2.5$ $\gamma = 3$ $\gamma = 35, 6$	P1 P8 P4 Chaotic	1X, 2X, 3X, etc. 1X/8, 1X/4, 3X/8, etc. 1X/4, 1X/2, 3X/4, etc. Continuous
$\gamma = \omega/\omega_{n1}$	$   \begin{array}{l} \gamma = 0.5, 1, \\ 1.5, 2 \\ \gamma = 3, 4, \\ 5, 6 \\ \gamma = 3.5 \\ \gamma = 4.5 \end{array} $	P1 Period-γ Chaotic motion P10	1 <i>X</i> , 2 <i>X</i> , 3 <i>X</i> , etc. 1 <i>X</i> /γ, 2 <i>X</i> /γ, etc. Continuous spectra 1 <i>X</i> /10, 1 <i>X</i> /5, 3 <i>X</i> /10, etc.	$\gamma = 0.5, 1,$ 1.5, 2 $\gamma = 2.5$ $\gamma = 3$ $\gamma = 3.5, 6$	P1 P8 P4 Chaotic motion	1X, 2X, 3X, etc. 1X/8, 1X/4, 3X/8, etc. 1X/4, 1X/2, 3X/4, etc. Continuous spectra
$\gamma = \omega/\omega_{n1}$	$\begin{array}{l} \gamma = 0.5, 1, \\ 1.5, 2 \\ \gamma = 3, 4, \\ 5, 6 \\ \gamma = 3.5 \\ \gamma = 4.5 \\ \gamma = 5.5 \end{array}$	P1 Period-γ Chaotic motion P10 P6	1 <i>X</i> , 2 <i>X</i> , 3 <i>X</i> , etc. 1 <i>X</i> /γ, 2 <i>X</i> /γ, etc. Continuous spectra 1 <i>X</i> /10, 1 <i>X</i> /5, 3 <i>X</i> /10, etc. 1 <i>X</i> /6, 1 <i>X</i> /3,	$\gamma = 0.5, 1,$ 1.5, 2 $\gamma = 2.5$ $\gamma = 3$ $\gamma = 3.5, 6$ $\gamma = 4.5$	P1 P8 P4 Chaotic motion	1X, 2X, 3X, etc. 1X/8, 1X/4, 3X/8, etc. 1X/4, 1X/2, 3X/4, etc. Continuous spectra 1X/10, 1X/5,
$\gamma = \omega/\omega_{n1}$	$   \begin{array}{l} \gamma = 0.5, 1, \\ 1.5, 2 \\ \gamma = 3, 4, \\ 5, 6 \\ \gamma = 3.5 \\ \gamma = 4.5 \\ \gamma = 5.5 \end{array} $	P1 Period-γ Chaotic motion P10 P6	1 <i>X</i> , 2 <i>X</i> , 3 <i>X</i> , etc. 1 <i>X</i> /γ, 2 <i>X</i> /γ, etc. Continuous spectra 1 <i>X</i> /10, 1 <i>X</i> /5, 3 <i>X</i> /10, etc. 1 <i>X</i> /6, 1 <i>X</i> /3, 1 <i>X</i> /2, etc.	$\gamma = 0.5, 1,$ 1.5, 2 $\gamma = 2.5$ $\gamma = 3$ $\gamma = 3.5, 6$ $\gamma = 4.5$	P1 P8 P4 Chaotic motion P10	1X, 2X, 3X, etc. 1X/8, 1X/4, 3X/8, etc. 1X/4, 1X/2, 3X/4, etc. Continuous spectra 1X/10, 1X/5, 3X/10, etc.
$\gamma = \omega/\omega_{n1}$	$             \gamma = 0.5, 1,                                   $	P1 Period-γ Chaotic motion P10 P6	1 <i>X</i> , 2 <i>X</i> , 3 <i>X</i> , etc. 1 <i>X</i> /γ, 2 <i>X</i> /γ, etc. Continuous spectra 1 <i>X</i> /10, 1 <i>X</i> /5, 3 <i>X</i> /10, etc. 1 <i>X</i> /6, 1 <i>X</i> /3, 1 <i>X</i> /2, etc.	$\gamma = 0.5, 1,$ 1.5, 2 $\gamma = 2.5$ $\gamma = 3$ $\gamma = 3.5, 6$ $\gamma = 4.5$ $\gamma = 5$	P1 P8 P4 Chaotic motion P10 P16	1X, 2X, 3X, etc. 1X/8, 1X/4, 3X/8, etc. 1X/4, 1X/2, 3X/4, etc. Continuous spectra 1X/10, 1X/5, 3X/10, etc. 1X/16, 1X/8,
$\gamma = \omega/\omega_{n1}$	$   \begin{array}{r} \gamma = 0.5, 1, \\ 1.5, 2 \\ \gamma = 3, 4, \\ 5, 6 \\ \gamma = 3.5 \\ \gamma = 4.5 \\ \gamma = 5.5 \end{array} $	P1 Period-γ Chaotic motion P10 P6	1 <i>X</i> , 2 <i>X</i> , 3 <i>X</i> , etc. 1 <i>X</i> /γ, 2 <i>X</i> /γ, etc. Continuous spectra 1 <i>X</i> /10, 1 <i>X</i> /5, 3 <i>X</i> /10, etc. 1 <i>X</i> /6, 1 <i>X</i> /3, 1 <i>X</i> /2, etc.	$\gamma = 0.5, 1,$ 1.5, 2 $\gamma = 2.5$ $\gamma = 3$ $\gamma = 3.5, 6$ $\gamma = 4.5$ $\gamma = 5$	P1 P8 P4 Chaotic motion P10 P16	1X, 2X, 3X, etc. 1X/8, 1X/4, 3X/8, etc. 1X/4, 1X/2, 3X/4, etc. Continuous spectra 1X/10, 1X/5, 3X/10, etc. 1X/16, 1X/8, 3X/16, etc.
$\gamma = \omega/\omega_{n1}$	$   \begin{array}{r} \gamma = 0.5, 1, \\ 1.5, 2 \\ \gamma = 3, 4, \\ 5, 6 \\ \gamma = 3.5 \\ \gamma = 4.5 \\ \gamma = 5.5 \end{array} $	P1 Period-γ Chaotic motion P10 P6	1 <i>X</i> , 2 <i>X</i> , 3 <i>X</i> , etc. 1 <i>X</i> /γ, 2 <i>X</i> /γ, etc. Continuous spectra 1 <i>X</i> /10, 1 <i>X</i> /5, 3 <i>X</i> /10, etc. 1 <i>X</i> /6, 1 <i>X</i> /3, 1 <i>X</i> /2, etc.	$\gamma = 0.5, 1,$ 1.5, 2 $\gamma = 2.5$ $\gamma = 3$ $\gamma = 3.5, 6$ $\gamma = 4.5$ $\gamma = 5$ $\gamma = 5$	P1 P8 P4 Chaotic motion P10 P16 P6	$\begin{array}{c} 1X, 2X, \\ 3X, \text{ etc.} \\ 1X/8, 1X/4, \\ 3X/8, \text{ etc.} \\ 1X/4, 1X/2, \\ 3X/4, \text{ etc.} \\ \hline \\ \text{Continuous} \\ \text{spectra} \\ 1X/10, 1X/5, \\ 3X/10, \text{ etc.} \\ 1X/16, 1X/8, \\ 3X/16, \text{ etc.} \\ 1X/6, 1X/3, \\ \hline \end{array}$

Table 4. Comparison of system motion forms under two loading conditions

#### 4. Conclusions

In this paper, nonlinear dynamic characteristics of a rotor system with pedestal looseness fault are simulated under two loading conditions by considering the effects of the stiffness of non-loosened bolts, looseness clearance and rotating speeds. The following conclusions can be drawn from the present study:

(1) When the pedestal displacement is less than the looseness clearance, with the decrease of the stiffness of non-loosened bolts, the system motion under the first loading condition changes from P1 through P2, chaotic motion to P1. However, the system motion under the second loading condition changes from P1 through P4, chotic motion to P1.

(2) When the pedestal displacement is greater than the looseness clearance, with the decreasese of the stiffness of non-loosened bolts, chaotic motion appears under the first loading condition, but quasi-periodic motion under the second loading condition.

(3) When the pedestal displacement is less than the looseness clearance, the system motion is P1 under the rotating speed less than or equal to two times of the first critical speed. Periodic motion and chaotic motion appear alternately when the rotating speed is greater than two times of the first critical speed. And periodic motion forms under condition 2 are more complicated than those under condition 1.

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#### References

- [1] Wen B. C., Wu X. H., Ding Q., et al. Theory and Experiment of Nonlinear Dynamics for Rotating Machinery with Faults. Science, 2004.
- [2] Muszynska A., Goldman P. Chaotic responses of unbalanced rotor bearing stator systems with looseness or rubs. Chaos, Solitons and Fractals, Vol. 5, Issue 9, 2009, p. 1683-1704.
- [3] Muszynska A. Rotordynamics. Boca Raton, CRC, 2005.
- [4] Chu F., Tang Y. Stability and non-linear responses of a rotor-bearing system with pedestal looseness. Journal of Sound and Vibration, Vol. 241, Issue 5, 2001, p. 879-893.
- [5] Lu W., Chu F. Experimental investigation of pedestal looseness in a rotor-bearing system. Key Engineering Materials, Vols. 413-414, 2009, p. 599-605.
- [6] Duan J. A., Huang Z. K. Nonlinear model of rotor system with loose fault. Journal of Central South University of Technology, Vol. 33, Issue 1, 2002, p. 78-81.
- [7] Chen G. Nonlinear dynamics of unbalance-looseness coupling faults of rotor-ball bearing-stator coupling system. Chinese Journal of Mechanical Engineering, Vol. 44, Issue 3, 2008, p. 82-88.
- [8] Kim Y. B., Choi S. K. A multiple harmonic balance method for the internal resonant vibration of a non-linear Jeffcott rotor. Journal of Sound and Vibration, Vol. 208, Issue 5, 1997, p. 745-761.
- [9] Hossain M. Z., Mizutani K., Sawai H., et al. Preloading effects on clearance problem in rotorcoupling vibration system: experimentation and simulation. Chaos, Solitons and Fractals, Vol. 14, Issue 9, 2002, p. 1371-1378.
- [10] Gonsalves D. H., Neilson R. D., Barr A. D. S. A study of the response of a discontinuously nonlinear rotor system. Nonlinear Dynamics, Vol. 7, Issue 4, 1995, p. 451-470.
- [11] Karpenko E. V., Wiercigroch M., Cartmell M. P. Regular and chaotic dynamics of a discontinuously nonlinear rotor system. Chaos, Solitons and Fractals, Vol. 13, Issue 6, 2002, p. 1231-1242.
- [12] Karlberg M., Aidanpaa J. O. Investigation of an unbalanced rotor system with bearing clearance and stabilizing rods. Chaos, Solitons and Fractals, Vol. 20, Issue 2, 2004, p. 363-374.

- [13] Gao S. H., Long X. H., Meng G. Nonlinear response and nonsmooth bifurcations of an unbalanced machine-tool spindle-bearing system. Nonlinear Dynamics, Vol. 54, Issue 4, 2008, p. 365-377.
- [14] Ren Z. H., Chen H., Ma H., et al. Pedestal looseness fault analysis of a vertical dual-disk over-hung rotor-bearing system. Transactions of the Chinese Society for Agricultural Machinery, Vol. 38, Issue 11, 2007, p. 196-201.
- [15] Ma H., Zhao X. Y., Teng Y. N., et al. Analysis of dynamic characteristics for a rotor system with pedestal looseness. Shock and Vibration, Vol. 18, Issues 1-2, 2011, p. 13-27.
- [16] Behzad M., Asayeshthe M. Numerical and experimental investigation on vibration of rotors with loose discs. Proceedings of the Institution of Mechanical Engineers, Part C: Journal of Mechanical Engineering Science, Vol. 224, Issue 1, 2010, p. 85-94.
- [17] API 617 Axial and Centrifugal Compressors and Turboexpanders for Petroleum. Chemical and Gas Industry Services, American Petroleum Institute, Washington D. C., 2002.
- [18] Ma H., Tai X. Y., Sun J., et al. Analysis of dynamic characteristics for a dual-disk rotor system with single rub-impact. Advanced Science Letters, Vol. 4, Issues 8-10, 2011, p. 2782-2789.