1891. Several rows of piles as barriers to isolate shear waves in saturated soils

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Abstract. There are three types of waves propagating in the saturated soils, which are fast pressure waves (P_1 waves for short), slow pressure waves (P_2 waves for short) and shear waves (S waves for short), and their wave numbers are firstly given in this paper. The lengths of elastic piles are much larger than the diameters in order to obtain ideal isolation effectiveness, so the isolation of SV waves by discontinuous barriers composed of several rows of piles can be simplified as a 2-dimensional plane problem. The conformal mapping method of complex functions and expansion method of wave functions are adopted, the incident, scattered and refracted waves are all expanded to the special functions with infinite series, the stresses and displacements are considered as continuous at the boundaries between the piles and adjacent saturated soils, and finally the theoretical solutions about the complex coefficients of the expanded potential functions are obtained. The normalized vertical displacement amplitudes are introduced, which are the ratios of displacement amplitudes of saturated soils caused by both the incident and scattered waves to those only by the incident SV waves, the isolation effectiveness of the barriers composed of one, two and three rows of piles are compared, and the factors that influence the isolation effectiveness are analyzed, which provides some references on the vibration isolation design.

Keywords: saturated soils, several rows of piles, shear waves, isolation effectiveness, vibration isolation design.

1. Introduction

Man-made vibrations such as pile driving construction, railway and highway traffic vehicles, heavy industry and engineering explosion are becoming more and more frequent, and they have caused adverse influences on precise instrument processing and human's life. In general, the previous approaches used to reduce the ground vibration strength near the man-made vibration sources can be categorized as the active and passive approaches. The active isolation measures are commonly used to isolate the vibration source from the ground; therefore, they are usually arranged around the vibration source in a very close distance. On the other hand, the passive isolation measures are normally used to reduce the energy transmitted to the protected structures, and consequently, they are usually arranged surrounding protected structures. Continuous and discontinuous barriers are the most two types of passive measures: open or in-filled trenches, and one row or several rows of piles or cylindrical cavities. Among these, discontinuous barriers are the advantageous barriers for the fields with higher groundwater level.

Continuous barriers such as open and in-filled trenches have been studied in the recent thirty years [1-8], and the relative results all reveal that ideal vibration isolation effectiveness can be obtained from a deeper open trench while the trench width has less influence on the vibration isolation effectiveness. But as to man-made vibrations with low frequencies, an open trench should be constructed too deep to obtain better vibration isolation effectiveness, and so to maintain the stability of the four sides of the trench is much important for guaranteeing the isolation effectiveness in the future. In order to reduce the maintenance fees of open trenches and obtain the appropriate vibration isolation effectiveness, discontinuous barriers composed of a row of piles and cylindrical cavities are designed to replace the open trenches where the soils are not suitable for excavations and maintenance such as soft soils with higher water tables.

Nowadays some geotechnical engineers are interested in discontinuous barriers especially

those composed of a row of solid or pipe piles [9-18], and some important conclusions have been drawn out: the shear modulus of the piles, the spacing distances between the piles and the wall thickness of the pipe piles are the main factors that influence the isolation effectiveness.

From the above-mentioned studies we can find that, open trenches have better vibration isolation effectiveness but their applications are limited to the excavation depth, while one row or several rows of piles have less isolation effectiveness than open trenches but they are not limited to the depth and can be applied in all soils.

The soft soils with higher water table can be studied as saturated soils, and there are three types of waves propagating in the saturated soils, which are fast pressure waves (P_1 waves for short), slow pressure waves (P_2 waves for short) and shear waves (S waves for short) [19], while there are only two types of waves in elastic media, which are pressure waves (P waves for short) and shear waves (S waves for short). In this paper, the SV waves that cause the vertical vibration of the ground are taken as example, the isolation of SV waves by several rows of piles in saturated soils is studied, and the factors that influence the isolation effectiveness are analyzed.

2. Wave functions of the saturated soils

According to Biot's wave theory [19], the control equations in saturated soils are mainly include 4 equations as follows:

1) Stress-strain relationships:

$$\sigma_{ij} = \lambda u_{i,i} \delta_{ij} + 2\mu \varepsilon_{ij} - \alpha p_f \delta_{ij}.$$
 (1)

2) Seepage equation of continuity:

$$-p_f = Mw_{i,i} + \alpha Mu_{i,i}.$$
(2)

3) Motion equation regardless the body forces:

$$\sigma_{ii,i} = \rho \ddot{u}_i + \rho_f \ddot{w}_i. \tag{3}$$

4) Fluid motion equation:

$$-p_{f,i} = \rho_f \ddot{u}_i + m \ddot{w}_i + b \dot{w}_i, \tag{4}$$

where, $\rho = f \rho_f + (1 - f) \rho_s$, $m = \rho_f / f$, $b = \eta / k_d$, p_f is the pore water pressure, λ and μ are the Lamé elastic constants of the solid skeleton, u is the displacement of solid skeleton, w is the displacement of water compared to the solid skeleton, M and α are two compaction coefficients that reveal the soil particles and pore water, δ_{ij} is Kronecker Delta, ρ is the density of the whole saturated soils, ρ_s is the density of solid particles, ρ_f is the density of pore water, f is the porosity of the saturated soils, and η and k_d are viscous and permeability coefficients of the pore water.

Substituting Eqs. (1) and (2) into Eq. (3), and Eq. (2) to Eq. (4), the vector equations about the saturated soils are obtained:

$$\mu \nabla^2 \mathbf{u} + (\lambda_c + \mu) \nabla u_{i,i} + \alpha M \nabla w_{i,i} = \rho \ddot{\mathbf{u}} + \rho_f \ddot{\mathbf{w}},$$
(5a)
$$M (\alpha \nabla u_{i,i} + w_{i,i}) = \rho_f \ddot{\mathbf{u}} + m \ddot{\mathbf{w}} + b \dot{\mathbf{w}},$$
(5b)

where, $\lambda_c = \lambda + \alpha^2 M$, and **u** and **w** are the two displacement vectors.

Scalar potential functions of φ and ϕ and vector potential functions of ψ and χ are introduced, and the displacement vectors of **u** and **w** are parted as follows:

$$\mathbf{u} = \nabla \varphi + \nabla \times \psi, \tag{6a}$$

(6b)

 $\mathbf{w} = \nabla \phi + \nabla \times \chi.$

Put Eq. (6) into Eq. (5), the wave functions of saturated soils are obtained in the forms of matrix:

$$\begin{bmatrix} \lambda_c + 2\mu & \alpha M \\ \alpha M & M \end{bmatrix} \begin{pmatrix} \nabla^2 \varphi \\ \nabla^2 \phi \end{pmatrix} + \begin{bmatrix} \rho \omega^2 & \rho_f \omega^2 \\ \rho_f \omega^2 & Q \end{bmatrix} \begin{pmatrix} \varphi \\ \phi \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix},$$

$$\begin{bmatrix} \mu & 0 \\ \nabla^2 \chi \end{pmatrix} \begin{pmatrix} \nabla^2 \psi \\ \nabla^2 \chi \end{pmatrix} + \begin{bmatrix} \rho \omega^2 & \rho_f \omega^2 \\ \rho_f \omega^2 & Q \end{bmatrix} \begin{pmatrix} \psi \\ \chi \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix},$$

$$(7a)$$

where, $Q = m\omega^2 + i\omega b$ and ω is the frequency of incident waves.

The complex wave numbers k_1 , k_2 and k_s of P₁, P₂ and S waves can be obtained from Eq. (7):

$$k_{1,2}^{2} = \frac{B \mp \sqrt{B^{2} - 4AC}}{2A},$$
(8a)
$$k_{c}^{2} = \frac{C}{2A},$$
(8b)

$$\mu Q'$$

where, $A = (\lambda + 2\mu)M$, $B = (\lambda_c + 2\mu)Q + \rho\omega^2 M - 2\rho_f \omega^2 \alpha M$ and $C = -\rho_f^2 \omega^4 + \rho\omega^2 Q$.

In the Eq. (6), potential functions of φ_1 , φ_2 and ψ are used to denote the potential functions of P₁, P₂ and S waves propagated in the solid skeletons of the saturated soils, while ϕ_1 , ϕ_2 and χ to denote those propagated in the liquids, and these potential functions have the relationships as follows:

$$\phi_1 = \gamma_1 \varphi_1, \quad \phi_2 = \gamma_2 \varphi_2, \quad \chi = \gamma_s \psi, \tag{9}$$

where, γ_1 , γ_2 and γ_s are:

$$\gamma_{1,2} = \frac{-(\lambda_c + 2\mu)k_{1,2}^2 + \rho\omega^2}{\alpha M k_{1,2}^2 - \rho_f \omega^2},$$
(10a)

$$\gamma_s = -\frac{\rho_f \omega^2}{Q}.$$
 (10b)

3. Potential functions of waves

3.1. Potential functions of the incident SV waves

The waves generated by traffic loads and other man-made vibrations are mainly propagated in the above 10 meters under the ground surface, which is less than the Rayleigh wave length, while the length of the piles that used to isolate the vibrations are usually larger than 10 meters, so the waves generated by the vibrations can be scattered by the piles in large scale and then reduce the vibrations, and the scattering problem can be simplified as a 2-dimensional plane problem.

 P_2 waves would be attenuated quickly in the saturated soils and they accounted for the negligible proportion of the total energy of elastic waves, and so incident P_2 waves are ignored as for the passive vibration isolation for the protecting structures that are usually far away from the vibration sources. The incident P_1 waves or SV waves would be coupling scattered, that is to say that the scattering waves all include P_1 waves, P_2 waves and SV waves, the studying method and procedures are similar for the incident P_1 waves or SV waves, and so the incident SV waves are only considered.

The saturated soils are taken as infinite homogeneous media, and there buried N piles whose lengths are much larger than their diameters, so the isolation problem can be simplified as a

2-dimensional plane scattering problem. N local Cartesian coordinate systems of (x_j, y_j) $(1 \le j \le N)$ are introduced, as shown in Fig. 1. The point M is an arbitrary point in the saturated soils.



Fig. 1. Distribution of piles and coordinate systems

The potential function of incident SV waves with frequency of ω in the reference Cartesian coordinate systems of (x, y) can be written as:

$$\psi^{inc} = \psi_0 \exp[ik_s(x\cos\beta + y\sin\beta)],\tag{11}$$

where, superscript of "*inc*" denotes incident, ψ_0 is the displacement amplitude of incident SV waves, β is the angle between incident SV waves and horizontal direction, and k_s is wave number of SV waves in the saturated soils that can be calculated from Eq. (8b). In order to discuss and compute easily, the public time factor of $e^{-i\omega t}$ is omitted in Eq. (11) and the lower equations.

According to the relationships of z = x + iy and $e^{i\beta} = \cos\beta + i\sin\beta$, Eq. (11) can be represented in the form of reference complex coordinate system (z, \bar{z}) :

$$\psi^{inc} = \psi_0 \exp\left[\frac{ik_s}{2} \left(ze^{-i\beta} + \bar{z}e^{i\beta}\right)\right]. \tag{12}$$

The complex coordinate system (z, \overline{z}) and *j*th local complex coordinate system (z_j, \overline{z}_j) have the relationship of $z = z_j + d_j$, so Eq. (1) can be further written as:

$$\psi_j^{inc} = \psi_0 \exp\left\{\frac{ik_s}{2} \left[(z_j + d_j)e^{-i\beta} + (\bar{z}_j + \bar{d}_j)e^{i\beta} \right] \right\}.$$
(13)

3.2. Potential functions of the scattered waves

Similar to Eq. (13), the potential function of scattering P₁ waves by *k*th pile that satisfies with the steady-state wave motion and Sommerfeld radiant condition in *k*th local complex coordinate system (z_k, \bar{z}_k) can be written as:

$$\varphi_{1k}^{sc}(z,\bar{z}) = \sum_{n=-\infty}^{\infty} A_n^k H_n^{(1)}(k_1|z_k|) \left(\frac{z_k}{|z_k|}\right)^n,\tag{14}$$

where, superscript of "*sc*" denotes scattering, A_n^k ($1 \le k \le N$) are the complex coefficients to be determined from the boundary conditions, and $H_n^{(1)}(\cdot)$ is the first kind of Hankel function with order *n*.

The location of observation point *M* can be written in *j*th and *k*th local complex coordinate systems as follows: $z = z_j + d_j = z_k + d_k$, $z_j = d_{jk} + z_k$, $z - d_k = z_j - d_{jk}$, so Eq. (14) can be written in reference coordinate system (z, \bar{z}) and *j*th local complex coordinate system (z_i, \bar{z}_i) :

$$\varphi_{1jk}^{sc} = \sum_{n=-\infty}^{\infty} A_n^k H_n^{(1)} (k_1 | z_j - d_{jk} |) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^n.$$
(15)

The total P_1 waves at the observation point M that scattered by all the N piles are:

$$\varphi_1^{sc} = \sum_{n=-\infty}^{\infty} \sum_{k=1}^{N} A_n^k H_n^{(1)} \left(k_1 | z_j - d_{jk} | \right) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^n.$$
(16a)

And the total scattered P₂ waves and SV waves are given directly based on Eq. (16a):

$$\varphi_2^{sc} = \sum_{n=-\infty}^{\infty} \sum_{k=1}^{N} B_n^k H_n^{(1)} (k_2 | z_j - d_{jk} |) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^n,$$
(16b)

$$\psi^{sc} = \sum_{n=-\infty}^{\infty} \sum_{k=1}^{N} C_n^k H_n^{(1)} (k_s | z_j - d_{jk} |) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^n,$$
(16c)

where, B_n^k and C_n^k ($1 \le k \le N$) are the complex coefficients to be determined from the boundary conditions.

3.3. Potential functions of the refracted waves in the piles

The piles are taken as infinite homogeneous media, and there are two kinds of waves propagate in them, which are P and S waves. Similar to the piles in elastic soils [20], the potential functions of P and SV waves in the *j*th pile are:

$$\varphi_j^{re} = \sum_{\substack{n=-\infty\\\infty}}^{\infty} D_n^j J_n(\tilde{k}_{pj}|z_j|) \left(\frac{z_j}{|z_j|}\right)^n, \tag{17a}$$

$$\psi_j^{re} = \sum_{n=-\infty} E_n^j J_n(\tilde{k}_{sj}|z_j|) \left(\frac{z_j}{|z_j|}\right)^n$$
(17b)

where, superscript of re denotes refracting, $D_n^j - E_n^j$ $(1 \le j \le N)$ are the complex coefficients to be determined from the boundary conditions, $J_n(\cdot)$ is the first kind of Bessel function with number of n, \tilde{k}_{pj} and \tilde{k}_{sj} are the wave numbers of P and SV waves respectively of the *j*th pile, and the superscript of \sim is adopted to denote the parameters and variables about the piles that can be easily distinguished from those of the saturated soils.

4. Theoretical solutions about the complex coefficients

Conformal mapping equation is introduced: $z_j = r_j \exp(i\theta_j)$, and stresses, displacements and water pressure in the saturated soils can be calculated from the above-mentioned potential functions:

$$u_r = e^{-i\theta} \frac{\partial(\varphi_1 + \varphi_2 - i\psi)}{\partial \bar{z}} + e^{i\theta} \frac{\partial(\varphi_1 + \varphi_2 + i\psi)}{\partial z},$$
(18a)

1891. Several rows of piles as barriers to isolate shear waves in saturated soils. Ping Xu, Ying Tie

$$u_{\theta} = ie^{i\theta} \frac{\partial(\varphi_1 + \varphi_2 + i\psi)}{\partial z} - e^{-i\theta} \frac{\partial(\varphi_1 + \varphi_2 - i\psi)}{\partial \bar{z}},$$
(18b)

$$w_r = e^{-i\theta} \frac{\partial(\gamma_1 \varphi_1 + \gamma_2 \varphi_2 - i\gamma_s \psi)}{\partial \bar{z}} + e^{i\theta} \frac{\partial(\gamma_1 \varphi_1 + \gamma_2 \varphi_2 + i\gamma_s \psi)}{\partial z},$$
(18c)

$$\sigma_r = -(\lambda_c + \gamma_1 \alpha M + \mu)k_1^2 \varphi_1 - (\lambda_c + \gamma_2 \alpha M + \mu)k_2^2 \varphi_2 + 2\mu \left[e^{2i\theta} \frac{\partial^2(\varphi_1 + \varphi_2 + i\psi)}{\partial z^2} + e^{-2i\theta} \frac{\partial^2(\varphi_1 + \varphi_2 - i\psi)}{\partial \bar{z}^2} \right],$$
(18d)

$$\tau_{r\theta} = 2i\mu \left[e^{2i\theta} \frac{\partial^2(\varphi_1 + \varphi_2 + i\psi)}{\partial z^2} - e^{-2i\theta} \frac{\partial^2(\varphi_1 + \varphi_2 - i\psi)}{\partial \bar{z}^2} \right], \tag{18e}$$

$$p_f = (\alpha + \gamma_1) M k_1^2 \varphi_1 + (\alpha + \gamma_2) M k_1^2 \varphi_2.$$
(18f)

The stresses and displacements at the boundaries between the piles and adjacent saturated soils are considered as continuous and the piles are water-proof, and the boundary conditions are depicted as:

$$\sigma_r^{inc+sc} = \tilde{\sigma}_r^{re}, \quad \tau_{r\theta}^{inc+sc} = \tilde{\tau}_{r\theta}^{re}, \quad u_r^{inc+sc} = \tilde{u}_r^{re}, \quad u_{\theta}^{inc+sc} = \tilde{u}_{\theta}^{re}, \quad w_r^{inc+sc} = 0.$$
(19)

Put the potential functions of Eqs. (13), (16) and (17) into Eqs. (18) and (19), and the infinite linear equation set about the complex coefficients of A_n^k - E_n^k are obtained:

$$\sum_{i=1}^{5} \sum_{n=-\infty}^{\infty} \sum_{k=1}^{N} \sum_{l=1}^{5} P_{il}^{kn} X_{ln}^{k} = Y_{i},$$
(20)

where, $X_{1n}^k = A_n^k$, $X_{2n}^k = B_n^k$, $X_{3n}^k = C_n^k$, $X_{4n}^k = D_n^k$ and $X_{5n}^k = E_n^k$. And the expressions about the matrix elements of P_{11}^{kn} - P_{55}^{kn} and Y_1 - Y_5 are:

$$\begin{split} P_{11}^{nk} &= -\lambda_{1}^{*}k_{1}^{2}H_{n}^{(1)}(k_{1}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n} + \frac{k_{1}^{2}e^{2i\theta_{j}}}{2}H_{n-2}^{(1)}(k_{1}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n-2} \\ &+ \frac{k_{1}^{2}e^{-2i\theta_{j}}}{2}H_{n+2}^{(1)}(k_{1}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n} + \frac{k_{2}^{2}e^{2i\theta_{j}}}{2}H_{n-2}^{(1)}(k_{2}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n-2} \\ &+ \frac{k_{2}^{2}e^{-2i\theta_{j}}}{2}H_{n+2}^{(1)}(k_{2}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n+2}, \\ P_{13}^{nk} &= -\lambda_{2}^{*}k_{2}^{2}H_{n-2}^{(1)}(k_{2}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n} + \frac{k_{2}^{2}e^{2i\theta_{j}}}{2}H_{n-2}^{(1)}(k_{2}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n-2} \\ &+ \frac{k_{2}^{2}e^{-2i\theta_{j}}}{2}H_{n+2}^{(1)}(k_{2}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n-2} \\ &- \frac{ik_{2}^{2}e^{-2i\theta_{j}}}{2}H_{n+2}^{(1)}(k_{s}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n+2}, \\ P_{13}^{nk} &= \delta_{jk}(\tilde{\lambda}_{j}^{*} + \tilde{\mu}_{j}^{*})\tilde{k}_{pj}^{2}J_{n}(\tilde{k}_{pj}|z_{j}|) \left(\frac{z_{j}}{|z_{j}|}\right)^{n} - \delta_{jk}\frac{\tilde{\mu}_{j}^{*}\tilde{k}_{pj}^{2}e^{2i\theta_{j}}}{2}J_{n-2}(\tilde{k}_{pj}|z_{j}|) \left(\frac{z_{j}}{|z_{j}|}\right)^{n-2} \\ &- \delta_{jk}\frac{\tilde{\mu}_{j}^{*}\tilde{k}_{2j}^{2}e^{2i\theta_{j}}}{2}J_{n+2}(\tilde{k}_{pj}|z_{j}|) \left(\frac{z_{j}}{|z_{j}|}\right)^{n-2} + \delta_{jk}\frac{i\tilde{\mu}_{j}^{*}\tilde{k}_{2j}^{2}e^{-2i\theta_{j}}}{2}J_{n+2}(\tilde{k}_{sj}|z_{j}|) \left(\frac{z_{j}}{|z_{j}|}\right)^{n+2}, \end{split}$$

$$\begin{split} &Y_{1} = -k_{s}^{2} \sin 2(\theta_{j} - \beta) \psi_{0} e^{ik_{s}[(z_{j} + d_{j})e^{-i\theta_{j}(z_{j} + d_{j})e^{i\theta_{j}/2}}, \\ &P_{21}^{nk} = i \frac{k_{1}^{2} e^{2i\theta_{j}}}{2} H_{n+2}^{(1)}(k_{1}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n-2} \\ &- i \frac{k_{1}^{2} e^{-2i\theta_{j}}}{2} H_{n+2}^{(1)}(k_{1}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n+2}, \\ &P_{22}^{nk} = i \frac{k_{2}^{2} e^{2i\theta_{j}}}{2} H_{n+2}^{(1)}(k_{2}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n-2} \\ &- i \frac{k_{2}^{2} e^{-2i\theta_{j}}}{2} H_{n+2}^{(1)}(k_{2}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n-2} \\ &- \frac{k_{2}^{2} e^{-2i\theta_{j}}}{2} H_{n+2}^{(1)}(k_{s}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n-2} \\ &- \frac{k_{2}^{2} e^{-2i\theta_{j}}}{2} H_{n+2}^{(1)}(k_{s}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n-2} \\ &- \frac{k_{2}^{2} e^{-2i\theta_{j}}}{2} H_{n+2}^{(1)}(k_{s}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n-2} \\ &- \frac{k_{2}^{2} e^{-2i\theta_{j}}}{2} H_{n+2}^{(1)}(k_{s}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n-2} \\ &+ \delta_{jk} \frac{t_{j}^{2} k_{j}^{2} e^{-2i\theta_{j}}}{2} J_{n+2}(k_{s}|z_{j}|z_{j}|) \left(\frac{z_{j}}{|z_{j}|}\right)^{n-2} \\ &+ \delta_{jk} \frac{t_{j}^{2} k_{j}^{2} e^{-2i\theta_{j}}}{2} J_{n+2}(k_{s}|z_{j}|z_{j}|) \left(\frac{z_{j}}{|z_{j}|}\right)^{n+2} \\ &+ k_{2} e^{-2i\theta_{j}}} H_{n+1}^{(1)}(k_{1}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}{|z_{j} - d_{jk}|}\right)^{n-1} \\ &- \frac{k_{1} e^{-2i\theta_{j}}}{2} H_{n+1}^{(1)}(k_{1}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}}{|z_{j} - d_{jk}|}\right)^{n+1} \\ &- \frac{k_{1} e^{-2i\theta_{j}}}{2} H_{n+1}^{(1)}(k_{s}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}}{|z_{j} - d_{jk}|}\right)^{n-1} \\ &- \frac{k_{2} e^{-i\theta_{j}}}}{2} H_{n+1}^{(1)}(k_{s}|z_{j} - d_{jk}|) \left(\frac{z_{j} - d_{jk}}}{|z_{j} - d$$

$$\begin{split} P_{42}^{nk} &= i \frac{k_2 e^{i \theta_j}}{2} H_{n+1}^{(1)} (k_2 | z_j - d_{jk} |) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^{n-1} \\ &+ i \frac{k_2 e^{-i \theta_j}}{2} H_{n+1}^{(1)} (k_2 | z_j - d_{jk} |) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^{n+1}, \\ P_{43}^{nk} &= -\frac{k_s e^{i \theta_j}}{2} H_{n-1}^{(1)} (k_s | z_j - d_{jk} |) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^{n-1} \\ &+ \frac{k_s e^{-i \theta_j}}{2} H_{n+1}^{(1)} (k_s | z_j - d_{jk} |) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^{n+1}, \\ P_{44}^{nk} &= -i \frac{\tilde{k}_{pj} e^{i \theta_j}}{2} J_{n-1} (\tilde{k}_{pj} | z_j |) \left(\frac{z_j}{|z_j|} \right)^{n-1} - i \frac{\tilde{k}_{pj} e^{-i \theta_j}}{2} J_{n+1} (\tilde{k}_{pj} | z_j |) \left(\frac{z_j}{|z_j|} \right)^{n+1}, \\ P_{45}^{nk} &= \delta_{jk} \frac{\tilde{k}_{sj} e^{i \theta_j}}{2} J_{n-1} (\tilde{k}_{sj} | z_j |) \left(\frac{z_j}{|z_j|} \right)^{n-1} - \delta_{jk} \frac{\tilde{k}_{sj} e^{-i \theta_j}}{2} J_{n+1} (\tilde{k}_{sj} | z_j |) \left(\frac{z_j}{|z_j|} \right)^{n+1}, \\ Y_4 &= i k_s \cos(\theta - \beta) \psi_0 e^{i k_s [(z_j + d_j) e^{-i \beta_j + (\tilde{x}_j + \tilde{d}_j) e^{i \beta}]/2}, \\ P_{51}^{nk} &= \frac{\gamma_1 k_1 e^{i \theta_j}}{2} H_{n-1}^{(1)} (k_1 | z_j - d_{jk}|) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^{n-1} \\ - \frac{\gamma_1 k_1 e^{-i \theta_j}}{2} H_{n+1}^{(1)} (k_2 | z_j - d_{jk}|) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^{n-1}, \\ P_{52}^{nk} &= \frac{\gamma_2 k_2 e^{i \theta_j}}{2} H_{n+1}^{(1)} (k_2 | z_j - d_{jk}|) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^{n-1} \\ - \frac{\gamma_2 k_2 e^{-i \theta_j}}{2} H_{n+1}^{(1)} (k_2 | z_j - d_{jk}|) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^{n-1}, \\ P_{53}^{nk} &= i \frac{\gamma_s k_s e^{i \theta_j}}{2} H_{n+1}^{(1)} (k_s | z_j - d_{jk}|) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^{n-1} \\ + i \frac{\gamma_s k_s e^{-i \theta_j}}{2} H_{n+1}^{(1)} (k_s | z_j - d_{jk}|) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^{n-1} \\ + i \frac{\gamma_s k_s e^{-i \theta_j}}{2} H_{n+1}^{(1)} (k_s | z_j - d_{jk}|) \left(\frac{z_j - d_{jk}}{|z_j - d_{jk}|} \right)^{n+1}, \\ Y_5 &= i \gamma_s k_s \sin(\theta - \beta) \psi_0 e^{i k_s [(z_j + d_j) e^{-i (k_j + (z_j + d_j) e^{i \beta}/2}, \\ \end{array}$$

where, $\lambda_1^* = (\lambda_c + \gamma_1 \alpha M)/\mu + 1$ and $\lambda_2^* = (\lambda_c + \gamma_2 \alpha M)/\mu + 1$. The coefficients matrix of Eq. (20) is $5N \times [5N(2\bar{n}+1)]$ (where, \bar{n} is the effective number of n during the calculation), which is not a square matrix and cannot be calculated. To multiply the left and right parts of Eq. (20) with $e^{-im\theta_j}$ ($m = 0, \pm 1, \pm 2,...$) and make the integral at the interval of $[-\pi, \pi]$ at the same time, a square matrix about the coefficient matrix is obtained, and the final theoretical solutions about the complex coefficients of $A_n^k - E_n^k$ are as follows:

$$\sum_{i=1}^{5} \sum_{n=-\infty}^{\infty} \sum_{k=1}^{N} \sum_{l=1}^{5} P_{il}^{knm} X_{ln}^{k} = Y_{i}^{m},$$
(21)

$$P_{il}^{knm} = \int_{-\pi}^{\pi} P_{il}^{kn} e^{-im\theta_j} d\theta_j,$$
(22a)

$$Y_i^m = \int_{-\pi}^{\pi} Y_i e^{-im\theta_j} d\theta_j.$$
(22b)

When calculating Eq. (21), the effective \bar{n} is determined by the receivable error:

$$\left|\frac{u_{\bar{n}+1} - u_{\bar{n}}}{u_{\bar{n}}}\right| \le 5 \%,\tag{23}$$

where, $u_{\bar{n}}$ and $u_{\bar{n}+1}$ are the total displacement in the saturated soils caused by all the P₁, P₂ and SV waves when the order n is valued \overline{n} and $\overline{n} + 1$ respectively.

5. Numerical results

The piles are assumed as uniform with radius of a, the spacing distance between the two adjacent piles are all d, and the SV waves are assumed to propagate vertically to axis x as shown in Fig. 2, and so $\beta = \pi/2$.



The normalized displacements of $|u_x/u_0|$ is introduced, where u_x is the displacement along axis x caused by both incident and scattered waves behind the barriers, and u_0 is the displacement caused by only the incident SV waves. The isolation effectiveness can be evaluated by the value of $(1-|u_x/u_0|) \times 100$ %, and so $|u_x/u_0|$ is more small and the isolation effectiveness is better.

In order to verify the reasonability of the theoretical methods s adopted in this paper, the theoretical results are compared with FEM, and the vibration isolation model of two rows of piles in elastic soils is selected as an example, as shown in Fig. 3 [11].



Fig. 3. Two rows of piles as barrier for vibration isolation

The vibration isolation fields are selected as a cube with each edge of 60.0 m, and the FEM model is built with software of ABAQUS, as shown in Fig. 4. The top of the model is free, the bottom and four sides of the model are defined as absorbing boundaries to avoid the wave reflection and they are divided into 21600 infinite elements, the inner soils are divided into 18000 eight node linear hexahedral elements with each edge of 1.0 m, the two rows of piles are assigned as beam elements with radius a = 0.2 m and adjacent distance d = 2.0 m, and the two rows of piles are tied to the soils at the interfaces.

There are some difficulties to reflect the water boundary conditions of the saturated soils with FEM, and so the elastic soils without water are adopted to simulate the vibration isolation problem to verify the reasonability of the theoretical results in this paper. The main physical and mechanical parameters of the elastic soils include elastic modulus E = 30 MPa, density $\rho = 1850$ kg/m³ and possio's ratio $\nu = 0.33$, and the ratio of the relative parameters of the piles to saturated soils are $\mu^* = 50$, $\rho^* = 1.35$ and $\nu^* = 0.75$. The distance between the load and piles is R = 2.0 m, and 17 piles are designed with the first row of 8 and second of 9 as shown in Fig. 2. The load is simplified as a harmonic load $F(t) = 800 \sin 20\pi t$, and acts on the top of the FEM model with action area of 1.0 m². The calculated displacements at the center of 20.0 m behind the two rows of piles are given in Fig. 5.





Fig. 5. Simulated displacements at the center of 20.0 m behind the two rows of piles



Fig. 6. Comparison curves between theoretically calculated and FEM simulated results

Fig. 5 shows that the maximum simulated displacement at the same place is 0.192 mm with

no piles and 0.137 mm with two rows of piles respectively, and so the displacement ratio $|u_x/u_0|$ is 0.715. As to the incident SV waves with frequency of 25π , the displacement ratio $|u_x/u_0|$ is 0.680, which is slightly less than the FEM simulated results. And the theoretically calculated curves of $|u_x/u_0|$ at the center of behind the two rows of piles along y axis ($0 \le y/a \le 300$) are compared with those simulated with FEM, as shown in Fig. 6.



Fig. 7. Contours of $|u_x/u_0|$ behind the barriers with pile row number when $k_s a = 1.0$

Fig. 6 shows that theoretically calculated results are almost near to FEM simulated results especially when y/a is farer that 60.0 with the maximum err of about 6%, and the theoretical

methods in elastic soils are proved to be reasonable, so the similar theoretical methods can be adopted to solve the vibration isolation problems in the saturated soils.

The physical and mechanical parameters are selected based on the silt soils [21] as shown in Table 1.



Fig. 8. Contours of $|u_x/u_0|$ behind the barriers with pile row number when $k_s a = 2.0$

The normalized incident frequency $k_s a$ is valued 1.0 and 2.0 respectively, the ratio of modulus

and density of piles to saturated soils are $\mu^* = 50$ and $\rho^* = 1.35$, three kind of arrangements of piles are selected, which are one row (8 piles), two rows (1st row of 8 piles and 2nd row of 9 piles) and three rows (1st row of 8 piles, 2nd row of 9 piles and 3rd row of 8 piles), and contours of $|u_x/u_0|$ in the areas of $(0 \le x/a \le 21, 0 \le y/a \le 500)$ behind the barriers are calculated and drawn as shown in Fig. 3 and 8, and the curves of $|u_x/u_0|$ at the position of x/a = 10.5 (which is just the center of the barriers) changing with y/a are given as shown in Fig. 9 and 10.



Table 1. Physical and mechanical parameters of the saturated soils

Fig. 10. Curves of $|u_x/u_0|$ at the center of the barriers when $k_s a = 2.0$

y/a

Fig. 7 and 8 show that: (1) the barriers compared of several rows of piles have better isolation effectiveness for higher incident frequencies such as $k_s a = 2.0$ than lower incident frequencies such as $k_s a = 1.0$, this is because that the incident waves with larger frequencies would be scattered and refracted by the piles more than those with lower frequencies, and so the incident waves would be consumed in large degree; (2) as to the incident SV waves with different frequencies, the displacements would be magnified by the piles and $|u_x/u_0|$ are larger than 1.0 at some locations near the barriers, and this phenomena would be obvious for the barriers with three rows of piles, which is caused by back and forth multiple scattering; (3) as to the incident SV waves with different frequencies and all barriers composed of different rows of piles, the center of the barriers have better isolation effectiveness than the two sides, this is because that the

incident waves would diffract from the two sides of the barriers; (4) as to the incident SV waves with different frequencies, the isolation effectiveness increase with the row number of piles.

Fig. 9 and 10 also show that, the row number of piles have different influence on the isolation effectiveness for different incident SV waves, as to the lower incident frequencies such as $k_s a = 1.0$, the isolation effectiveness increases with the row number of piles increase from one to three, while as to the larger incident frequencies such as $k_s a = 2.0$, the isolation effectiveness increases obviously when row number of piles increase from one to two and not obviously when row number of piles increase from two to three, and this is because that the incident with higher frequencies have shorter wavelengths and would be scattered by the piles intensely.

Take two rows of piles (1st row of 8 piles and 2nd row of 9 piles) as example. The modulus ratio μ^* is valued 10, 50, 100 and 500 respectively, other parameters are same to Figs. 7-10, and the curves of $|u_x/u_0|$ at the position of x/a = 10.5 of the barriers changing with y/a are given as shown in Figs. 11 and 12.



Fig. 11. Curves of $|u_x/u_0|$ with different modulus ratio when $k_s a = 1.0$



Fig. 12. Curves of $|u_x/u_0|$ with different modulus ratio when $k_s a = 2.0$

Figs. 11 and 12 show that as to the incident SV waves with different frequencies, the μ^* has the similar influence on the isolation effectiveness: the isolation effectiveness increases obviously when μ^* increases from 10 to 50 while not obvious when μ^* increases from 50 to 500, and so $\mu^* = 50$ is a better selection based on both isolation effectiveness and construction costs.

6. Conclusions

The conformal mapping method of complex functions and expansion method of wave

functions are adopted, the theoretical solutions about several rows of piles as barriers for incident SV waves in the saturated soils are obtained, and some conclusions are drawn out as follows:

1) The barriers composed of several rows piles are better for incident SV waves with higher frequencies than those with lower frequencies.

2) The pile number has different influence on the isolation effectiveness for the incident SV waves with different frequencies: as to incident SV waves with lower and higher frequencies, 3 and 2 rows of piles are the best selection respectively;

3) The stiffness of the piles has similar influence on the isolation effectiveness, and piles with modulus ratio of $\mu^* = 50$ is a better selection based on both isolation effectiveness and construction costs.

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