2587. Research on nonlinear inertance calculation method of ball-screw inerter

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Abstract. This paper proposes an inertance calculation method of the ball-screw inerter when considering the nonlinear factors, which will be called as the nonlinear inertance. With the inerter increasingly used in vibration damping systems, the effects that caused by the nonlinear factors, especially by friction, have attracted a lot of attention. The inertance is deemed as the inherent property of the inerter. In most systems, an analysis with the application of the inertance of the ideal linear inerter (which will be called as the ideal linear inertance) may produce a considerable error, thus the consideration of nonlinear factors becomes increasingly significant when calculating the inertance. Based on the Frenet frame theory and Hertz contact theory, the relative sliding speed and sliding friction torque between balls and raceways of the ball-screw inerter have been deduced, thus obtaining the sliding friction efficiency of the balls, and then introducing the rolling friction efficiency of the balls, and finally obtaining the method of the nonlinear inertance calculation after friction consideration. Based on the jobs above, an experimental study on the inertance of ball-screw inerter was carried out. In addition, the ideal linear inertance, nonlinear inertance and experimental value were compared with each other. The results show that the nonlinear inertance is closer to the experimental value; it demonstrates the accuracy of the calculation method proposed in this paper. Besides, the influences of other parameters of the ball-screw mechanism on the inertance are also analyzed, including the nominal radius, lead, contact angle, ball radius and number of balls, the comparative results show that the lead and ball radius have a great influence on the inertance. These results provide a more accurate inertance calculation method for designing a ball-screw inerter and its application in engineering vibration damping systems, meanwhile the theory and approach about inerter will also be developed.

Keywords: ball-screw inerter, Frenet frame, Hertz contact, friction efficiency, nonlinear inertance.

1. Introduction

Inertance, measured in kilograms, is the constant of proportionality of inerter which is a two-terminal mechanical device with the property that the applied force at its two terminals is proportional to the relative acceleration between them [1, 2]. In the past few years, the ideal inerters were applied to vehicle, motorcycle and train suspension systems, where significant performance improvements were achieved [3]. However, with the inerter increasingly used in the vibration damping systems, effects caused by nonlinear factors, especially the friction, attracted a lot of attention. In most systems, carrying on analysis with the application of ideal linear inertance may produce a considerable error, thus the consideration of nonlinear factors becomes increasingly significant when calculating the inertance.

The concept of inerter was first proposed by Smith at 2001 when researching the synthesis of mechanical networks, and two basic forms of inerter were given: ball-screw inerter and rack-pinion inerter [1, 4]. After that, Smith and his team carried out constant researches on inerter and its application. In 2004, the inerter was used by Smith team in vehicle suspension and high-performance motorcycle control systems [5]. In 2005, McLaren Eq. (1) Team won the championship because of the installation of an inerter [2] in the car damping system, which

enhanced the grip performance of the car. Starting from 2006, Fu-Cheng Wang has explored the application of inerter on train suspensions and buildings [6-11]. Since 2007, the Jiangsu University team in China verified that the inerter had the characteristics of passing high-frequency and resisting low-frequency. Besides, the ideal ceiling and floor damping of the vehicle suspension are realized by using the inerter. With this achievement, the Jiangsu University team applied for several international PCT patents [12-14]. In 2013, researchers of China North Vehicle Research Institute carried out an application research of inerter and the ISD on the high mobility crawler vehicle [15]. In 2015, Yinglong Hu et al. considered the problems of analysis and optimization of the inerter-based isolators based on an "anti-axial" single-degree-of-freedom isolation system [16].

With deepening the research, the influences of nonlinear factors (including friction, backlash, elastic effects) on the inerter were continually put forward. In 2008, considering that the gaps, friction and elastic forces exist in the inerter structure, Fu-Cheng Wang established the inerter nonlinear mechanic model and studied the influence of these nonlinear factors on the ISD suspension calculation process [17]. In 2009, the Smith team carried out a series of bench tests on a ball-screw inerter and rack-pinion inerter. They found that the backlash and elastic force existing in the mechanical inerter had a direct impact on its mechanical properties [18]. In 2016, a Jiangsu University team established a nonlinear inerter model and carried out mechanical properties tests based on the consideration of the friction, gap and screw elastic force. And finally, they had obtained the mechanical model structure of the inerter which was able to identify parameters [19]. In the same year, they proposed a new type of fluid inerter and analyzed the nonlinearities including friction and nonlinear damping force caused by the viscosity of fluid [20].

However, no one has yet proposed how to calculate the inertance taking into account of the nonlinear factors. The inertance calculation equation of the ideal ball-screw inerter is:

$$b_{liner} = \left(\frac{2\pi}{p}\right)^2 \cdot J,$$

where p represents the lead of ball-screw mechanism and s represents the rotary inertia of flywheel. Obviously, it didn't take nonlinear factors into account. [17-20] show that nonlinear factors have a great influence on the kinetic properties of inerter and damping performance of ISD suspension. With the increasingly use of inerter, inertance, inherent inerter property, its calculating accuracy gets particularly important, especially in low-frequency and precise damping systems. Thus, the inertance calculation method of the inerter, which influenced by nonlinear factors, needed to be further studied in the engineering applications.

The objective of this paper is to study the calculation method of nonlinear inertance of ballscrew inerter, which deduce the differential sliding velocity between the balls and raceways, thus obtaining the sliding friction efficiency of the balls, and finally obtaining the calculation method of the nonlinear inertance. Based on the jobs above, an experimental study on the inertance of ball-screw inerter was carried out. In addition, the ideal linear inertance, nonlinear inertance and experimental value are compared with each other. The results show that the nonlinear inertance is closer to the experimental value, it demonstrates the accuracy of the calculation method proposed in this paper. Besides, the influences of other parameters of the ball-screw mechanism on the inertance are also analyzed, including nominal radius, lead, contact angle, ball radius and number of balls. The organization of this paper is as follows. The assumptions in this paper are described in Section 2. Section 3 deduces the differential sliding velocity between balls and raceways. In Section 4, the sliding friction efficiency is deduced, and the rolling friction efficiency is quoted. Through the previous derivations, the inertance of ball-screw inerter is ultimately obtained in Section 5. In Section 6, the accuracy of the nonlinear inertance calculating method that mentioned in this paper was verified trough experiments, and the influences of some parameters of the ball-screw mechanism on the inerter are analyzed. Conclusions are drawn in Section 7.

2. Ball-screw inerter modal

Fig. 1 is a model of ball-screw inerter, including screw, nut, balls, flywheel, bearing and support frame; where the ball-screw mechanism is composed of a screw and ball nut, the nut is permanently connected with the inner bearing ring, and the outer bearing ring is permanently connected with the support frame, the flywheel is permanently connected with the nut in the support frame, a lifting lug is connected with the support frame, and another lifting lug is connected with the screw. The force acts on two lifting lugs, promoting the linear motion of the screw, and making the nut and the flywheel rotate at the same time.

To better deduce the nonlinear inertance of the ball-screw inerter, the following assumptions are made for the ball-screw inerter model shown in Fig. 1:

(1) Assume that the force of every working ball in the raceway is the same;

(2) Assume that two contact angles which between the ball and two raceways are approximately equal;

(3) Assume that the ball-screw mechanism model is fit for the Hertz contact model;

(4) Assume that the truncation error produced by the calculating process of frictional moment between ball and raceway can be negligible.

(5) Assume that the friction coefficient between the balls and raceways is constant.



Fig. 1. Model of ball-screw inerter

3. Differential sliding speed of contact point

3.1. Establishment of coordinate system

In order to carry out a dynamic analysis on the ball screw assembly in inerter, a World Coordinate System (O-XYZ) is established as shown in Fig. 2, in which the Z axis coincides with the central axis of the screw; Moreover, a Frenet frame is established based on the motion curve of the ball center to study the sliding phenomenon of the contact area between balls and raceway.



Fig. 2. Establishment of coordinate system

Set the base vector of the World Coordinate System as:

$$W = \begin{pmatrix} i \\ j \\ k \end{pmatrix}. \tag{1}$$

And the base vector of the Frenet frame as:

$$F = \begin{pmatrix} t \\ n \\ b \end{pmatrix},\tag{2}$$

where the basic vector of the Frenet frame can be expressed as [21]:

$$t = \frac{dR}{d\theta} / \left| \frac{dR}{d\theta} \right|,\tag{3}$$

$$b = \left(\frac{dR}{d\theta} \times \frac{d^2R}{d\theta^2}\right) / \left|\frac{dR}{d\theta} \times \frac{d^2R}{d\theta^2}\right|,$$
(4)
$$n = b \times t.$$
(5)

The curve equation of the ball center in the World Coordinate System can be expressed as:

$$R = \{r_o \cdot \cos\theta, r_o \cdot \sin\theta, r_o \cdot \theta \cdot \tan\alpha\}W,\tag{6}$$

where r_o represents the nominal radius of ball screw assembly, θ represents ball center track which is relative to the World Coordinate System turning angle, α represents spiral angle of the ball screw assembly.

To figure out the ratio between World Coordinate System and Frenet frame, take the derivative of Eq. (6) to obtain:

$$\frac{dR}{d\theta} = \{-r_o \cdot \sin\theta, r_o \cdot \cos\theta, r_o \cdot \tan\alpha\}W,\tag{7}$$

$$\frac{d^2R}{d\theta^2} = \{-r_o \cdot \cos\theta, -r_o \cdot \sin\theta, 0\}W.$$
(8)

Substitute Eqs. (6)-(8) into Eq. (3)-(5) and get:

$$F = \begin{bmatrix} -\sin\theta \cdot \cos\alpha \ \cos\theta \cdot \cos\alpha \ \sin\alpha \\ \cos\theta \ \sin\theta \cdot \sin\alpha \ -\cos\theta \cdot \sin\alpha \ \cos\alpha \end{bmatrix} W,$$
(9)

where define:

10

$$T_m^{-1} = \begin{bmatrix} -\sin\theta \cdot \cos\alpha & \cos\theta \cdot \cos\alpha & \sin\alpha \\ -\cos\theta & -\sin\theta & 0 \\ \sin\theta \cdot \sin\alpha & -\cos\theta \cdot \sin\alpha & \cos\alpha \end{bmatrix}.$$
 (10)

According to Eq. (9) and (10), one can know that:

$$W = T_m \cdot F. \tag{11}$$

And then by performing the matrix inverse of Eq. (10), one can obtain:

2587. RESEARCH ON NONLINEAR INERTANCE CALCULATION METHOD OF BALL-SCREW INERTER. HUABING WEN, JUNHUA GUO, YANG LI, YUE LIU, YEE PIN GOH

$$T_m = \begin{bmatrix} -\sin\theta \cdot \cos\alpha \ \cos\theta \ \sin\alpha \cdot \sin\theta \\ \cos\alpha \cdot \cos\theta \ \sin\theta \ -\sin\alpha \cdot \cos\theta \\ \sin\alpha \ 0 \ \cos\alpha \end{bmatrix}.$$
(12)

3.2. Contact point movement on ball

As shown in Fig. 3, in the Frenet frame, A can be defined as the contact point between the ball and the nut, while β_A as the contact angle, and B can be defined as the contact point between the ball and screw, while β_B as the contact angle.



Fig. 3. Schematic of contact point between ball and raceway

Transforming Eq. (7) into the Frenet frame:

$$\frac{dR}{dt} = \left\{ \frac{r_o}{\cos\alpha} \cdot \frac{d\theta}{dt}, 0, 0 \right\} F.$$
(13)

Assume that the instantaneous rotation angular velocity of the ball in the Frenet frame is as: $\omega_{ball} = \{\omega_t, \omega_n, \omega_b\}F$, then the instantaneous velocity of contact point A and B can be denoted as:

$$V_{A,ball} = \frac{dR}{dt} + \omega_{ball} \times \overrightarrow{O'A} = \begin{bmatrix} \frac{r_o}{\cos\alpha} \cdot \frac{d\theta}{dt} + r_b(\omega_b \cdot \cos\beta_A - \omega_n \cdot \sin\beta_A) \\ r_b \cdot \omega_t \cdot \sin\beta_A \\ -r_b \cdot \omega_t \cdot \cos\beta_A \end{bmatrix}^T F,$$
(14)

$$V_{B,ball} = \frac{dR}{dt} + \omega_{ball} \times \overline{O'B} = \begin{bmatrix} \frac{r_o}{\cos\alpha} \cdot \frac{d\theta}{dt} + r_b(\omega_n \cdot \sin\beta_B - \omega_b \cdot \cos\beta_B) \\ -r_b \cdot \omega_t \cdot \sin\beta_B \\ r_b \cdot \omega_t \cdot \cos\beta_B \end{bmatrix}^T F.$$
(15)

In the two equations above, r_b represents the ball radius, $\overrightarrow{O'A} = \{0, -r_b \cdot \cos\beta_A, -r_b \cdot \sin\beta_A\}F$, $\overrightarrow{O'B} = \{0, r_b \cdot \cos\beta_B, r_b \cdot \sin\beta_B\}F$.

3.3. Contact point movement on nut side

According to the drive principle of ball-screw mechanism, when the displacement of screw is L, the nut rotation angular displacement will be $2\pi L/p$, where p represents the screw lead. So, when the velocity of screw is dL/dt, the angular nut velocity can be written as:

$$\omega_{nut} = \left\{0, 0, \frac{d\theta}{dt}\right\} W = \left\{0, 0, \frac{2\pi}{p} \cdot \frac{dL}{dt}\right\} W.$$

Then transform it into the Frenet frame:

$$\omega_{nut} = \frac{d\theta}{dt} \{\sin\alpha, 0, \cos\alpha\}F.$$

Therefore, the velocity of contact point *A* on the screw can be denoted as:

$$V_{A,nut} = \omega_{nut} \times \overrightarrow{O'A} = \begin{bmatrix} r_b \cdot \frac{d\theta}{dt} \cdot \cos\alpha \cdot \cos\beta_A \\ r_b \cdot \frac{d\theta}{dt} \cdot \sin\alpha \cdot \sin\beta_A \\ -r_b \cdot \frac{d\theta}{dt} \cdot \sin\alpha \cdot \cos\beta_A \end{bmatrix}^T F.$$
(16)

3.4. Contact point movement of on screw side

The velocity of contact point B in the World Coordinate System can be denoted as:

$$V_{B,screw} = \left\{0,0,\frac{dL}{dt}\right\}W.$$
(17)

Transform the Eq. (17) into the Frenet frame:

$$V_{B,screw} = \left\{ \sin\alpha \cdot \frac{dL}{dt}, 0, \cos\alpha \cdot \frac{dL}{dt} \right\} F.$$
⁽¹⁸⁾

3.5. Relative sliding speed of contact point

For contact point A, according to the velocity $V_{A,ball}$ on the screw and the velocity $V_{A,nut}$ on the nut side, the relative sliding speed of contact A can be obtained as:

$$V_{A} = V_{A,ball} - V_{A,nut}$$

$$= \begin{bmatrix} \frac{r_{o}}{\cos\alpha} \cdot \frac{d\theta}{dt} + r_{b}(\omega_{b} \cdot \cos\beta_{A} - \omega_{n} \cdot \sin\beta_{A}) - r_{b} \cdot \frac{d\theta}{dt} \cdot \cos\alpha \cdot \cos\beta_{A} \\ r_{b} \cdot \omega_{t} \cdot \sin\beta_{A} - r_{b} \cdot \frac{d\theta}{dt} \cdot \sin\alpha \cdot \sin\beta_{A} \\ - \left(r_{b} \cdot \omega_{t} \cdot \cos\beta_{A} - r_{b} \cdot \frac{d\theta}{dt} \cdot \sin\alpha \cdot \cos\beta_{A}\right) \end{bmatrix}^{T} F.$$
(19)

Similarly, the relative sliding speed of contact *B* can be obtained as:

$$V_B = V_{B,ball} - V_{B,screw} = \begin{bmatrix} \frac{r_o}{\cos\alpha} \cdot \frac{d\theta}{dt} + r_b(\omega_n \cdot \sin\beta_B - \omega_b \cdot \cos\beta_B) - \sin\alpha \cdot \frac{dL}{dt} \\ -r_b \cdot \omega_t \cdot \sin\beta_B \\ r_b \cdot \omega_t \cdot \cos\beta_B - \cos\alpha \cdot \frac{dL}{dt} \end{bmatrix}^T F.$$
(20)

As shown in the Fig. 4 above, due to the existence of contact angle in the ball-screw mechanism, while the ball rolls along the curve of the ball center, a spin-slip phenomenon around the common

normal of contact point will also simultaneously exist.

The inerter is commonly used in the low and medium frequency workplaces, and the common normal of contact point A and B are both inside the normal plane b - n, so it can be approximately deemed as $\beta_A = \beta_B = \beta$. And according to the Plane Vector Geometry: $\overrightarrow{O'A} / \overrightarrow{O'B}$, it means the common normal of two contact points are duplicated. Assume the angular velocity when the ball is rotating around common normal as $\omega_{o'}$, and then obtain the following equation set:

$$\begin{cases} \omega_n = \omega_{o'} \cdot \cos\beta, \\ \omega_b = \omega_{o'} \cdot \sin\beta, \\ \omega_t = 0. \end{cases}$$

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Substitute the equation set Eq. (21) into the expression of V_A and V_B , and then convert it into angular velocity rotating around common normal:

$$\omega_{A} = \frac{2\pi}{p} \cdot \frac{dL}{dt} \begin{bmatrix} \frac{r_{o}}{r_{b} \cdot \cos\alpha} - \cos\alpha \cdot \cos\beta \\ -\sin\alpha \cdot \sin\beta \\ \sin\alpha \cdot \cos\beta \end{bmatrix}^{T} F,$$
(22)
$$\omega_{B} = \frac{1}{r_{b}} \cdot \frac{dL}{dt} \begin{bmatrix} \frac{r_{o}}{\cos\alpha} \cdot \frac{2\pi}{p} - \sin\alpha \\ 0 \\ -\cos\alpha \end{bmatrix}^{T} F.$$
(23)

4. Friction efficiency

4.1. Interface stress

As shown in Fig. 5, according to the Hertz contact theory, the shape of the contact surface between the ball, nut, and screw is ellipse. Assume that the major semi-axis of ellipse between the ball and nut is as a_A , and the minor semi-axis of ellipse is as b_A . Then, create a Cartesian coordinate system $(O_A - X_A, Y_A, Z_A)$ by taking the ellipse center as its origin, where X_A coincides with the major axis of the ellipse, Y_A coincides with the minor axis, and Z_A is perpendicular to the $X_A - Y_A$ plane; Similarly, the major semi-axis of the ellipse between the ball and screw is assumed as a_B , and the minor semi-axis is assumed as b_B . Then, create a Cartesian coordinate system $(O_B - X_B, Y_B, Z_B)$ by taking the ellipse center as its origin O_B , where X_B coincides with the major axis of the ellipse, Y_B coincides with the minor axis, and Z_B is perpendicular to the $X_B - Y_B$ plane. Due to space limitation, the calculating methods of major semi-axis a_A and a_B , minor semi-axis b_A and b_B can be referred to the methods in literature [22].



Fig. 5. Schematic of contact area between ball and raceway



Fig. 6. Schematic of contact stress distribution

As shown in Fig. 6, take the contact area between the ball and screw as an example, inside the ellipse contact area, the contact stress of every point is different. On axis Z_B , due to the maximum deformation, its contact stress reaches the maximum value too. The contact stresses of other points are distributed according to the half ellipsoid rule, and these values are [22]:

$$q = q_{\max} \sqrt{1 - \frac{x^2}{a_B^2} - \frac{y^2}{b_B^2}},$$
(24)

where, q represents the stress of random point inside the ellipse contact area, q_{max} represents the maximum stress inside the ellipse contact area.

The stress of the ellipse contact area is distributed according to the half ellipsoid rule, so the positive pressure acting on a single working ball, and the stress of every point has an integral ratio as follows [22]:

$$P_{in} = \int_{S} q dS = \frac{2\pi \cdot a_B \cdot b_B \cdot q_{\max}}{3},$$
(25)

where, S represents the elliptical contact area.

The Eq. (24) shows that in order to get the stress of random point inside the ellipse contact area, the maximum stress q_{max} inside it has to be calculated first. The following shows the solution of maximum stress q_{max} .

As shown in Fig. 7, due to the existence of lead angle α , the axial excitation load F_{in} acting on the single working ball is $F_{in} \cdot \cos \alpha$, so the positive pressure acts on a single working ball can be obtained as:

$$P_{in} = \frac{F_{in} \cdot \cos\alpha \cdot \sin\beta}{Z},\tag{26}$$

where Z represents the ball number in the ball-screw mechanism;

Synthesize Eq. (25) and Eq. (26), the maximum stress inside the ellipse contact area:

$$q_{\max} = \frac{3 \cdot F_{in} \cdot \cos\alpha \cdot \sin\beta}{2\pi \cdot Z \cdot a_B \cdot b_B}.$$
(27)

Fig. 7. Schematic of stress analysis

4.2. Friction torque of contact area

As shown in Fig. 9, take the contact between the screw and ball as an example, take a random point Q(x, y) in the coordinate system, and its area element dS = dxdy, then its friction can be obtained as $dF_f = \mu_f \cdot q_Q dxdy$, where μ_f refers to the friction coefficient of contact area, and q_Q refers to the positive pressure of point Q, continue to get its moment of the area relating to axis Z_B as $dM_B = \sqrt{x^2 + y^2} dF_f$ based on the length between Q and B.



Fig. 9. Schematic of contact surfaces infinitesimal

Thus, the moment of the whole area relating to axis Z_B :

$$M_B = \iint\limits_{S} \sqrt{x^2 + y^2} \cdot \mu_f \cdot q_Q dx dy.$$
⁽²⁸⁾

Decompose M_B to the three-coordinate axis of Frenet frame:

$$M_{BF} = \{0, M_B \cdot \cos\beta, M_B \cdot \sin\beta\}F,\tag{29}$$

where M_{BF} represents the friction torque of contact point B in the Frenet frame.

4.3. Sliding friction efficiency

The angular velocity and friction moment of two contact points *A* and *B* in the Frenet frame have been obtained above, so the its friction power is respectively shown as bellow:

$$P_{fA} = M_{AF} \cdot \omega_A, \tag{30}$$

$$P_{fB} = M_{BF} \cdot \omega_B, \tag{31}$$

where M_{AF} represents the friction moment of contact point in the Frenet frame;

The sliding friction power of single ball is a sum of two equations above:

$$P_f = P_{fA} + P_{fB}. (32)$$

Since the physical quantity power is a scalar, its value will not change along with the change of coordinate system, thus the excitation power while in the Frenet frame can be expressed as:

$$P_{total} = F_{in} \cdot \frac{dL}{dt}.$$
(33)

Define the energy consumption of rolling friction that account for the percentages of excitation as sliding friction efficiency:

$$\eta_1 = \frac{P_f}{P_{total}} \times 100\%. \tag{34}$$

4.4. Rolling friction efficiency

Define the energy consumption of rolling friction that account for the percentages of excitation as rolling friction efficiency [23]:

$$\eta_2 = \frac{\tan(\alpha - \arctan\mu_f')}{\tan\alpha} \times 100 \,\%,\tag{35}$$

where μ'_{f} represents the rolling friction coefficient, generally within the range of 0.0036-0.0038.

5. Nonlinear inertance

If this research does not take the consideration of nonlinear factors, according to the literature [1], the calculation approach of ideal linear inertance of the ball-screw inerter can be expressed as:

$$b_{linear} = \left(\frac{2\pi}{p}\right)^2 \cdot J. \tag{36}$$

To reduce the inertance calculation error of ball-screw inerter in the engineering damping system, based on the friction efficiency derived above, including the sliding friction efficiency and rolling friction efficiency. Introduce the nonlinear factors on the basis of ideal linear inertance of the ball-screw inerter, and finally the nonlinear inertance calculation method can be deduced in this paper, which is:

$$b_{nonlinear} = \left(\frac{2\pi}{p}\right)^2 \cdot J \cdot (1 - \eta_1 - \eta_2). \tag{37}$$

6. Experimental verification of nonlinear inertance and parameter evaluation

6.1. Experimental verification

To verify the derived nonlinear inertance and the accuracy of this calculation method, this research has performed a testing experiment to the inertance of the ball-screw inerter on the ES-10-240 vibration test bench, the parameters of ball-screw mechanism are shown in Table 1 and the setup of vibration test bench is shown in Fig. 10. This experiment has also installed a force sensor at the connection between the ball-screw inerter and the vibration test bench, and installed several vibration acceleration sensors on the vibration table at the same time, and then the DASP V10 model analyzer was used for simultaneous vibration data acquisition and signal spectrum analysis, by loading different frequencies of square wave excitation to the ball screw inerter, which are within the frequency range from 1 Hz to 12 Hz, for every rises of 1 Hz per test, and the experimental results of force and acceleration are recorded in Table 2.

Table 1. Independent geometric parameters of ball-screw mechanism

Name	Scale	
Nominal diameter (r_o)	6 mm	
Screw lead (p)	10 mm	
Contact angle (β)	45°	
Ball diameter (r_b)	3 mm	
Ball number (Z)	18	









b) Living photos

Fig. 10. Experimental device and its live photos

I able 2. Experimental test force and acceleration							
Frequency (Hz)	Force (N)	Acceleration	Frequency (Hz)	Force (N)	Acceleration		
1	6.64	0.18	7	105.47	2.56		
2	9.92	0.26	8	130.49	3.29		
3	16.79	0.42	9	168.41	4.12		
4	35.65	0.90	10	208.03	5.21		
5	54.41	1.35	11	248.80	6.25		
6	79.09	1.94	12	308.17	7.69		

Table 2 is organized into Fig. 11, with the frequency rises, both ends of the force and acceleration have the increasing trend, according to the definition of inertance: $\Delta F = b \cdot \Delta a$, where ΔF is the force of two ends, Δa is the acceleration of two ends, the experimental results of inertance can be obtained, as shown in Table 3.

2587. RESEARCH ON NONLINEAR INERTANCE CALCULATION METHOD OF BALL-SCREW INERTER. HUABING WEN, JUNHUA GUO, YANG LI, YUE LIU, YEE PIN GOH

Frequency (Hz)	Experiment-b (kg)	Frequency (Hz)	Experiment-b (kg)
1	37.6	7	41.2
2	38.1	8	39.7
3	39.6	9	40.9
4	39.4	10	39.9
5	40.3	11	39.8
6	40.8	12	40.1

Table 3. Experimental results of inertance

After that, fit the test points by using MTALAB cftool tool box. In addition, the experimental fitting curve is compared with the nonlinear inertance and the ideal linear inertance, thus obtaining the following graph Fig. 12.

These experimental results show that the nonlinear inertance derived by this paper is slightly larger than the experimental fitting results. When the frequency is less than 12 Hz, the error between the ideal linear inertance and the experimental inertance is greater than 10.0 %. It can be known that the ideal linear inertance is not accurate enough in a low frequency range of the ball-screw inerter, in some demanding vibration isolation systems; the use of the ideal linear inertance will introduce large errors. However, when the frequency is less than 12 Hz, the error between the nonlinear inertance and the experimental inertance is less than 5.0 %, especially when the frequency is greater than 3 Hz, the error between the nonlinear inerterace and experimental inertance is less than 2.0 %. Obviously, compare with the ideal linear inertance of the ball-screw inerter, the nonlinear inertance derived in this paper is more accurate in the low frequency range of the inerter.



6.2. Parameter evaluation

Fig. 13 shows the comparison of the influence of different lead on the nonlinear inertance, the

other parameters are shown in Table 1. At the same frequency, the smaller is the lead, the greater is the inertance, the stronger is the nonlinearity of inertance of the ball-screw inerter. Additionally, as the frequency increases, the nonlinear inertance tends to be stable.



Fig. 13. Influence of different lead on inertance

Fig. 14-Fig. 17 show the comparison of the effects of nominal radius, ball radius, ball number, and contact angle on the nonlinear inertance respectively, the other parameters are shown in Table 1. At the same frequency, the larger is the nominal screw radius, the smaller is the ball radius, the larger is the number of balls and the larger is the contact angle, so the nonlinear inertance will be smaller, and the nonlinearity of the inertance will be stronger. In addition, within 3 Hz, the frequency has a great influence on the nonlinearity of inertance, the smaller is the frequency, the stronger will be the nonlinearity of inertance and the smaller will be the inertance; As the frequency increases, the nonlinear influence of the frequency on the inertance is weakened, and the inertance tends to be stable.



This paper only considers the influence of the parameters of the ball-screw mechanism in the inerter, and the coupling between the friction coefficient and the friction torque is not taken into account. [24] shows that the change in friction torque will change the viscous resistance of the lubrication medium to which the ball is subjected, which will result in a different friction coefficient. With the increase of the frequency, the difference in the friction coefficient will result in the nonlinear inertance close to the ideal linear inertance in a fluctuating shape.

Wen's main task is to put forward the main idea of the article, research the project and guide the article writing. Guo research the project and write article based on research results. Li and Liu modify the article according to the results of the study. Goh's main task is to polish the article.

7. Conclusions

With the Frenet frame theory and Hertz contact theory, this paper has validated the calculation method accuracy of the nonlinear inertance based on the proper simplification of ball-screw inerter model and some experiments. In addition, the parameters of the ball-screw mechanism have some influences on the non-linear inertance too. Some conclusions are summarized as follows:

1) At low frequencies especially below 3 Hz, the error between the ideal linear inertance and the actual inertance of the ball-screw inerter is greater than 10 %. In some demanding vibration isolation systems, the use of ideal linear inertance will produce a large error.

2) The error of the nonlinear inertance of ball-screw inerter deduced in this paper is less than 5 % at low frequencies compared with the actual inertance, especially when the frequency is larger than 3 Hz, its error is less than 2 %. Therefore, the inertance of ball-screw inerter deduced in this paper is relatively accurate.

3) The smaller is the ball-screw inerter lead, the greater is the inertance, and the stronger is the nonlinearity of inertance. In addition, at the same frequency, the larger is the nominal radius of the screw, the smaller is the ball radius, the larger is the number of balls and the larger is the contact angle, so the nonlinear inertance will be smaller, and the nonlinearity of the inertance will be stronger.

4) Within 3 Hz, the frequency has a great influence on the nonlinearity of inertance, the smaller is the frequency, the stronger will be the nonlinearity of inertance, and the smaller will be the inertance; As the frequency increases, the nonlinear influence of the frequency on the inertance is weakened, and the inertance tends to be stable.

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