2485. The experimental investigation of nonlinear impact behaviors of partial rolling bearing with failure collision

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Abstract. The influence of nonlinear impact behavior of partial failure collision rolling bearing is studied when the pitfall occurs in the inner ring. The piecewise non-smooth experiment model of three-degree-of-freedom rolling bearing system with fault in inner ring is established, and the experiment method is used to further reveal the bifurcations and chaos of bearing system through establishing the Poincaré mapping on the collision plane. The nonlinear impact behaviors are studied with the vibration force increasing, when the vibration frequency is determined. The impact signal spectrum characteristics are analyzed through changing the quality of the mass M_1 after adding a force sensor. The study of nonlinear impact behaviors of the fault bearing system provides reliable basis for the design and fault diagnosis and provides theoretical guidance and technical support for the actual design in the safe and stable operation of large high-speed rotating machinery.

Keywords: bearing, nonlinear behavior, collision, Poincaré mapping.

1. Introduction

Bearing is one of the most important part of the rotating machinery and is widely used in aviation aircraft, spacecraft, rocket, ship, nuclear reactor and so on. It plays a key role for the safe running of the equipment. Once the bearing is broken, it will cause more serious consequences including security challenge and increasing downtime. The nonlinear vibration behaviors caused by fault are very complicated, which has significant effect on the design and fault diagnosis of bearing. Utkarshkumar Arvindbhai Patel et al. established the inner and outer ring fault mathematical model to research the nonlinear dynamic behavior of bearing fault according to Hertz contact theory. Because of the contact force and process exhibit nonlinear stiffness characteristic, the nonlinear behavior of the system can be analyzed through n-dimension nonlinear simultaneous equations and numerical simulation [1]; N. Tandon and A. Choudhury studied the impact of the bearing by assuming different shape signals, such as sine wave, square wave and triangle wave respectively in the bearing inner ring and outer ring, rolling elements and the accuracy of the numerical simulation is proved by comparing the signal frequency response of the simulation and experimental data [2]; N. Tandon et al. described the fault characteristics through the frequency spectrum characteristics by considering the rolling bearing as a three degree of freedom fault model which is consisted of inner and outer ring raceway waviness, size irregular rolling body and so on [3]; N Sawalhi and R. B. Randall analyzed the process of rolling element in and out of fault location and studied the vibration of the bearing behavior pricewise by assuming the fault position signal as a step type which stirs up a more broad band frequency response and the cestrum analysis method is used to determine the fault location and fault type [4]; Alireza Moazenahmadi et al. [5] studied impact type and the signal response of the fault signal by

establishing the failure nonlinear dynamic bearing model and assuming the fault is the rectangular shape with the size of the fault changing from small to large; He et al. [6] studied the dynamics of rolling bearing system with gaps under variable load. Zhang et al. [7] investigated the nonlinear dynamic characteristics and stability of rolling bearing system under the influence of non-equilibrium force. Tang et al. [8] obtained the nonlinear bearing force in work situation based on Hertz elastic contact theory and the kinesiology of rolling bearing, and analyzed vibration features of rolling bearing system. Gao et al. [9] established the dynamical equations of the spindle system based on Hertz contact force model and studied mechanism and pathway of the instability.

The rolling bearing with failure may be categorized into vibro-impact systems. The strong nonlinear characteristic caused by instantaneity and discontinuity of impact can exhibit rich dynamic behavior. Many researchers have investigated the bifurcation and chaos of vibro-impact systems by theoretical and numerical analyses. Shaw and Holmes [10] studied the oscillator system with single piece rigid constraint under the action of a harmonic force by using modern dynamics theoretical methods and analyzed the bifurcation behaviors of periodic motions by using center manifold and normal formal theory. Jin and Hu [11] studied the vibro-impact response of Hertz contact model of flexible beam system in theory and carried on the corresponding experiment verification. Luo and Xie [12] studied the complexity co-dimension-two bifurcation behaviors of a two-degree-of-freedom vibro-impact system. Ding and Xie [13] studied the path from torus bifurcations to chaos in a three-degree-of-freedom vibro-impact system. Jin and Lu [14] studied the Lyapunov exponent of vibro-impact system with constraining stops through building compound map and obtained the determinant indexes to differentiate periodic motion and chaos. Yue et al. [15] studied the co-dimension-two bifurcation and chaos of symmetric period motion in a three-degree-of-freedom vibro-impact system and pointed out that symmetry can suppress some kinds of bifurcations of system.

Though these researches acquire many achievements, most of the researches center on theoretical and numerical analyses of dynamical behaviors in the rolling bearing with failure. This paper studies the dynamic behavior of partial rolling bearing with failure collision via experiment. The design of the testing rig and experimental procedures are described in detail. The rich dynamical behaviors are revealed by changing excitation frequencies.

2. The theoretical model design

The three-degree-of-freedom vibro-impact physical model of rolling element bearing with fault in inner ring is shown in Fig. 1, in it (a) is the whole bearing with failure, (b) is the particle rolling bearing with failure collision, (c) is the experimental simplified figure of partial rolling bearing with failure collision. The purpose is to study the nonlinear behaviors of partial rolling bearing with failure collision, and the Fig. 1(c) is designed to make sure the collision is point collision to simulate the collision in real bearing.



Fig. 1. The three-degree-of-freedom vibro-impact model of bearing

If there is only small fatigue pitfall on the inner ring, the system is described as rigid impact when the rolling elements pass through this fault location. The three-degree-of-freedom dynamic model with gap between M_1 and M_2 is shown in Fig. 1, in which, M_1 , M_2 , M_3 are equivalent mass block of shaft with inner ring, rolling elements and outer ring with base coordinate respectively. The corresponding vertical displacements are X_1 , X_2 and X_3 respectively. The damping coefficients are C_1 , C_2 and C_3 , the stiffness parameters are K_1 , K_2 and K_3 . The exitation force is $F_1 \sin(\Omega_1 t)$ acting on M_1 , Ω_1 is the revolution frequency of the shaft. The gap between M_1 and M_2 in the vibration system is D, and the recovery coefficient is R. When the relative displacement between the M_1 and M_2 is more than D, the rigid collision occurs.

3. Design of the experimental rig

3.1. The design of springs

The spring is one of the most important component in the experiment. The material used is 65 Mn, and the parameters designed of the three kinds of springs are in Table 1. In order to be convenient to install, the design of the spring is set to lap two coils up and down respectively, so the total number of coils is 4 more than the working coils.

(1) Allowable shearing stress $[\tau]$ (kg/mm²) of the spring is classified with three classes, I, II and III. The stress values are 0.3 σ_b , 0.4 σ_b , 0.5 σ_b respectively. Stress σ_b is 130 (kg/mm²) when the steel wire diameter *d* is 6 mm.

(2) The shear modulus G (N/m²) and the modulus of elasticity E (N/m²). When the diameter d of the steel wire is $4 \ge d \ge 0.5$ (mm), the range of G is 8300-8000 (N/m²), and E is 20750-20500 (N/m²), while d > 4 (mm), the value of G is 8000 (N/m²), and E is 20000 (N/m²).

(3) Spring index C, $C = D_2/d$, where D_2 is the mean diameter of the coil.

(4) Spring curvature K, K = (4C - 1)/(4C - 4) + (0.615/C).

(5) Effective coils $n, n = F_n G d^4 / 8P_n D_2^3$.

(6) Spring stiffness P', $P' = Gd^4/8D_2^3n$.

(7) Gap between coils, δ_1 , $\delta_1 \ge 0.1d$.

(8) Pitch of spring $t = d + F_j/n + \delta_1$, generally, $t \approx D_2/3 - D_2/2$, and space of coils is $\delta = t - d$.

(9) Free height of spring $H_0 = n * t + 3.5d$.

| Spring type | Type I | Type II | Type III |
|--------------------------|------------|-----------|-----------|
| Outer diameter of spring | 36 mm | 33.9 mm | 40.3 mm |
| Steel wire diameter | 6 mm | 6 mm | 6 mm |
| Spring stiffness | 8 kg/mm | 12 kg/mm | 4 kg/mm |
| Working coils | 6 coils | 5 coils | 8 coils |
| Total number of coils | 10 coils | 9 coils | 12 coils |
| Spring natural length | 81 mm | 71 mm | 105 mm |
| Maximal load | 140.234 kg | 147.47 kg | 126.57 kg |
| Maximal stroke | 17.529 mm | 12.357 mm | 31.528 mm |
| Working stroke | 10 mm | 10 mm | 10 mm |

Table 1. The designed parameters of the three kinds of springs

3.2. Measurement of the spring stiffness

The stiffness values of springs are measured in lab. The designed theoretical values of the springs are $k_1 = 8$ kg/mm, $k_2 = 12$ kg/mm, $k_3 = 4$ kg/mm. The spring force versus displacement curves are shown in Figs. 2, 3 and 4. The stiffness values are calculated from the slope of the curves. The slopes of this lines are $k_1 = 7.1270$ kg/mm, $k_2 = 12.3905$ kg/mm, $k_3 = 3.8927$ kg/mm. The unit kg/mm is transformed into N/m, and we can get that the stiffness of the springs are $k_1 = 69845$ N/m, $k_2 = 121427$ N/m, $k_3 = 38148$ N/m, which are used in the

following experiment.



Fig. 4. The diagram of the weight and the deformation of Spring 3

3.3. Components of the experimental rig

The experimental rig is composed of springs, collision ball, signal generator, data acquisition system, as shown in Fig. 5. A part of magnified figure is shown in Fig. 6. The collision structure comprises three mass blocks, the M_1 , M_2 and M_3 , and the springs K_1 , K_2 and K_3 . The linear rails which are vertical to the base are applied to ensure the mass blocks slide freely on the road. The collision part is designed between the M_1 and M_2 , and the gap between the M_1 and collision ball is D, as shown in Fig. 7.



Fig. 5. The experimental rig figure: (1) is the vibration generator, (2) is the collision test bench, (3) is the foundation bed, (4) is the guide rail



Fig. 6. The partial magnified figure of the main body of the experimental bench: (1) is the laser displacement sensor, (2) is the mass block M₁, (3) is the mass block M₂, (4) is the mass block M₃, (5) is the spring k₁, (6) is the spring k₂, (7) is the spring k₃, (8) is the force sensor

2485. THE EXPERIMENTAL INVESTIGATION OF NONLINEAR IMPACT BEHAVIORS OF PARTIAL ROLLING BEARING WITH FAILURE COLLISION. Xing He, Qiang Wang, YongBao Liu, Huidong Xu, Shuyong Liu



Fig. 7. The collision position setting: (1) is the collision point

The force sensor is installed to connect the mass block M_1 and the vibration generator. The base is heavy enough to ensure its stability without vertical displacement when the collision occurs. The signal generator shown in Fig. 8 generates sinusoid, triangle, rectangular waveform. The power amplifier is shown in Fig. 9. Three laser displacement sensors whose linear accuracies are 5 um are applied to measure the vertical displacement of M_1 , M_2 and M_3 . The force sensor is used to obtain the excitation force curves and the type of multi-channel data acquisition system is the NI PXIE-1062Q.



Fig. 8. The signal generator



Fig. 9. The power amplifier for exciter

4. Experiment study of the bifurcation and chaos responses in the system

4.1. The influence of increasing excitation force on the system

The chaotic response of the three-degree-of-freedom impact collision system under different excitation amplitude and constant frequency is researched in the section. The mass of the blocks are $M_1 = 1.370$ kg, $M_2 = 1.700$ kg, $M_3 = 3.260$ kg, respectively, and their test stiffness values are $K_1 = k_1 = 69845$ N/m, $K_2 = k_2 = 121427$ N/m, $K_3 = k_3 = 38148$ N/m. The gap *D* between M_1 and the top of the collision ball is 1.96 mm at the stationary state.

When the excitation frequency is 34.59 Hz, the power amplifier gain is adjusted to increase the amplitude of the excitation force, and the response is observed. When the gain is 0.5, the system response of M_1 is shown in Fig. 10 and Fig. 11.

When the power amplifier gain is 0.65, the response is shown in Fig. 12 and Fig. 13.

2485. THE EXPERIMENTAL INVESTIGATION OF NONLINEAR IMPACT BEHAVIORS OF PARTIAL ROLLING BEARING WITH FAILURE COLLISION. XING HE, QIANG WANG, YONGBAO LIU, HUIDONG XU, SHUYONG LIU



Fig. 13. The phase diagram and Poincaré section of M_2

When the power amplifier gain is increased continuously to 1.7, the system response of M_2 is shown in Fig. 14 and Fig. 15.

When the amplifier gain is increased continuously to 1.8, the system response of M_2 is shown in Fig. 16 and Fig. 17.

From the Figs. 10-17, we can see that the nonlinear behaviors occur in the system with the excitation force increasing when the frequency of the excitation is kept as a constant. When the force is small, that is the amplifier gain is 0.5, the system is in the stable period-1 motion initially, and we can see it in the reconstructed phase diagram of M_2 in Fig. 11, and the time response and

the signal spectrum diagram is shown in Fig. 10. When the power amplifier gain is increased gradually and passes through the bifurcation point, the period-2 motion appears, just like the power amplifier gain is 0.65 which is near the critical gain, the period doubling bifurcation occurs. The phase diagram and Poincaré section of M_2 is shown in Fig. 13, and we can see that there is a little gap in the phase diagram, and the Poincaré mapping appears two points, from which we can fix it that the system is in stable period-2 motion. With the increasing of the amplifier gain, the gap is increasing in the phase diagram of M_2 in Fig. 15, and the Poincaré section of it appears two points, and the time response and the signal spectrum diagram of M_2 at the gain is in Fig. 14. When the amplifier gain is increased continuously to 1.8, the period-4 motion of the system appears. There are four mainly points in the Poincaré section of M_2 in Fig. 10, Fig. 12, Fig.14, and Fig. 16. It is interesting that the attractor and Poincaré section of the reconstructed attractor is more sensitive to the change of the force than the spectra diagram.







Fig. 16. The time response and the signal spectrum diagram of M_2

According to the analysis above, the system is in the stable period-1 motion initially, and with

80

w(Hz)

b)

t(s) a) the decreasing of the amplifier gain, the critical bifurcation point appears, and then the period-2 motion and the period-4 motion of the system appear.



4.2. The influence of increasing excitation force of the system after increasing the mass of M_1

The force sensor is added to connect the exciter and the mass M_1 and the mass M_1 is increasing to $M_1 = 1.475$ kg, and the other parameters remain the same. The variation of the spectrum structure is researched from the force waveform signal. The excitation frequency is fixed at 32.6 Hz, the bifurcation behavior is observed. When the power amplifier gain is 0.5, the system has a single periodic motion, and the response diagrams of M_2 are shown in Fig. 18.

The spectrum and time response diagrams of the force are shown in Fig. 19. The collision frequency is the same as the exciter frequency in the single periodic state, which means the spectral characteristic is confirmed by experimental data.









When the gain is increased to 1.0, there is a deformation on the collision point, and the time response diagram, the phase diagram and the signal spectrum diagram of M_2 are shown in Fig. 20. The time response and the signal spectrum diagrams of force are shown in Fig. 21.

When the frequency 32.6 Hz, and the gain is increased to 1.2, the bifurcation behavior M_2 occurs, and the responses are shown in Fig. 22, from which we can see the bifurcation behaviors of the system. The corresponding dynamical force time series and spectrum are shown in Fig. 23.



Fig. 20. The time response diagram, the phase diagram and the signal spectrum diagram of M_2



Fig. 21. The time response diagram and the signal spectrum diagram of force



Fig. 22. The time response diagram, the phase diagram and the signal spectrum diagram of M_2

When the frequency keeps constant, and the amplitude of force is increased further to 1.7, the bifurcation occurs clearly, and the responses are shown in Fig. 24. The corresponding dynamical force time series and spectrum are shown in Fig. 25.







Fig. 24. The time response diagram, the phase diagram and the signal spectrum diagram of M_2



Fig. 25. The time response diagram and the signal spectrum diagram of force

When the amplifier gain is increased further to 2.0, the chaotic vibration is obtained, and the phase attractor, time series, and spectrum are observed, as shown in Fig. 26. The corresponding dynamical force time series and spectrum are shown in Fig. 27.

According to the research, it is obvious that when the force sensor is added, the mass of M_1 is increased to $M_1 = 1.475$ kg, the chaotic response of the system changes as well. When the force amplitude increases, the system behaviors varies from the single periodic motion to the double periodic and chaotic vibration. The peak of the 2X spectrum also increases, as shown in Figs. 18, 20, 22, 24, 26. When the frequency is invariant, the peak of the harmonics varies slowly, as shown in Figs. 19, 21, 23, 25, 27. The response is very different from the responses in section 3.1, and

we can see that the change of the quality parameters of the system will have a huge impact on the nonlinear behavior of the system.



Fig. 26. The time response diagram, the phase diagram and the signal spectrum diagram of m_2



Fig. 27. The time response diagram and the signal spectrum diagram of force

Qiang Wang put forward the theme and idea of this manuscript and wrote this manuscript. Yongbao Liu conducted the simulation. Huidong Xu determined the experiment scheme. Shuyong Liu help to build the experimental platform. Xing He conducted the experimental data.

5. Conclusions

The experimental rig of three-degree-of-freedom bearing local fault simulation system is designed, and the nonlinear behaviors of the system are researched with the variation of the excitation force and mass of the block. Results show that when the excitation frequency keeps constant, the system exhibits the single periodic motion, double bifurcation and period-4 response with the increasing of the force amplitude. When the force sensor is installed, both the block mass and the bifurcation dynamical characteristic are changed. The single period, period-2 and chaotic response are obtained. The displacement, force time series and spectrum diagram are analyzed. It is discovered that the main frequency peak of the response is changed when the force amplitude is increased, even if the excitation frequency keep constant, which means the complexity of the vibro-impact system.

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