Analysis of transverse vibration of wire rope in flexible hoisting system

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Abstract. Taking the flexible hoisting system as the research object, the transverse vibration mathematical model of the steel wire rope of the winding hoisting system is established, and the transverse vibration displacement, velocity and acceleration of the steel wire rope at 18 m away from one end of the wire rope under the sine excitation at other end is obtained through the comparison and verification of the results from the mathematical model, mechanical model and experiment. The results show that: for the length of the wire rope of 20 m, under the sine excitation at the end of amplitude of 0.002 m and frequency of 20 Hz, the transverse vibration displacement at 2 m from one end of the wire rope with length of 20 m is basically between -0.005 m and 0.005 m, the transverse vibration velocity at 2 m from one end of the wire rope is basically between -0.5 m/s² and 0.5 m/s². And the transverse vibration displacement decreases near the end, meanwhile the transverse vibration frequency increases with the reduction of length of vertical section of wire rope.

Keywords: flexible, wire rope, transverse, vibration.

1. Introduction

During the running process of mine hoisting system, the steel wire rope has property of elasticity, so there are longitudinal vibration and transverse vibration of the steel wire rope. For longitudinal vibration, the study of its vibration can obtain the stress of the steel wire rope in the running process, and then adjust the parameters to reduce its dynamic tension, then preventing the wire rope from breaking up, and so improving the safety factor of the hoisting system. For the transverse vibration, the transverse vibration of the steel wire rope can be obtained by studying its vibration under a certain excitation, and then the transverse vibration of the steel wire rope can be reduced by adjusting the parameters, so as to prevent the problem of rope skipping and improve the operation stability of the hoisting system. Bao Jihu [1] and Kou Baofu [2] established the mechanical model of the flexible lifting system based on Hamilton principle and carried out the numerical solution. Wang Wen [3] transformed the lateral vibration partial differential equation of elevator wire rope into a variable coefficient ordinary differential equation by using the finite difference method and solved it. A. Mankowski [4] studied the internal damping characteristics in the transverse vibration of the steel wire rope of the mine hoisting system, and analyzed the influence of internal damping on the transverse vibration. Blodgett R. E. [5] studied the transverse vibration of the tail rope of the elevator lifting system, and proposed the method of introducing damper on the upper end of the wire rope for its transverse vibration control. H. M. Irvine [6] introduced the modeling, statics and dynamics of Cable in his book Cable Structures. G. Tagata [7] studied the transverse vibration of the steel wire rope driven by axial harmonics with one end fixed and the other end driven by axial harmonics. J. A. Wickert. [8] made A modal analysis of the string and beam models with typical axial motion, and obtained the relationship between their motion speed and vibration frequency. Shabana A. A. [9] specifically introduced the freedom and forced lateral vibration of strings in Vibration of Discrete and Continuous Systems. Bilbao S. [10]

took the transverse vibration equation of nonlinear string as the starting point, and respectively uses the Kirchhoff carrier equation, digital waveguide equation and linear wave equation to solve it and analyzes the solution results. In this paper, the flexible hoisting system was took as the research object, the transverse vibration mathematical model of the steel wire rope of the flexible hoisting system was established, and through the comparison and verification of the solution results of the mechanical model, obtaining the transverse vibration displacement of the steel wire rope at two positions under the sine excitation at one end. The above research mainly adopts the variational method to obtain the transverse vibration equation of steel wire rope in flexible hoisting system. In this paper, Euler equation is adopted to derive the transverse vibration equation, which simplifies the derivation process. At the same time, this paper adopts the mechanical model and experimental model to verify the numerical model, which effectively verified the mathematical model, and can realize the further modification and improvement of the mathematical model, and thus provide better date support for the reasonable establishment of the mathematical model of transverse vibration of wire rope in hoisting system. The research and analysis of this paper has some reference basis for the modeling and analysis of the super deep flexible hoisting system, and has some reference and application value for the design of the super deep flexible hoisting system.

2. Establishment of mathematical model

For the flexible hoisting system, the wire rope is divided into the vertical rope, the suspended rope and the wire rope wound on the drum. Since the vertical part of the steel wire rope is directly responsible for the load bearing and lifting in hoisting system, and its length is constantly changing during the hoisting process, its transverse vibration displacement needs to be paid attention to, and its performance will directly affect the stability and safe operation of the hoisting system. Therefore, the vertical rope is selected to analyze the transverse vibration displacement, velocity and acceleration of steel wire rope of the hoisting system. Before establishing the transverse vibration model of the hoisting system, there are 2 assumptions: 1) The elasticity modulus of steel wire rope is constant during the hoisting process of the system. 2) The longitudinal vibration of the steel wire rope is ignored in order to analyze the transverse vibration of the steel wire rope. The established model of flexible hoisting system is shown in Fig. 1. In order to analyze the transverse vibration displacement of wire rope of the hoisting system, the mathematical model should be established first. The Hamilton function of steel wire rope is obtained from the energy of hoisting system, and its mathematical equation is obtained from Euler equation. Establishing the cartesian coordinate system as shown in Fig. 1. The length of the vertical rope of the hoisting system is set as l(t), the mass of the hoisting weight is m, and there are guide rails on both sides of the hoisting weight to limit its transverse vibration of load. The hoisting system has sine excitation at the wheel, the excitation amplitude is h, and the excitation frequency is w. The details are shown in Fig. 1.

First, find the kinetic energy of the steel wire rope with a length of l(t). Set the speed of each point of the wire rope as v. Then we can obtain the following expression:

$$v_x = v, \quad v_y = \frac{\partial y(x,t)}{\partial t} = vy_x + y_t, \tag{1}$$

$$T = \frac{1}{2}\rho \int_0^{l(t)} [v^2 + (vy_x + y_t)^2] dx.$$
 (2)

Then, find the potential energy of the steel wire rope. The potential energy of wire rope includes stretching strain energy and bending strain energy of wire rope. The potential energy is:

$$V = \frac{1}{2} EA \int_0^{l(t)} \varepsilon^2 \, dx + \frac{1}{2} EI \int_0^{l(t)} y_{xx} \, dx.$$
(3)

Within that, *EA* and *EI* are respectively the stretching and bending stiffness of the steel wire rope, ε and y_{xx} are the stretching strain and the bending strain. The stretching strain ε can be obtained from the equation below:

$$\varepsilon = \frac{ds - dx}{dx} = \frac{\sqrt{1 + \left(\frac{dy}{dx}\right)^2 dx - dx}}{dx} = \frac{\left[1 + \frac{1}{2}\left(\frac{dy}{dx}\right)^2\right] dx - dx}{dx} = \frac{1}{2}y_x^2.$$
(4)

Within that, ds is the arc length of steel wire rope corresponding to dx during the process of transverse vibration of steel wire rope. Then the Taylor formula is used to expand $\sqrt{1 + \left(\frac{dy}{dx}\right)^2}$ and ignore the higher order term. So the potential energy can be expressed as:



Fig. 1. Transverse vibration model of hoisting system

Then the Lagrange function is:

$$L = T - V = \frac{1}{2}\rho \int_{0}^{l(t)} [v^{2} + (vy_{x} + y_{t})^{2}] dx - \frac{1}{8}EA \int_{0}^{l(t)} y_{x}^{4} dx - \frac{1}{2}EI \int_{0}^{l(t)} y_{xx} dx$$

$$= \int_{0}^{l(t)} \left[\frac{1}{2}\rho(v^{2} + v^{2}y_{x}^{2} + y_{t}^{2} + 2vy_{x}y_{t}) - \frac{1}{8}EAy_{x}^{4} - \frac{1}{2}EIy_{xx}^{2}\right] dx.$$
(6)

Then the Hamilton function is:

$$S = \int_{t_1}^{t_2} \int_0^{l(t)} \left[\frac{1}{2} \rho(v^2 + v^2 y_x^2 + y_t^2 + 2vy_x y_t) - \frac{1}{8} EAy_x^4 - \frac{1}{2} EIy_{xx}^2 \right] dxdt.$$
(7)

According to the Hamilton principle, the solution of the above formula is the solution when the Hamilton function is taken to the minimum value, that is, the variation of S is 0. Then the following expression can be acquired:

$$\delta S = 0. \tag{8}$$

The differential equation satisfying the solution is obtained by using Euler equation. From Eq. (7) set:

$$F = \frac{1}{2}\rho(v^2 + v^2y_x^2 + y_t^2 + 2vy_xy_t) - \frac{1}{8}EAy_x^4 - \frac{1}{2}EIy_{xx}^2.$$
(9)

From Euler equation:

$$F_{y} - \frac{\partial}{\partial t}(F_{y_{t}}) - \frac{\partial}{\partial x}(F_{y_{x}}) + \frac{\partial}{\partial x^{2}}(F_{y_{xx}}) = 0.$$
(10)

The partial differential equation can be obtained through calculation of y_x , y_t and y_{xx} , then the Eq. (10) can be converted to the following form:

$$\rho(v^2 y_{xx} + 2v y_{xt} + v_t y_x + y_{tt}) - \frac{3}{2} EA y_x^2 y_{xx} + EI y_{xxxx} = 0.$$
(11)

The boundary conditions are:

$$y(0,t) = e(t), \quad y(l(t),t) = 0.$$
 (12)

Then it constitutes a definite solution problem:

$$\rho(v^2 y_{xx} + 2v y_{xt} + v_t y_x + y_{tt}) - \frac{3}{2} EAy_x^2 y_{xx} + EIy_{xxxx} = 0, \quad 0 < x \le l(t), \quad 0 < t \le T.$$
(13)

The boundary conditions are:

$$y(0,t) = e(t), \quad y(l(t),t) = 0, \quad 0 < t \le T.$$
 (14)

Thus, the establishment of the transverse vibration mathematical model of the steel wire rope in the variable tension hoisting system was completed. By adopting Galerkin discrete method, the partial differential equation problem could also be transformed into an ordinary differential equation problem, and then the problem could be solved. Finally, the transverse vibration of the steel wire rope could be obtained.

3. Solution of mathematical model

For this definite solution problem, as the boundary conditions are inhomogeneous, the inhomogeneous boundary conditions need to be converted into homogeneous boundary conditions first. So y(x, t) should be set to two parts as below:

$$y(x,t) = y_1(x,t) + y_2(x,t).$$
 (15)

And $y_1(x, t)$ satisfies the following conditions:

$$y_1(0,t) = 0, \quad y_1(l(t),t) = 0.$$
 (16)

While $y_2(x, t)$ satisfies the following conditions:

$$y_2(0,t) = e(t), \quad y_2(l(t),t) = 0.$$
 (17)

Then the graph of $y_2(x, t)$ can be set as a line passing through (0, e(t)) and (l(t), 0) in the plane. So $y_2(x, t)$ can be expressed as following:

$$y_2(x,t) = -\frac{e(t)}{l(t)}x + e(t).$$
(18)

Substituting $y(x, t) = y_1(x, t) + y_2(x, t)$ into Eq. (13), the following expression is obtained:

$$\rho(v^2 y_{1xx} + v^2 y_{2xx} + 2v y_{1xt} + 2v y_{2xt} + v_t y_{1x} + v_t y_{2x} + y_{1tt} + y_{2tt}) - \frac{3}{2} EA(y_{1x} + y_{2x})^2 (y_{1xx} + y_{2xx}) + EIy_{1xxxx} + EIy_{2xxxx} = 0.$$
⁽¹⁹⁾

In order to simplify the calculation, the nonlinear equation is linearized to expression below:

$$\rho(v^2 y_{1xx} + 2vy_{1xt} + v_t y_{1x} + y_{1tt}) + EIy_{1xxxx}$$

= -\rho(v^2 y_{2xx} + 2vy_{2xt} + v_t y_{2x} + y_{2tt}) - EIy_{2xxxx}. (20)

And the boundary conditions are:

$$y_1(0,t) = 0, \quad y_1(l(t),t) = 0.$$
 (21)

Therefore, the original problems of definite solution of homogeneous equation of Eq. (13) with inhomogeneous boundary conditions of Eq. (14) can be transformed into problems of definite solution of homogeneous equation of Eq. (20) with inhomogeneous boundary conditions of Eq. (21).

Since the definite solution problem is an infinite dimensional dynamical system problem, the Galerkin discrete method can be used to approximate the original solution $y_1(x, t)$ with finite dimensional solution, and then the infinite dimensional partial differential equation problem can be transformed into a finite dimensional ordinary differential equation problem, so as to solve the definite solution problem. Since the boundary condition is a variable boundary condition, the definite solution problem is a variable domain problem. In order to obtain the solution, it can be set that $x/l(t) = \varepsilon$, so the domain of ε is [0, 1], and the variable domain problem is transformed into a fixed domain problem. So, $y_1(x, t)$ can be assumed as following form:

$$y_1(x,t) = \sum_{i=1}^n \varphi_i(\varepsilon) q_i(t).$$
(22)

Within that:

$$\varphi_i(\varepsilon) = \sqrt{2} \sin(i\pi\varepsilon).$$
 (23)

According to Eq. (18) and Eq. (22), the following expressions can be obtained:

$$y_{1x} = \frac{1}{l(t)} \sum_{i=1}^{n} \varphi_{i}^{"} q_{i}, \quad y_{1xt} = \sum_{i=1}^{n} \left[-\frac{\varepsilon v}{l^{2}(t)} \varphi_{i}^{"} q_{i} + \frac{1}{l(t)} \varphi_{i}^{'} q_{i}^{'} - \frac{v}{l^{2}(t)} \varphi_{i}^{'} q_{i} \right],$$

$$y_{1xx} = \frac{1}{l^{2}(t)} \sum_{i=1}^{n} \varphi_{i}^{"} q_{i}, \quad y_{1xxxx} = \frac{1}{l^{4}(t)} \sum_{i=1}^{n} \varphi_{i}^{""} q_{i}, \quad y_{1t} = \sum_{i=1}^{n} \left[-\frac{\varepsilon v}{l(t)} \varphi_{i}^{'} q_{i} + \varphi_{i} q_{i}^{'} \right],$$

$$y_{1tt} = \sum_{i=1}^{n} \left\{ \frac{\varepsilon^{2} v^{2}}{l^{2}(t)} \varphi_{i}^{"} q_{i} - 2 \frac{\varepsilon v}{l(t)} \varphi_{i}^{'} q_{i}^{'} + \left[\frac{2\varepsilon v^{2}}{l^{2}(t)} - \frac{\varepsilon v_{t}}{l(t)} \right] \varphi_{i}^{'} q_{i} + \varphi_{i} q_{i}^{"} \right\}, \quad y_{2x} = -\frac{e(t)}{l(t)},$$

$$y_{2xt} = -\frac{e(t)v - e'(t)l(t)}{l^{2}(t)}, \quad y_{2xx} = 0, \quad y_{2xxxx} = 0,$$

$$y_{2t} = \frac{e(t)v - e'(t)l(t)}{l^{2}(t)}x + e'(t),$$

$$y_{2tt} = \frac{-e''(t)l^{2}(t) + 2e'(t)l(t)v + e(t)v_{t}l(t) - 2e(t)v^{2}}{l^{3}(t)}x + e''(t).$$
(24)

Substituting Eq. (24) into Eq. (20), so the Eq. (20) can be expressed as following:

$$\sum_{i=1}^{n} \rho \left\{ \frac{(1-\varepsilon)^{2} v^{2}}{l^{2}(t)} \varphi_{i}^{\prime\prime} q_{i} + 2 \frac{(1-\varepsilon)v}{l(t)} \varphi_{i}^{\prime} q_{i}^{\prime} - \left[\frac{2(1-\varepsilon)v^{2}}{l^{2}(t)} - \frac{(1-\varepsilon)v_{t}}{l(t)} \right] \varphi_{i}^{\prime} q_{i} + \varphi_{i} q_{i}^{\prime\prime} \right\} + \frac{EI}{l^{4}(t)} \sum_{i=1}^{n} \varphi_{i}^{\prime\prime\prime\prime} q_{i} = \rho (1-\varepsilon) \left\{ e^{\prime\prime}(t) - \frac{2v}{l(t)} e^{\prime}(t) + \left[\frac{2v^{2}}{l^{2}(t)} - \frac{v_{t}}{l(t)} \right] e(t) \right\}.$$
(25)

Multiply both sides and integrate over the interval of [0, 1], then the Eq. (25) can be written in matrix form as below:

$$MQ'' + CQ' + KQ = F, (26)$$

where:

$$m_{ij} = \rho \int_{0}^{1} \varphi_{i}^{\prime\prime} \varphi_{j} d\varepsilon, \quad c_{ij} = \frac{2\rho v}{l(t)} \int_{0}^{1} (1-\varepsilon) \varphi_{i}^{\prime} \varphi_{j} d\varepsilon,$$

$$k_{ij} = \frac{\rho v^{2}}{l^{2}(t)} \int_{0}^{1} (1-\varepsilon)^{2} \varphi_{i}^{\prime\prime} \varphi_{j} d\varepsilon - \rho \left[\frac{2v^{2}}{l^{2}(t)} - \frac{v_{t}}{l(t)}\right] \int_{0}^{1} (1-\varepsilon) \varphi_{i}^{\prime} \varphi_{j} d\varepsilon$$

$$+ \frac{EI}{l^{4}(t)} \int_{0}^{1} \varphi_{i}^{\prime\prime\prime\prime} \varphi_{j} d\varepsilon,$$

$$f_{j} = \rho \left\{ e^{\prime\prime}(t) - \frac{2v}{l(t)} e^{\prime}(t) + \left[\frac{2v^{2}}{l^{2}(t)} - \frac{v_{t}}{l(t)}\right] e(t) \right\} \int_{0}^{1} (1-\varepsilon) \varphi_{j} d\varepsilon.$$
(27)

The calculation of integration are as following:

$$m_{ij} = \begin{cases} \rho, & i = j, \\ 0, & i \neq j, \end{cases} \quad c_{ij} = \begin{cases} 0, & i = j, \\ -\frac{4\rho v i j}{l(t)(i^2 - j^2)} [1 - 2(-1)^{i-j}], & i \neq j, \end{cases}$$

$$k_{ij} = \begin{cases} \frac{EI}{l^4(t)} i^4 \pi^4, & i = j, \\ \frac{2\rho i j}{l^2(t)(i^2 - j^2)} \left\{ -\frac{4v^2 i^2}{i^2 - j^2} + [2v^2 - l(t)v_t] [1 - 2(-1)^{i-j}] \right\}, & i \neq j. \end{cases}$$
(28)

And then $q_i(t)$ can be obtained through solving the Eq. (26). Then by substituting it into Eq. (22), $y_1(x, t)$ can be aquired, and finally the transverse vibration displacement of the wire rope can be got by adding $y_2(x, t)$ to $y_1(x, t)$.

4. Analysis of results of mathematical model

The vertical length of wire rope of the hoisting system changes as following expressions:

$$0 - t_{1}: l(t) = l - \frac{1}{2}at^{2},$$

$$t_{1} - t_{2}: l(t) = l + \frac{1}{2}at_{1}^{2} - at_{1}t,$$

$$t_{2} - t_{3}: l(t) = l - \frac{1}{2}a(-t^{2} + 2t_{3}t - t_{1}^{2} + t_{1}t_{2} + t_{1}t_{3} - t_{3}^{2}).$$
(29)

The speed of the wire rope changes as is shown in Fig. 2.



Fig. 2. Speed curve of wire rope of hoisting system

Then setting the parameters of hoisting system as Table 1.

Variable	Value
Initial length of wire rope <i>l</i>	20 m
Unit length mass of steel wire rope ρ	0.4 kg
Elastic modulus of wire rope E	1e11 Pa
Section bending resistance coefficient of steel wire rope I	4.9e-8
Time t_1	2 s
Time t_2	8 s
Time t_3	10 s
Acceleration of hosting system a	0.1 m/s^2
Amplitude of sine excitation h	0.002 m
Frequency of sine excitation w	20 rad/s

The transverse vibration displacement of wire rope in hoisting system can be obtained by solving the mechanical model established on computer simulation. Also, the transverse vibration displacement, velocity and acceleration of wire rope of the hoisting system at positions 18 m are analyzed, as are shown in the Fig. 3, Fig. 4 and Fig. 5.

From the observation of Fig. 3, Fig. 4 and Fig. 5, it can be found that when the excitation amplitude of the steel wire rope is 0.002 m and the excitation frequency is 20 Hz, the transverse vibration displacement of the steel wire rope at 18 m is between -0.005 m and 0.005 m, the transverse vibration displacement of the steel wire rope at 18 m is between -0.005 m and 0.005 m, the transverse vibration velocity is between -0.03 m/s and 0.03 m/s, and the transverse vibration acceleration is between -0.75 m/s² and 0.75 m/s². At the same time, the vibration frequency

increases at the later operation process of the hoisting system, which is because of the reduction of the vertical length of the steel wire rope.



Fig. 3. Transverse vibration displacement of wire rope at position 2 m



Fig. 4. Transverse vibration velocity of wire rope at position 2 m



Fig. 5. Transverse vibration acceleration of wire rope at position 2 m

5. Establishment of mechanical model

The transverse vibration analysis of steel wire rope in flexible hoisting system can be described not only by mathematical model, but also by mechanical model directly, and then the transverse vibration displacement can be analyzed. In the process of establishing the mechanical model of the transverse vibration of the steel wire rope of the flexible hoisting system, it is necessary to establish the mechanical model of the steel wire rope. As the steel wire rope is a continuous flexible elastic body, some methods should be adopted in the construction process. The common methods to build the mechanical model of steel wire rope are the beam model and the mass ball model. In the beam model, the two sections of the beam are coupled by bushing forces to form a continuous beam model, as is shown in Fig. 7. While in the mass ball model, the two mass balls are coupled by spring forces to form a discrete mass ball model. The mass ball model is shown in Fig. 6. The two methods can be used separately or in combination, depending on the system requiring being analyzed.





Fig. 7. Beam model

In order to ensure the reliability of the mechanical model of the steel wire rope and prevent the coupling failure of the beam model in the simulation process, a simplified method for building the steel wire rope model is adopted, that is, two sections of the cylinder are connected by hinges, and then the continuous beam is simulated. The simplified model is shown in Fig. 8.



Fig. 8. Simplified model of wire rope

	<u> </u>
Variable	Value
Length of wire rope <i>l</i>	20 m
Radius of wire rope r	0.05 m
Unit length mass of steel wire rope ρ	0.4 kg
Radius of drum R	0.6 m
Length of cylinder a	0.2 m
Number of cylinder segments p	100
Mass of cylinder m	0.08 kg
Weight of hoisting load M	100 kg
Amplitude of sine excitation h	0.002 m
Frequency of sine excitation w	20 rad/s

Table 2. Parameters of flexible hoisting system

In order to simulate the winding process of the steel wire rope on the drum, the steel wire rope should be divided into sections of cylinders so as to wind the drum smoothly. Therefore, in order to simulate the winding process and achieve efficient modeling at the same time, the length of the cylinder should be small. Here, the length of each steel cylinder is set as 0.2 m. At the same time, the total length of the simulated wire rope is 20 m to ensure the efficiency of modeling. Finally, in order to simulate the physical parameters of the actual steel wire rope as real as possible, the mass of the cylinder is consistent with the mass of the steel wire rope of actual length, which is 0.08 kg. And here, the mass of the steel wire rope of unit length is set to be 0.4 kg/m. The specific dynamic hoisting system parameters are shown in Table 2. In addition, the hoisting process of the drum adopts the four section curve motion diagram as shown in Fig. 2. The model of wire rope winding drum and the model of flexible hoisting system are shown in Fig. 10 and Fig. 9.

After the dynamic mechanical model of hoisting system is established, the transverse vibration displacement of wire rope can be analyzed.



Fig. 9. Model of dynamic hoisting system

Fig. 10. Model of wire rope winding drum

6. Analysis of mechanical model

In order to measure the transverse vibration displacement of the steel wire rope, 2 m away from the end of the steel wire rope is selected as the measuring point. The transverse vibration displacement and velocity of the steel wire rope at the position are measured to reflect the transverse vibration conditions of the steel wire rope with axial movement in the hoisting system.

Then the dynamic analysis is carried out on the above flexible hoisting system model, and the transverse vibration displacement and velocity of the steel wire rope at positions 18m are obtained, as are shown in Fig. 11 and Fig. 12.

It can be found from Fig. 11 and Fig. 12, when the excitation amplitude is 0.002 m and the excitation frequency is 20 Hz, the transverse vibration displacement of the steel wire rope is between -0.0045 m and 0.003 m at 18 m, and the transverse vibration velocity is between -0.075 m/s and 0.075 m/s. This is because there is some error including modeling residual, elastic modulus residual and other residuals in simulating the vibration of the steel wire rope in the simulation software, so the amplitude of displacement and velocity have differences with the results of mathematical model.



Fig. 11. Transverse vibration displacement of wire rope at position 18 m

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Fig. 12. Transverse vibration velocity of wire rope at position 18 m

7. Comparison of results between the mathematical model and the mechanical model

By comparing the transverse vibration displacement of the steel wire rope at 18m obtained from the mathematical analysis and the simulated analysis, it can be found that the mathematical analysis results are basically consistent with the simulated analysis results. Because there is a certain residual of modeling and elastic modulus in building the mechanical model of wire rope in the simulation environment, there is a small difference between the two results. However, this difference can be basically ignored to reflect the actual results. The results show that: under the excitation of amplitude of 0.002 m and frequency of 20 Hz at one end of the wire rope, the transverse vibration displacement at 18 m from one end of the wire rope is basically between -0.005 m and 0.005 m, and the transverse vibration velocity at 2 m from one end of the wire rope is basically between -0.04 m/s and 0.04 m/s. Finally, the transverse vibration frequency increases with the decrease of the length of the vertical part of the wire rope. The results are shown in Fig. 13 and Fig. 14.



Fig. 13. Comparison of transverse vibration displacement



Fig. 14. Comparison of transverse vibration velocity

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8. Experimental verification of a hoisting system

In order to further verify the validity of the mathematical model, the experiment is chosen to verify it. As a sine excitation is required, the head sheave is designed as an eccentric part with an eccentricity of 2 mm, as is shown in Fig. 15. The test bench is composed of programmable logic controller, motor, drum, head sheave, wire rope, guide and metal support, as is shown in Fig. 16. The test system consists of acceleration sensor, vibration test system, data processing software and PC, as shown in Fig. 17 and Fig. 18. The test bench parameters are shown in Table 3.

	2
Variable	Value
Length of wire rope <i>l</i>	1.8 m
Radius of wire rope r	0.0015 m
Radius of drum R	0.03 m
Weight of hoisting load M	0.5 kg
Amplitude of sine excitation h	0.002 m

Table 3. Parameters of flexible hoisting system of test bench



Fig. 15. The eccentric head sheave realizing sine excitation



Fig. 17. Acceleration sensor



Fig. 16. Test model of hoisting system



Fig. 18. Testing system and software

By running the hoisting system, the transverse vibration acceleration at 0.2 m away from the end of the wire rope is obtained, as shown in Fig. 19. Acceleration can be found that the wire rope in the starting phase change is very small, this is because when the motor starts, there is a process

to increase the torque of the motor, and during that time the wire rope does not move, so the horizontal vibration acceleration is very small. And it can be found that the transverse vibration acceleration of the wire rope is basically between -0.5 m/s^2 and 0.5 m/s^2 .



Fig. 19. Transverse vibration acceleration of wire rope at position 0.2 m



Fig. 20. Comparison between results of numerical model and test bench

9. Comparison of results between the mathematical model and the experimental model

To compare the results of transverse vibration acceleration of the wire rope between numerical method and experimental method, as is shown in Fig. 20. It can be found that both results are basically between -0.5 m/s^2 and 0.5 m/s^2 . So the validity of the mathematical model is verified.

10. Conclusions

This paper takes the flexible hoisting system as the research object, establishing the mathematical model and mechanical model of the transverse vibration of the steel wire rope of the flexible hoisting system, and analyzes the transverse vibration displacement, velocity and acceleration of the steel wire rope at 2 m away from one end of the rope when the other end of the steel wire rope is excited. Finally, the following four conclusions are drawn as below.

1) When the excitation amplitude at one end of the wire rope is 0.002 m and the excitation frequency is 20 Hz, the transverse vibration displacement at 2 m from one end of the wire rope with length of 20 m is basically between -0.005 m and 0.005 m, the transverse vibration velocity at 2 m from one end of the wire rope is basically between -0.04 m/s and 0.04 m/s, and the transverse vibration acceleration of the wire rope is basically between -0.5 m/s^2 and 0.5 m/s^2 .

2) The transverse vibration frequency of the steel wire rope increases with the decrease of the length of the vertical rope.

3) The mathematical model and mechanical model of the flexible hoisting system were established, and through comparing the data, the numerical results and mechanical results are found basically consistent, indicating the validity of the established mathematical model and mechanical model.

4) The experiment of the flexible hoisting system were established, and through comparing the data, the numerical results and experimental results are found basically consistent, which indicate the validity of the established mathematical model and test bench.

The research and analysis results in this paper have certain reference significance and application value for the modeling, designing and analyzing of the flexible hoisting system.

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