Significance of fin tip temperature on the heat transfer rate and thermal efficiency of a convective-radiative rectangular fin with variable thermal conductivity

M. G. Sobamowo¹, O. A. Adeleye², A. A. Yinusa³, M. O. Oyekeye⁴, M. A. Waheed⁵

^{1,3,4}Department of Mechanical Engineering, University of Lagos, Akoka, Lagos State, Nigeria
 ²Department of Biomedical Engineering, University of Lagos, Akoka, Lagos State, Nigeria
 ⁵Department of Mechanical Engineering, Federal University of Agriculture, Abeokuta, Ogun State, Nigeria
 ¹Corresponding author

E-mail: ¹*mikegbeminiyiprof@yahoo.com*, ²*rotimiadeleye1711@gmail.com*, ³*mynotebook2010@yahoo.com*, ⁴*moyekeye@unilag.edu.ng*, ⁵*akindoye@gmail.com*

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Abstract. In this work, effect of fin tip temperature on the rate of heat transfer and thermal efficiency of a rectangular convective-radiative fin with temperature-dependent thermal conductivity is analyzed using differential transformation method. The results of the power series solutions are verified numerically, and very good agreements are established. Also, the symbolic solutions are used to examine the effects of the conductive-convective and nonlinear thermal conductivity parameters on the thermal performance of the passive device. It is found that when the nonlinear thermal conductivity parameter increases, the fin tip temperature increases. However, the temperature at the tip of the fin decreases as the conductive-convective parameter increases. The thermal efficiency of the fin increases as the fin tip temperature and nonlinear thermal conductivity parameters are augmented but an increase conductive-convective parameter causes the fin tip temperature and the thermal efficiency of the extended surface to reduce. An increase in the conductive-convective parameter causes decrease the temperature distribution and thermal efficiency in the passive device. However, the efficiency of the fin increases as the nonlinear thermal conductivity parameter increases. When nonlinear thermal conductivity and conductive-convective parameters increase, the rate of heat transfer at the fin base increases. The developed analytical solutions provide a good platform for the nonlinear thermal analysis of the fin and proper design of the extended surfaces in thermal systems.

Keywords: analytical solution, rectangular fin, fin tip temperature, differential transform method.

Nomenclature

- A_{cr} Fin cross sectional area, m²
- B_o Magnetic field intensity, Tesla or kg/sec²Amp
- c_{pa} Specific heat capacity, J/kgK
- *h* Coefficient of convective heat transfer, W/m^2K
- J_c Conduction current intensity, A
- *k* Fin thermal conductivity, W/mK
- k_b Fin thermal conductivity at the base temperature, W/mK
- *L* Fin length, M
- *M_c* Adimensional convective parameter
- *N_r* Adimensional radiation parameter
- *P* Fin perimeter, m
- t Time, sec
- T Fin temperature, K
- T_{∞} Ambient temperature, K
- T_b Fin temperature at the base, K

- *x* Fin axial distance, m
- *X* Adimensional fin length
- δ Fin thickness, m
- θ Adimensional temperature
- θ_b Adimensional temperature at the fin base
- ρ Fin material density, kg/m³
- σ Stefan-Boltzmann constant, W/m²K⁴
- σ Electrical conductivity, $\Omega^{-1}m^{-1}$ or sec²Amp²/kgm³

1. Introduction

The enhancements of heat transfer by the use of fins in thermal and electronic systems serve the prime purpose of effective cooling of the engineering and industrial equipment [1]. Such an important applications has provoked so many research works [2-18], just to mention a few. In the bid of theoretical investigations, the thermal damage problems and heat transfer enhancement by the extended surfaces, the controlling thermal models of the passive devices are always nonlinear. Although, there are various approximate analytical and numerical solutions that have been used to solve the thermal problems [6-26], the quest for further studies is increasing demanded. In a recent work, Darvishi et al. [27] studied steady state thermal performance in convective-radiative porous radial fin while in the same year, Hoshyar [28] adopted homotopy perturbation method to developed series solution to steady state thermal performance longitudinal porous fins with variable internal heat generation. In the following year, Sobamowo [29] applied Galerkin's method of weighted residual to some simple but highly accurate analytical solutions for thermal performance of rectangular fin with variable internal heat generation and thermal conductivity. Also, with the use of various approximate analytical methods, the effects of magnetic field on the steady state thermal performance of solid and porous fins were explored by Hashar et al. [30], Oguntala et al. [31], Patel and Meher [32]. Additionally, Sobamowo [33] examined the optimum design and thermal performance of a cooling fin under convective and radiative heat transfer using finite element method.

The review of the past studies shows that the analytical study of nonlinear transient heat transfer analysis in extended surfaces have not extensively been presented in literature. Moreover, the obvious advantages of generating analytical solutions to the nonlinear problems are very much important. Such solutions provide proper physical insights and effective predictions of the thermal performance of the extended surface. Also, the analytical approach of differential transformation method reduces the complex mathematical analysis, high computational cost and time. Additional to the best of the authors' knowledge, the exploration of the effect of fin tip temperature on the rate of heat transfer and thermal efficiency of a fin has not been presented in literature. Therefore, the present work applies differential transformation method to examine the effect of temperature at the tip of the fin on the rate of heat transfer and thermal efficiency of a rectangular convective fin with temperature-dependent thermal conductivity. Also, the symbolic solutions are used to examine the effects of the conductive-convective and nonlinear thermal conductivity parameters on the thermal performance of the passive device.

2. Problem formulation

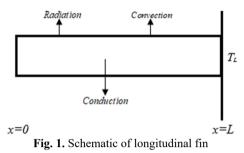
Consider a longitudinal straight fin with variable thermal conductivity and exposed on both faces to a convective-radiative environment at temperature T_{∞} and heat transfer co-efficient *h* as in Fig. 1. Assuming that the extended surface isotropic, homogeneous and with constant thermo-physical properties. It is taken that the heat transfer is one-dimensional along the length of the fin. The prime surface is perfectly in thermal contact with the base of the fin and there is no heat gain or loss through the tip of the fin.

Thermal energy equation based on model assumptions is expressed as:

$$\frac{d}{dx}\left(k(T)\frac{dT}{dx}\right) - \frac{hP}{A_{cr}}(T - T_{\infty}) - \frac{\sigma\varepsilon P}{A_{cr}}(T^4 - T_{\infty}^4) = 0,\tag{1}$$

where the temperature-dependent thermal conductivity is expressed as a linear law:

$$k(T) = k_b \left(1 + \gamma (T - T_{\infty}) \right). \tag{2}$$



Therefore, Eq. (1) can be written as:

$$\frac{d}{dx}\left(k_b\left(1+\gamma(T-T_{\infty})\right)\frac{dT}{dx}\right) - \frac{hP}{A_{cr}}(T-T_{\infty}) - \frac{\sigma\varepsilon P}{A_{cr}}(T^4-T_{\infty}^4) = 0.$$
(3)

The boundary conditions are:

$$x = 0, \quad \frac{\partial T}{\partial x} = 0, \tag{4}$$

$$x = L, \quad \tilde{T} = T_L. \tag{5}$$

When negligible temperature difference exists between fin base and fin tip, the term T^4 in Eq. (3) could be expressed as a linear function of temperature as:

$$T^{4} = T_{\infty}^{4} + 4T_{\infty}^{3}(T - T_{\infty}) + 6T_{\infty}^{2}(T - T_{\infty})^{2} + \dots \cong 4T_{\infty}^{3}T - 3T_{\infty}^{4}.$$
 (6)

On substituting Eq. (6) into Eq. (3), we have:

$$\frac{d}{dx}\left(\left(1+\gamma(T-T_{\infty})\right)\frac{dT}{dx}\right) - \frac{hP(T-T_{\infty})}{k_bA_{cr}} - \frac{4\sigma\varepsilon PT_{\infty}^3(T-T_{\infty})}{k_bA_{cr}} = 0.$$
(7)

On introducing the following dimensionless parameters in Eq. (8) into Eq. (7):

$$X = \frac{x}{L}, \quad \theta = \frac{T - T_{\infty}}{T_L - T_{\infty}}, \quad N^2 = \frac{PhL^2}{A_{cr}k_b}, \quad Nr = \frac{4\sigma\varepsilon PT_{\infty}^3 L^3}{A_{cr}k_b}, \quad \beta = \gamma(T_L - T_{\infty}), \tag{8}$$

we arrived at the dimensionless forms of the governing as follows:

$$\frac{d}{dX}\left[(1+\beta\theta)\frac{d\theta}{dX}\right] - (N^2 + Nr)\theta = 0.$$
(9)

On expanding Eq. (9), we have:

$$\frac{\partial^2 \theta}{\partial X^2} + \beta \theta \frac{\partial^2 \theta}{\partial X^2} + \beta \left(\frac{\partial \theta}{\partial X}\right)^2 - (N^2 + Nr)\theta = 0, \tag{10}$$

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and the dimensionless boundary conditions are:

$$X = 0, \quad \frac{d\theta}{dX} = 0, \tag{11}$$
$$X = 1, \quad \theta = 1. \tag{12}$$

3. Method of solution: differential transform method

The above governing differential equation is highly nonlinear, and such nonlinearity imposes some difficulties in the development of exact analytical method to generate closed form solution for the equation. Therefore, in this work, differential transformation method is used. The differential transformation method maps a governing equation into an algebraic domain and then obtain an inversion using a series summation method. This approximate analytical method generates solution with the controlling parameters adequately conserved.

The basic definitions of the method can be found in Moradi et al. [19]. However, the differential transformations of some functions are presented as follow:

1. If
$$z(x) = u(x) \pm v(x)$$
, then $Z(k) = U(k) \pm V(k)$ for all $k \ge 0$.
2. If $z(x) = \alpha u(x)$, then $Z(k) = \alpha U(k)$.
3. If $z(x) = \frac{d^{u}}{dx}$, then $Z(k) = (k+1)U(k+1)$.
4. If $z(x) = \frac{d^{2u}}{dx^{2}}$, then $Z(k) = (k+1)(k+2)U(k+2)$.
5. $z(x) = \frac{d^{nu}}{dx^{n}}$, then $Z(k) = (k+1)(k+2)(k+3)...(k+n)U(k+n)$.
6. If $z(x) = u(x)v(x)$, then $Z(k) = \sum_{l=0}^{k} U[l]V[k-l]$.
7. If $z(x) = u^{2}(x)$, then $Z(k) = \sum_{l=0}^{k} U[l]V[k-l]$.
8. If $z(x) = u(x)v(x)w(x)$, then $Z(k) = \sum_{l=0}^{k} \sum_{p=0}^{k-1} U[l]V[p]W[k-l-p]$.
9. If $z(x) = u^{3}(x)$, then $Z(k) = \sum_{l=0}^{k} U^{n-1}[l]U[k-l]$.
10. If $z(x) = u^{m}(x)$, then $Z(k) = \sum_{l=0}^{k} U^{n-1}[l]U[k-l]$.
11. $z(x) = u(x)\frac{du(x)}{dx}$, then $Z(k) = \sum_{l=0}^{k} U^{k-l}U[l]U[p][k-l-p+1]U[k-l-p+1]$.
13. If $z(x) = \left[\frac{du(x)}{dx}\right]^{2}$, then $Z(k) = \sum_{l=0}^{k} (l+1)U[l+1](k-l+1)U[k-l+1]$.
14. If $z(x) = \frac{du(x)dw(x)}{dx}$, then $Z(k) = \sum_{l=0}^{k} \delta[l-m](k-l+1)U[k-l+1]$.
15. $z(x) = x^{m}\frac{d^{2}u(x)}{dx^{2}}$, then $Z(k) = \sum_{l=0}^{k} \delta[l-m](k-l+2)(k-l+1)U(k-l+2)$.
16. $z(x) = x^{m}\frac{d^{2}u(x)}{dx^{2}}$, then $Z(k) = \sum_{l=0}^{k} \delta[l-m](k-l+n)(k-l+n-1)...(k-l+1)U(k-l+n)$.
17. $z(x) = x^{m}$, then $Z(k) = \delta(k-m) = \begin{bmatrix} 1, & k=m \\ 0, & otherwise \end{bmatrix}$.

4. Solution of the nonlinear differential transformation method

The differential transform of the dimensionless governing of equation in Eq. (10) is:

$$(k+1)(k+2)\theta(k+2) + \beta \sum_{l=0}^{k} \theta(l)(k-l+1)\theta(k-l+1) + \beta \sum_{l=0}^{k} (l+1)\theta[l+1](k-l+1)\theta[k-l+1] - (N^2 + Nr)\theta[k] = 0.$$
(13)

Which can be simplified further as:

$$(k+1)(k+2)\theta(k+2) + \beta \sum_{l=0}^{k} \theta(l)(k-l+1)\theta(k-l+1) + \beta \sum_{l=0}^{k} (l+1)\theta[l+1](k-l+1)\theta[k-l+1] - (N^2 + Nr)\theta[k] = 0 = 0,$$
(14)

where:

$$\delta(k) = \begin{cases} 1, & k = 0, \\ 0, & otherwise. \end{cases}$$
(15)

From the boundary conditions, one can write differential transforms:

$$\begin{aligned} \theta[0] &= \lambda, \\ \theta[1] &= 0. \end{aligned}$$

From the differential transform of the governing equation, we have:

$$\begin{split} \theta[2] &= \frac{(N^2 + Nr)\lambda}{2(1 + \beta\lambda)}, \\ \theta[3] &= 0, \\ \theta[4] &= \frac{(N^2 + Nr)^2\lambda(1 - 2\beta\lambda)}{24(1 + \beta\lambda)^3}, \\ \theta[5] &= 0, \\ \theta[6] &= \frac{(N^2 + Nr)^3\lambda(1 - 16\beta\lambda + 28\beta^2\lambda^2)}{720(1 + \beta\lambda)^5}, \\ \theta[7] &= 0, \\ \theta[8] &= \frac{(N^2 + Nr)^4\lambda(1 - 7\beta\lambda + 60\beta^2\lambda^2 - 896\beta^3\lambda^3)}{40320(1 + \beta\lambda)^7}, \\ \theta[9] &= 0, \\ \theta[10] &= \frac{(N^2 + Nr)^5\lambda(1 - 332\beta\lambda + 7812\beta^2\lambda^2 - 39896\beta^3\lambda^3 + 51184\beta^4\lambda^4)}{3628800(1 + \beta\lambda)^9}, \\ \theta[11] &= 0. \end{split}$$

From the definition of DTM, we have the solution of:

$$\theta[X] = \lambda + \left[\frac{(Nc^{2} + Nr)\lambda}{2(1 + \beta\lambda)}\right] X^{2} + \left[\frac{(N^{2} + Nr)^{2}\lambda(1 - 16\beta\lambda + 28\beta^{2}\lambda^{2})}{720(1 + \beta\lambda)^{5}}\right] X^{4} \\ + \left[\frac{(N^{2} + Nr)^{3}\lambda(1 - 16\beta\lambda + 28\beta^{2}\lambda^{2})}{720(1 + \beta\lambda)^{5}}\right] X^{6} \\ + \left[\frac{(N^{2} + Nr)^{4}\lambda(1 - 7\beta\lambda + 60\beta^{2}\lambda^{2} - 896\beta^{3}\lambda^{3})}{40320(1 + \beta\lambda)^{7}}\right] X^{8} \\ + \left[\frac{(N^{2} + Nr)^{5}\lambda(1 - 332\beta\lambda + 7812\beta^{2}\lambda^{2} - 39896\beta^{3}\lambda^{3} + 51184\beta^{4}\lambda^{4})}{3628800(1 + \beta\lambda)^{9}}\right] X^{10} + \dots,$$
(16)

where, the constant λ is determined from the terminal boundary conditions as:

$$\theta[1] = \lambda + \left[\frac{(N^2 + Nr)\lambda}{2(1 + \beta\lambda)}\right] + \left[\frac{(N^2 + Nr)^2\lambda(1 - 16\beta\lambda + 28\beta^2\lambda^2)}{720(1 + \beta\lambda)^5}\right] \\ + \left[\frac{(N^2 + Nr)^3\lambda(1 - 16\beta\lambda + 28\beta^2\lambda^2)}{720(1 + \beta\lambda)^5}\right] \\ + \left[\frac{(N^2 + Nr)^4\lambda(1 - 7\beta\lambda + 60\beta^2\lambda^2 - 896\beta^3\lambda^3)}{40320(1 + \beta\lambda)^7}\right] \\ + \left[\frac{(N^2 + Nr)^5\lambda(1 - 332\beta\lambda + 7812\beta^2\lambda^2 - 39896\beta^3\lambda^3 + 51184\beta^4\lambda^4)}{3628800(1 + \beta\lambda)^9}\right] + \dots + \dots = 1.$$
(17)

5. Heat transfer rate in the fin

The rate of heat transfer through the fin is given as:

$$q = k(T)A_{cr}\frac{dT}{dx}.$$
(18)

The dimensionless heat transfer rate is given as:

$$\bar{q} = \frac{qL}{kA_{cr}(T_L - T_{\infty})} = (1 + \beta\theta)\frac{d\theta}{dX}.$$
(19)

Substituting Eq. (16) into Eq. (19), we have:

$$\begin{split} \bar{q}_{b} &= \frac{qL}{A_{cr}k(T_{L} - T_{\infty})} \\ &= \left\{ \left(1 + \beta \left[\lambda + \left[\frac{(N^{2} + Nr)\lambda}{2(1 + \beta\lambda)} \right] X^{2} + \left[\frac{(N^{2} + Nr)^{2}\lambda(1 - 16\beta\lambda + 28\beta^{2}\lambda^{2})}{720(1 + \beta\lambda)^{5}} \right] X^{4} \\ &+ \left[\frac{(N^{2} + Nr)^{3}\lambda(1 - 16\beta\lambda + 28\beta^{2}\lambda^{2})}{720(1 + \beta\lambda)^{5}} \right] X^{6} \\ &+ \left[\frac{(N^{2} + Nr)^{4}\lambda(1 - 7\beta\lambda + 60\beta^{2}\lambda^{2} - 896\beta^{3}\lambda^{3})}{40320(1 + \beta\lambda)^{7}} \right] X^{8} \\ &+ \left[\frac{(N^{2} + Nr)^{5}\lambda(1 - 332\beta\lambda + 7812\beta^{2}\lambda^{2} - 39896\beta^{3}\lambda^{3} + 51184\beta^{4}\lambda^{4})}{3628800(1 + \beta\lambda)^{9}} \right] X^{10} + \ldots \right] \right) \quad (20) \\ &\times \left[\frac{\left[\frac{(N^{2} + Nr)\lambda}{(1 + \beta\lambda)} \right] X + \left[\frac{(N^{2} + Nr)^{2}\lambda(1 - 16\beta\lambda + 28\beta^{2}\lambda^{2})}{180(1 + \beta\lambda)^{5}} \right] X^{3} \\ &+ \left[\frac{(N^{2} + Nr)^{3}\lambda(1 - 16\beta\lambda + 28\beta^{2}\lambda^{2})}{120(1 + \beta\lambda)^{5}} \right] X^{5} \\ &+ \left[\frac{(N^{2} + Nr)^{4}\lambda(1 - 7\beta\lambda + 60\beta^{2}\lambda^{2} - 896\beta^{3}\lambda^{3})}{5040(1 + \beta\lambda)^{7}} \right] X^{7} \\ &+ \left[\frac{(N^{2} + Nr)^{4}\lambda(1 - 7\beta\lambda + 60\beta^{2}\lambda^{2} - 896\beta^{3}\lambda^{3})}{362880(1 + \beta\lambda)^{9}} \right] X^{9} + \ldots \right] \right\}. \end{split}$$

The rate of heat transfer at the base of the fin X = 1:

$$\begin{split} \bar{q}_{b}|_{X=1} &= \left\{ \left(1 + \beta \theta \right) \frac{d\theta}{dX} \Big|_{X=1} \\ &= \left\{ \left(1 + \beta \left[\lambda + \left[\frac{(N^{2} + Nr)\lambda}{2(1 + \beta\lambda)} \right] + \left[\frac{(N^{2} + Nr)^{2}\lambda(1 - 16\beta\lambda + 28\beta^{2}\lambda^{2})}{720(1 + \beta\lambda)^{5}} \right] \right. \\ &+ \left[\frac{(N^{2} + Nr)^{3}\lambda(1 - 16\beta\lambda + 28\beta^{2}\lambda^{2})}{720(1 + \beta\lambda)^{5}} \right] \\ &+ \left[\frac{(N^{2} + Nr)^{4}\lambda(1 - 7\beta\lambda + 60\beta^{2}\lambda^{2} - 896\beta^{3}\lambda^{3})}{40320(1 + \beta\lambda)^{7}} \right] \\ &+ \left[\frac{(N^{2} + Nr)^{5}\lambda(1 - 332\beta\lambda + 7812\beta^{2}\lambda^{2} - 39896\beta^{3}\lambda^{3} + 51184\beta^{4}\lambda^{4})}{3628800(1 + \beta\lambda)^{9}} \right] + \dots \right] \right) \end{split}$$
(21)
$$\times \left[\frac{\left[\frac{(N^{2} + Nr)\lambda}{(1 + \beta\lambda)} \right] + \left[\frac{(N^{2} + Nr)^{2}\lambda(1 - 16\beta\lambda + 28\beta^{2}\lambda^{2})}{180(1 + \beta\lambda)^{5}} \right] \\ &+ \left[\frac{(N^{2} + Nr)^{3}\lambda(1 - 16\beta\lambda + 28\beta^{2}\lambda^{2})}{120(1 + \beta\lambda)^{5}} \right] \\ &+ \left[\frac{(N^{2} + Nr)^{4}\lambda(1 - 7\beta\lambda + 60\beta^{2}\lambda^{2} - 896\beta^{3}\lambda^{3})}{5040(1 + \beta\lambda)^{7}} \right] \\ &+ \left[\frac{(N^{2} + Nr)^{4}\lambda(1 - 7\beta\lambda + 60\beta^{2}\lambda^{2} - 896\beta^{3}\lambda^{3})}{5040(1 + \beta\lambda)^{7}} \right] \\ &+ \left[\frac{(N^{2} + Nr)^{5}\lambda(1 - 332\beta\lambda + 7812\beta^{2}\lambda^{2} - 39896\beta^{3}\lambda^{3} + 51184\beta^{4}\lambda^{4})}{362880(1 + \beta\lambda)^{9}} \right] + \dots \right] \right\}. \end{split}$$

6. Thermal efficiency of the fin

The quantity of heat dissipated form the fin can be written as:

$$Q_f = \int_0^1 [Ph(T - T_\infty) + \sigma P\varepsilon (T^4 - T_\infty^4)] dx.$$
⁽²²⁾

The maximum heat dissipated form the fin can be expressed as:

$$Q_{\max} = PhL(T_L - T_{\infty}) + +\sigma PL\varepsilon(T_L^4 - T_{\infty}^4).$$
⁽²³⁾

Therefore, the thermal efficiency of the fin is given as:

$$\eta = \frac{Q_f}{Q_{\text{max}}} = \frac{\int_0^L (Ph(T - T_\infty) + \sigma P\varepsilon(T^4 - T_\infty^4))dx}{PhL(T_b - T_\infty) + \sigma PL\varepsilon(T_L^4 - T_\infty^4)}.$$
(24)

After applying Eqs. (6) and (8) to Eq. (24), the thermal efficiency of the fin can be expressed in dimensionless form as:

$$\eta = \int_0^1 \theta d X. \tag{25}$$

On substituting Eq. (16) into Eq. (25), after integrating the resulting equation, we have:

$$\eta = \lambda + \frac{1}{3} \left[\frac{(Nc^2 + Nr)\lambda}{2(1 + \beta\lambda)} \right] + \frac{1}{5} \left[\frac{(Nc^2 + Nr)^2\lambda(1 - 16\beta\lambda + 28\beta^2\lambda^2)}{720(1 + \beta\lambda)^5} \right] + \frac{1}{7} \left[\frac{(Nc^2 + Nr)^3\lambda(1 - 16\beta\lambda + 28\beta^2\lambda^2)}{720(1 + \beta\lambda)^5} \right] + \frac{1}{9} \left[\frac{(Nc^2 + Nr)^4\lambda(1 - 7\beta\lambda + 60\beta^2\lambda^2 - 896\beta^3\lambda^3)}{40320(1 + \beta\lambda)^7} \right] + \frac{1}{11} \left[\frac{(Nc^2 + Nr)^5\lambda(1 - 332\beta\lambda + 7812\beta^2\lambda^2 - 39896\beta^3\lambda^3 + 51184\beta^4\lambda^4)}{3628800(1 + \beta\lambda)^9} \right] + \dots$$
(26)

In order to verify the results of the present work, the nonlinear model was also solved numerically using finite difference method. The finite difference scheme for the nonlinear thermal model in Eq. (10) is:

$$\left(\frac{\theta_{i+1} - 2\theta_i + \theta_{i-1}}{2\Delta^2 X}\right) + \beta(\theta_i) \left(\frac{\theta_{i+1} - 2\theta_i + \theta_{i-1}}{2\Delta^2 X}\right) + \beta\left(\frac{\theta_{i+1} - \theta_{i-1}}{2\Delta X}\right)^2 - (N^2 + Nr)\theta_i = 0.$$
(27)

While for the boundary conditions are:

$$\frac{\theta_1 - \theta_{-1}}{2\Delta X} = 0 \quad \Rightarrow \quad \theta_1 = \theta_{-1},$$

$$\theta_M = 1.$$
(28)

From Eq. (19), the finite difference of the rate of heat transfer from the fin can be written as:

$$q_b = (1 + \beta \theta_i) \left(\frac{\theta_{i+1} - \theta_i}{\Delta X} \right).$$
⁽²⁹⁾

From Eq. (25), one can write the finite difference of the thermal efficiency as:

$$\eta = \sum_{i=1}^{N} \theta_i.$$
(30)

7. Results and discussion

The developed power series solutions which are presented in the previous section are simulated with the aid of MATHEMATICA. Also, the symbolic solutions are used to examine the effects of the conductive-convective and nonlinear thermal conductivity parameters on the thermal performance of the passive device as shown in Figs. 2-6. However, Table 1 shows the comparison of the results of DTM and FDM. It shows that DTM provides good agreements with the results of the numerical method.

Table 1. Comparison of results					
X	FDM	DTM	FDM-DTM		
0.0000	0.867912	0.867912	0.000000		
0.2000	0.872763	0.872762	0.000001		
0.4000	0.887512	0.887511	0.000001		
0.6000	0.912773	0.912772	0.000001		
0.8000	0.949668	0.949668	0.000000		
1.0000	1.000000	1.000000	0.000000		

 Table 1. Comparison of results

The effects of conductive-convective and nonlinear thermal conductivity parameters on the tip temperature of the fin are shown in Table 2. It is shown that when the nonlinear thermal conductivity parameter increases, the fin tip temperature also increases. However, the fin tip temperature decreases as the conductive-convective parameter increases.

	barameters on the dimensionless tip					
Ν	β	λ (Tip temperature)				
1	-0.2	0.60417				
1	0.0	0.64805				
1	0.4	0.71605				
1	0.8	0.76438				
1	1.2	0.79964				
1	1	0.78333				
1	2	0.84673				
2	1 2 -0.2	0.23687				
2	0.0	0.26580				
2	0.4	0.32443				
2	0.8	0.38078				
2	1.2	0.43268				
2	1 2 -0.2	0.40743				
2	2	0.51961				
3	-0.2	0.08812				
3	0.0	0.09934				
$ \begin{array}{c} 2\\ 2\\ 2\\ 2\\ 2\\ 2\\ 3\\ 3\\ 3\\ 3\\ 3\\ 3\\ 4\\ 4 \end{array} $	0.4	0.12613				
3	0.8	0.15387				
3	1.2	0.18456				
3	1.2 1	0.16840				
3	2 -0.2	0.25000				
4		0.03301				
4	0.0	0.03667				
4	0.4	0.04723				
4	0.8	0.06123				
4	1.2	0.07136				
4	1	0.06702				
4	2	0.08987				

Table 2. Effects of nonlinear thermal conductivity and convective parameters on the dimensionless tip temperature

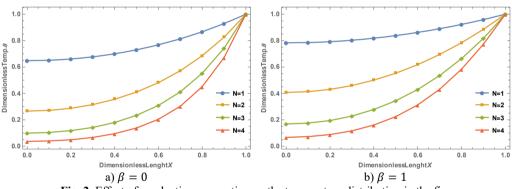
Tables 3 depict the impacts of fin tip temperature, conductive-convective and nonlinear thermal conductivity parameters on the thermal efficiency of the rectangular fin. The table shows that the thermal efficiency of the fin increases as the fin tip temperature and nonlinear thermal conductivity parameters are augmented. However, an increase conductive-convective parameter causes the fin tip temperature and the thermal efficiency of the fin to reduce.

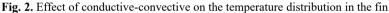
on	on the dimensionless tip temperature and efficiency for $\beta = -0.2$ to 0.2					
Ν	$\beta = -0.2$		$\beta = 0.0$		$\beta = 0.2$	
	α	η	α	η	α	η
1	0.60417	0.544452	0.64805	0.590776	0.68501	0.629962
2	0.23687	0.249533	0.26580	0.276327	0.29528	0.302383
3	0.08812	0.156813	0.09934	0.170802	0.11227	0.186138
4	0.03306	0.118420	0.03667	0.125771	0.04133	0.135573
5	0.01254	0.096811	0.01358	0.100752	0.01492	0.106018
6	0.00483	0.083083	0.00510	0.085161	0.00542	0.087526
7	0.00190	0.073960	0.00196	0.074685	0.00203	0.075718
8	0.00077	0.067905	0.00078	0.067624	0.00070	0.052991
9	0.00032	0.063126	0.00032	0.063083	0.00033	0.064163
10	0.00014	0.059395	0.00014	0.058893	0.00014	0.060066

 Table 3. Effects of nonlinear thermal conductivity and convective parameters

on t	on the dimensionless tip temperature and efficiency for $\beta = 0$ to 0.8					
Ν	$\beta = 0.0$		$\beta = 0.4$		$\beta = 0.8$	
	α	η	α	η	α	η
1.0	0.64805	0.59078	0.71605	0.66318	0.76438	0.71579
1.5	0.42510	0.39208	0.49878	0.45963	0.56102	0.51685
2.0	0.26580	0.27633	0.32443	0.32710	0.38078	0.37346
2.5	0.16308	0.21064	0.20379	0.24852	0.24578	0.28284
3.0	0.09934	0.17080	0.12613	0.20113	0.15387	0.22605
3.5	0.06037	0.14451	0.07756	0.17036	0.09617	0.19184
4.0	0.03667	0.12577	0.04723	0.14770	0.06123	0.17233
4.5	0.02230	0.11177	0.02831	0.12904	0.03904	0.15866
5.0	0.01358	0.10075	0.01672	0.11318	0.02351	0.14096
5.5	0.00830	0.09206	0.00983	0.10033	0.01295	0.11797
6.0	0.00510	0.08516	0.00581	0.09030	0.00704	0.09962
6.5	0.00315	0.07941	0.00347	0.08245	0.00396	0.08767
7.0	0.00196	0.07469	0.00211	0.07693	0.00230	0.07933
7.5	0.00123	0.07077	0.00130	0.07244	0.00138	0.07387
8.0	0.00078	0.06762	0.00081	0.06843	0.00084	0.06865
8.5	0.00050	0.06505	0.00051	0.06465	0.00053	0.06640
9.0	0.00032	0.06308	0.00033	0.06310	0.00033	0.06101
9.5	0.00021	0.05973	0.00022	0.06397	0.00022	0.06240
10.0	0.00014	0.05889	0.00014	0.05790	0.00014	0.05691

Table 4. Effects of nonlinear thermal conductivity and convective parameters on the dimensionless tip temperature and efficiency for $\beta = 0$ to 0.8





The effects of conductive-convective parameter on the temperature distribution in the fin is shown in Figs. 2(a) and 2(b). It is shown that as the conductive-convective parameter increases, the fin temperature distribution in the passive device decreases which mean that the heat transfer rate through the fin is augmented and hence, the fin thermal efficiency is improved. This is because, as the convective heat transfer increases, there is more effective convective cooling which consequently lowers the temperatures in the fin. By extension, it could be stated that increasing in the values of conductive-convective parameter produces a decrease in the value of the fin efficiency. However, the efficiency of the fin increases as the nonlinear thermal conductivity parameter increases. It should be noted that the profile has steepest temperature gradient at lower value of the conduction-convection term, but its much higher value gotten from the lower value of thermal conductivity than the other values of convective parameter in the profiles produces a lower heat-transfer rate. This shows that the thermal performance or efficiency of the fin is favoured at low values of convective parameter since the aim (high effective use of the fin) is to minimize the temperature decrease along the fin length, where the best possible scenario is when $T = T_h$ everywhere. It must be pointed out that a small value of N_c correspond to a relatively short and thick fins of poor thermal conductivity and high value of N_c implies a long fin or fin with low value of thermal conductivity. Since, the thermal performance or efficiency of the fin is favoured at low values of convective fin parameter, very long fins are to be avoided in practice.

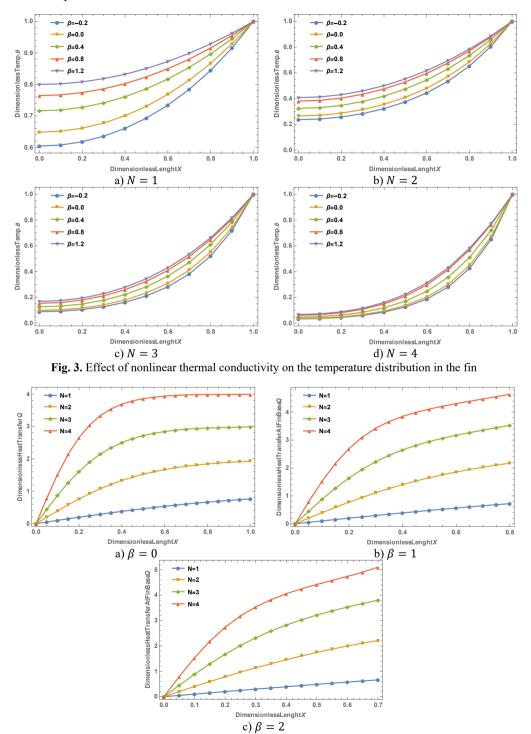


Fig. 4. Effect of conductive-convective on the rate of heat transfer in the fin

Figs. 3(a)-3(d) present the influence of the effect of nonlinear thermal conductivity paramter on temperature distributions in the fin. It should be noted that the positive value of the nonlinear thermal conductivity parameter implies fin materials whose thermal conductivity increases with increase in temperature while the negative value of the nonlinear thermal conductivity parameter corresponds to fin materials whose thermal conductivity decreases with increase in temperature. It can be inferred that the heat transfer rate through the fin increases as the nonlinear thermal conductivity parameter increases. This effect is as a result of the fact that the fin is able to transfer more heat due to increased thermal conductivity of the fin as the nonlinear thermal conductivity parameter increases.

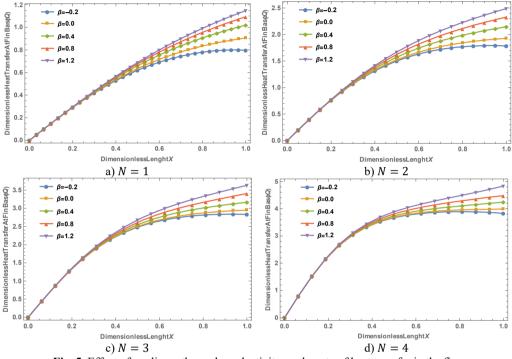


Fig. 5. Effect of nonlinear thermal conductivity on the rate of heat transfer in the fin

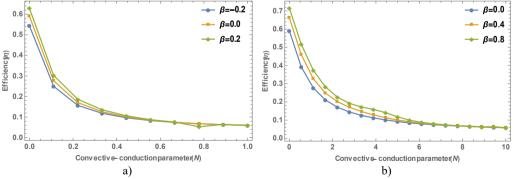


Fig. 6. Effect of nonlinear thermal conductivity on the efficiency in the fin

The influence of the nonlinear thermal conductivity and conductive-convective parameters on the rate of heat transfer at the base of the rectangular fin is shown in Figs. 4(a)-4(c) and 5(a)-5(d). The figure presents that the as nonlinear thermal conductivity and conductive-convective parameters increase, the rate of heat transfer at the fin base increases. This is due to the fact when

the nonlinear thermal conductivity and conductive-convective parameters increase, the ability of the fin material to conduct heat to the fin base increases. When the nonlinear thermal conductivity parameter increases, the temperature-dependent thermal conductivity of the fin material becomes more sensitive to the fin temperature which causes an increase in the heat transfer rate in the fin.

Fig. 6(a)-6(b) show the impact of the nonlinear thermal conductivity and conductive-convective parameters on the thermal efficiency of the longitudinal fin. From the figure, it is depicted that when the nonlinear thermal conductivity is increased, the thermal efficiency of the fin increases but an increase in the conductive-convective parameters causes the thermal efficiency of the fin to decrease.

8. Conclusions

In this work, effect of temperature at the tip of the fin on the rate of heat transfer and thermal efficiency of a rectangular convective fin with temperature-dependent thermal conductivity is analyzed using differential transformation method. The results of the approximate analytical solutions are verified numerically, and very good agreements are established. Also, the symbolic solutions are used for further parametric studies and the following results were established:

1) When the nonlinear thermal conductivity parameter increases, the fin tip temperature increases. However, the fin tip temperature decreases as the conductive-convective parameter increases.

2) The thermal efficiency of the fin increases as the fin tip temperature and nonlinear thermal conductivity parameters are augmented but an increase conductive-convective parameter causes the fin tip temperature and the thermal efficiency of the fin to reduce.

3) An increase in the conductive-convective parameter causes decrease the temperature distribution and thermal efficiency in the passive device. However, the efficiency of the fin increases as the nonlinear thermal conductivity parameter increases.

4) When nonlinear thermal conductivity and conductive-convective parameters increase, the rate of heat transfer at the fin base increases.

The developed analytical solution has given some physical insight into the thermal performance of the extended surfaces. The parametric study provides a good platform for proper design of the radiating fin for heat transfer enhancement in thermal systems.

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Dr. **Sobamowo M. Gbeminiyi** received OND and H.N.D form The Polytechnic, Ibadan in 1998 and 2002, respectively. He also obtained B.Sc., M.Sc. and Ph.D. in 2006, 2009 and 2013, respectively in the Department of Mechanical Engineering, University of Lagos, Nigeria. Where he lectures currently. Dr. Sobamowo has published over 240 research papers in various international journals. He is an author of Engineering textbooks and a reviewer for many international and local journals. His research interests include energy systems modelling, simulation and design, fluid flow and heat transfer analysis, renewable energy systems analysis, and thermal fluidic-induced vibration in energy systems. He is a member of the Nigerian Institution of Mechanical Engineers, Nigeria Society of Engineers, Council of Regulation of Engineering and International Association of Engineers. Contributions in this paper: ideas and formulation of the research. model development. writing – original draft preparation. writing – review and editing. application of numerical method to analyze the problem.



Dr. **Olurotimi A. Adeleye** obtained his Ph.D. degree in the area of computational mechanics in systems engineering, from the University of Lagos, Nigeria. He has worked as a Lecturer and Researcher in the Department of Systems Engineering, University of Lagos, but now he works as a Senior Lecturer and Researcher in the Department of Biomedical Engineering, in the same University. Dr. Adeleye has over 40 research publications in various international journals, and he is also a reviewer for some local and international journals. His research interests include Computational Mechanics, Biomechanics, Biomaterials, Nanotechnology and the applications of analytical, numerical, and Artificial Intelligent tools in obtaining solutions to nonlinear problems in these areas. He has published about forty (40) research articles in reputable journals. Contributions in this paper: application of mathematical method to analyze the problem. implementation of the computer code and supporting algorithms; testing of existing code components.



Yinusa Ahmed Amoo is a Lecturer and a Ph.D. candidate in the Department of Mechanical Engineering, University of Lagos, Nigeria. He obtained his B.Sc. and M.Sc. in the same department in 2014 and 2017, respectively. He has published over 80 papers in various local and international journals, and he is also a reviewer for some local and international journals. His research interests include Nanotechnology and its applications, Adaptation to climate change, thermal fluidic-induced vibration in energy systems and applied mechanics. He is a member of the Nigerian Institution of Mechanical Engineers and Council of Regulation of Engineering in Nigeria. Contributions in this paper: verification of the overall replication/reproducibility of results. writing of original draft preparation.



Dr. **Manasseh Olusegun Oyekeye** lectures in the department of Mechanical Engineering, University of Lagos, Nigeria. He obtained the Bachelor of Engineering with Second class honors, Upper division from University of Ilorin in 1986. He also obtained Master of Engineering degree from the University of Ibadan in 2003 and Ph.D. in 2018 in University of Lagos, Nigeria. He was involved in production and maintenance engineering between 1987 and 2004. After which, he joined the academics at the University of Lagos in 2004 and he has been involved in teaching and research work till date. He has published some research papers in some international journals. Contributions in this paper: verification of the overall replication/reproducibility of results. writing of original draft preparation



M. A. Waheed is a professor of thermofluids at the Department of Mechanical Engineering, College of Engineering, Federal University of Agriculture, Abeokuta, Nigeria. He undertook his academic training at the University of Ilorin (for his B.Eng. and M.Eng. degrees) and Aachen University of Technology (RWTH), Aachen, Germany (for his Ph.D.). His area of specialization is Thermofluids, Heat and Mass Transfer, Computational Fluid Dynamics and Energy Studies. He is a Fellow, Nigerian Society of Engineers (FNSE), Fellow, Nigerian Institution of Mechanical Engineers (FNIMECHE), Fellow, the Institute of Management Consultants (FIMC) and a Nigerian COREN Registered Engineer. He has served in many capacities in the University System as Head of Department, Head of Research Group, Deputy Dean of a College, Director of Academic Planning and Deputy Vice-Chancellor Academic at the Federal University of Agriculture, Abeokuta-Nigeria, FUNAAB. Contributions in this paper: management, coordination, oversight and leadership responsibilities for the research activity, planning and execution. provision of study materials and analysis tools.