Time domain response analysis of multi-flow channel hydraulic mount

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Abstract. In order to analyze the effect of the combination of long and short inertia channels and orifice flow channels on the time domain response of hydraulic mounts. Firstly, six hydraulic mounts with different combinations of inertia channels and orifice flow channels are proposed. And then, the transfer functions of dynamic stiffness and upper chamber pressure for six structures of hydraulic mounts are derived using the lumped parameter method. Next, the time domain analytic formulas for the transfer force and upper chamber pressure for six structural hydraulic mounts under steady-state excitation and step excitation are obtained using the convolution method. Finally, the analytical formula is compared with the hydraulic mount's model built by AMEsim; Meanwhile, the effects of inertia terms of inertia channels, damping, and damping of orifice flow channels on hydraulic mounts transfer forces are analyzed; Analyze the effect of transfer force variation and excitation amplitude on hydraulic mounts damping for different configurations of structures. Research shows that inertia channels and orifice flow channels directly affect the low-frequency dynamic characteristics of hydraulic mounts. At the same time, the effective damping height of the hydraulic mounts depends on the excitation amplitude.

Keywords: long and short inertia channel; orifice flow channel; hydraulic mount; lumped parameters; transfer force.

1. Introduction

The mounts are vibration isolation elements connected between the engine and the vehicle frame, supporting the static engine load-bearing capacity, isolating the engine vibration transmitted to the vehicle frame, and reducing the impact of road impact on the engine, limiting the engine movement space. The ideal powertrain components should exhibit large stiffness and large damping at low frequencies and low stiffness and low damping at high frequencies to achieve vehicle vibration isolation performance under different operating conditions [1]. Due to the strong frequency and amplitude dependence of the mounts themselves, passive hydraulic mounts that are superior to passive rubber mounts are proposed [2]. Meanwhile, semi-active or active mounts have recently become a major research topic [3-5]. However, passive hydraulic mounts are still widely used due to their simple design and low cost.

Singh [6] proposed a linear time-invariant model with aggregate mechanics and fluid cells and validated the model by comparing dynamic stiffness predictions with experimental data for a frequency range of 1-50 Hz, emphasizing the modeling of free and fixed types of decouplers. Tiwari [7] conducted an experimental study on the hydraulic mount of an inertial channel and free-floating decoupler, identifying the inertial channel resistance, the top and bottom chamber compliance, and the effective piston area of the top chamber, and characterizing it in the frequency and time domains. Yoon and Singh [8, 9] proposed an indirect method based on the quasi-linear fluid system formulation to estimate the power transmitted by a fixed or free decoupled hydraulic

mount on a rigid base to compensate for the inaccuracy of the mechanical system model. Farzad [10] proposed a variable stiffness decoupler membrane structure to replace the conventional decoupled membrane channel, and the new decoupled membrane structure can adjust the membrane stiffness according to the excitation frequency to achieve the optimal vibration isolation performance of the mounts. Liao [11] used ABAQUS software to build a fluid-structure coupling model of multi-inertial channel hydraulic mounts and to identify the structural parameters.

Studies on multi-inertia channel hydraulic mounts or bushings are mainly conducted by Zhang [12] to study the effects of the number, size, and length of inertia channels on the low-frequency dynamic performance of hydraulic mounts, revealing that different numbers of inertia channels can change the relationship between the stiffness and loss angle of hydraulic mounts and the excitation frequency. Yang [13, 14] established a multi-inertia channel-multi-throttle orifice type hydraulic bushing set parameter model, and derived formulas for stiffness and loss angle to analyze the relationship between the dynamic characteristics of multi-channel hydraulic bushings and the number of channels. It was shown that increasing the number of inertia channels and/or orifices can significantly improve the dynamic stiffness, the amplitude of loss angle, and the corresponding frequency of the bushing, thus greatly improving the performance of the bushing. Tan Chai [15-17] developed a controllable conceptual hydraulic bushing model with two parallel unequal inertia channels and a flow control unit to analyze the dynamic and time-domain characteristics of a multi-flow channel bushing numerically and experimentally. Lu [18] studied the n = 2 inertial channels mainly to derive a linear set of total parameters model for the hydraulic mount with n = 2 and analyzed its dynamic characteristics. Li [19] proposed an inertial channel optimization method based on a linearized low-frequency model excluding the effect of decoupling membranes, considering multiple arrangements of inertial channel and orifice combinations. Also, the predictions of the hydraulic mounts model were referred to the relevant references by Shishegaran [20-30]. It is shown that the use of two configurations can reduce the relative engine and body displacement transmissibility. Although studies on the time and frequency domain characteristics of multi-inertia channel hydraulic bushings have been reported, there are still some differences between the hydraulic bushing and hydraulic mount characteristics [31], and the time-domain characteristics of multi-flow channel hydraulic mounts cannot be well understood by hydraulic bushing time-domain analysis. Meanwhile, although the analysis of dynamic characteristics of multi-flow channel hydraulic mounts has been studied, the frequency domain characteristics and transient characteristics of multi-flow channel hydraulic mounts considering various combinations simultaneously have not been reported yet.

Based on the above review, the specific research objectives of this paper are as follows: To propose a lumped parameter model of the hydraulic mounts for 6 configurations of long and short inertia channels combined with orifice flow channels, and to solve the corresponding dynamic stiffness and upper chamber pressure transfer functions. Solve the time domain expressions for the steady-state response and ideal step response of the transfer force and upper chamber pressure based on the derived dynamic stiffness and upper chamber pressure transfer functions. The effects of upper chamber compliance, fluid viscosity, inertia, and damping of long and short inertia channels and orifice flow channel damping on hydraulic mount transfer force are analyzed based on time-domain expressions, and the effect of excitation amplitude on damping is also analyzed.

2. Hydraulic mount prototype study

2.1. Hydraulic mount structure

Fig. 1 shows a schematic diagram of the hydraulic mount, where Fig. 1(a) shows the structure of the passive hydraulic mount, and Fig. 1(b) shows the lumped parameter model of the passive hydraulic mount. In this case, the hydraulic mount structure includes an inertial channel and an orifice flow channel. The flow rates in the inertial and orifice flow channels are denoted by q_i and

 q_o , respectively. The rubber unit is modeled using the rubber stiffness K_r and the viscous damping coefficient B_r . The compliances of the upper and lower chambers of the hydraulic mount are C_1 and C_2 respectively, and the pressure of the upper chamber #1 is P_1 , and the pressure of the lower chamber #2 is P_2 . where the inertial channel and the orifice flow channel are represented by the fluid inertia I_i and I_o and the fluid resistance R_i and R_o , respectively.



Fig. 1. Hydraulic mount: a) structure diagram, b) lumped parameter model

2.2. Hydraulic mount flow channel design with multiple configurations

From the literature [12, 32], it is known that increasing the number of hydraulic mount inertia channels can improve the peak frequency of the dynamic stiffness and loss angle of the mount, and the combination of different flow channels can broaden the range of hydraulic mount vibration isolation frequency. However, these literature studies are all from the mount frequency characteristics and rarely involve the analysis of multi-flow channel hydraulic mount time domain characteristics. To investigate common multi-flow hydraulic mounts, six prototype configurations as shown in Table 1 and Fig. 2 were analyzed.



Fig. 2. Multiple configurations of the hydraulic mount prototype a) single long inertia channel, configuration Z₁, (L₁ = 212 mm, D₁ = 8.53 mm), configuration Z₂ (L₂ = 4L₁, D₂ = D₁); b) parallel combination of long inertia channel and orifice flow channel, configuration Z₃ (L₃ = 4L₁, D₃ = D₁, L_{o1} = 2 mm, D_{o1} = 0.2 mm); configuration Z₄ (L₄ = 4L₁, D₄ = D₁, L_{o2} = L_{o1}, D_{o2} = 5D_{o1});
c) parallel combination of three long inertia channels, configuration Z₅ (L₅₁ = L₅₂ = L₁, L₅₃ = 4L₁, D₅₁ = D₅₂ = D₅₃ = D₁), configuration Z₆ (L₅₁ = L₅₂ = L₅₃ = L₁, D₅₁ = D₅₂ = D₅₃ = D₁)

The first configuration (named Z_1 , Z_2) hydraulic mount upper and lower chambers are connected only by long inertia channels, this structure is the most common at present. Where, the length of Z_1 inertia channel $L_1 = 212$ mm, the length of Z_2 inertia channel $L_2 = 4L_1$, the diameter of Z_1 and Z_2 inertia channels are equal $D_1 = D_2 = 8.53$ mm. In the second configuration (named Z_3 and Z_4), the upper and lower chambers of the hydraulic mount are connected in parallel by long inertia channels and orifice flow channel. Where Z_3 and Z_4 inertia channels have equal lengths and diameters, $L_3 = L_4 = 4L_1$, $D_3 = D_4 = 8.53$ mm, respectively; The length of the Z_3 configuration orifice flow channel and Z_4 are equal to $L_{o1} = L_{o2} = 2$ mm, the diameter of the Z_3 orifice flow channel is $D_{o1} = 0.2$ mm, and the diameter of the Z_4 orifice flow channel is $D_{o2} = 1$ mm. The third configuration (named Z_5, Z_6) is a parallel combination of the number of long inertia channels n = 3. Where the diameters of the Z_5 and Z_6 inertial channels are equal $D_5 = D_6 = 8.53$ mm, $L_{51} = L_{52} = L_1$, $L_{53} = 4L_1$ in the Z_5 configuration; the lengths of the three inertial channels in the Z_6 configuration are equal $L_{61} = L_{62} = L_{63} = L_1$.

Fig. 3 shows the effect of variation in inertial channel length and cross-sectional area on the flow rate of the decoupler membrane channel. The decoupler membrane channel flow increases slightly with increasing inertial channel length. The flow rate of the decoupler membrane channel decreases slightly as the cross-sectional area of the inertial channel increases.

with inertia channel and orifice flow channel						
Design	п	Description (also see Fig. 2)				
Z_1	1	Baseline configuration (A long inertia channel, L_1 , D_1)				
Z_2	1	A long inertia channel ($L_2 = 4L_1, D_2 = D_1$)				
Z_3	2	Parallel combination of long inertia channel and orifice flow channel $(L_3 = 4L_1,$				
		$D_3 = D_1, L_{o1}, D_{o1}$				
Z_4	2	Parallel combination of long inertia channel and orifice flow channel $(L_3 = 4L_1,$				
		$D_3 = D_1, L_{o2} = L_{o1}, D_{o2} = 5D_{o1}$				
Z_5	3	Long inertia channel parallel combination ($L_{51} = L_{52} = L_1, L_{53} = 4L_1$,				
		$D_{51} = D_{52} = D_{53} = D_1$				
Z_6	2	Long inertia channel parallel combination ($L_{61} = L_{62} = L_{63} = L_1$,				
	3	$D_{51} = D_{52} = D_{53} = D_1$				

Table 1.	. Study configuration: prototype hydraulic engine mounts
	with inertia channel and orifice flow channel



Fig. 3. Effect of hydraulic mount inertia channel length and cross-sectional area variation on decoupling membrane channel flow: a) length variation; b) cross-sectional area variation

2.3. AMEsim-based hydraulic mount study

Fig. 4 shows the hydraulic mounts models of six configurations built using AMEsim software. Fig. 4(a) shows the hydraulic mounts of Z_1 and Z_2 configurations, Fig. 4(b) shows the hydraulic mounts of Z_3 and Z_4 configurations, and Fig. 4(c) shows the hydraulic mounts of Z_5 and Z_6 configurations.

Build 6 hydraulic mounts configurations using AMEsim, and the dynamic characteristics of the combined inertial channel and orifice flow channel hydraulic mounts were simulated for the above six configurations. The test uses a sinusoidal displacement excitation $x_e(t) = A\sin\omega t$, where A is the amplitude of the harmonic input and $f = \omega/2\pi$ is the frequency of the excitation, where the unit is Hz. The transfer forces of the hydraulic mounts and the dynamic pressures $P_1(t)$ and $P_2(t)$ within the upper and lower chambers are calculated in 1 Hz increases between 1 and 30 Hz in $X_e = 0.3$ mm, $X_e = 0.9$ mm and, where $X_e = 2A$ is the peak-to-peak (P-P) value. The calculated dynamic stiffness and loss angle for a given excitation amplitude $X_e = 0.3$ mm and $X_e = 0.9$ mm hydraulic mounts can be obtained as shown in Fig. 4. It is worth noting that configurations Z_5 and Z_6 exhibit greater dynamic stiffness and loss angle peak frequency. While configurations Z_3 and Z_4 show the smallest peak frequencies for dynamic stiffness and hysteresis angle, this is due to the inclusion of orifice flow channels in both configurations. Finally, the dynamic stiffness and loss angle of Z_1 and Z_2 are relatively small relative to the peak, due to the number of flow channels of one in the configuration. When the excitation amplitude X_e is increased from 0.3 mm to 0.9 mm, the dynamic stiffness and hysteresis angle of the hydraulic mounts decrease, which is more obvious for Z_1 , Z_2 , Z_5 and Z_6 .



Fig. 4. AMEsim hydraulic mount model: a) single inertia channel; b) combination of inertia channel and orifice flow channel; c) combination of 3 inertia channels



Fig. 5. Dynamic stiffness and loss angle of six configurations of hydraulic mount: a) $X_e = 0.3$ mm; b) $X_e = 0.9$ mm

3. Multi-flow channel hydraulic mount analysis model

3.1. Model I single inertia channel hydraulic mount

When the hydraulic mount is excited by $x_e(t)$, applying the continuity equation to the fluid control hydraulic unit in combination with Fig. 1 and Fig. 1(a), the following Eq. (1)-(5) are obtained. Where the fluid inertia I_i and the fluid resistance R_i are expressed:

$$P_{1}(t) - P_{2}(t) = I_{i}q_{i}(t) + R_{i}q_{i}(t),$$
(1)

$$C_{1}P_{1}(t) = A_{p}\dot{x}_{e}(t) - q_{i}(t),$$

$$C_{i}\dot{P}_{i}(t) = q_{i}(t)$$
(2)
(3)

$$I_{i} = \frac{\rho L}{A_{i}},$$
(4)

$$R_i = \frac{128 visL}{\pi D_h^4},\tag{5}$$

where ρ is the density of the fluid, *L* is the length of the flow channel, A_i is the cross-sectional area of the flow channel, and ν is the viscosity of the fluid, D_h is the equivalent "hydraulic diameter" of the flow channel. The first term on the right side of Eq. (1) is the pressure drop due to the fluid in the inertial channel due to inertia; the second term on the right side is the pressure drop due to the viscous damping of the fluid in the inertial channel.

Therefore, the hydraulic mount transfer force F_T under external excitation X_e is defined as follows:

$$F_T = B_r \dot{x}_e(t) + K_r x_e(t) + A_p P_1(t).$$
(6)

Then, by Laplace transformation of Eq. (6), the expressions for the variation of the transfer complex stiffness and the upper chamber pressure with the excitation amplitude for a single-flow channel hydraulic mount are obtained as follows:

$$K_{dyn}(s) = \frac{F_T}{x_e}(s) = K_r + B_r s + A_p^2 K_1 \frac{I_i s^2 + R_i s + K_2}{I_i s^2 + R_i s + K_2 + K_1},$$
(7)

$$\frac{P_1}{x_e}(s) = A_p K_1 \frac{I_i s^2 + R_i s + K_2}{I_i s^2 + R_i s + K_2 + K_1}.$$
(8)

Since $1/C_1 = K_1$, $1/C_2 = K_2$, and $K_1 \gg K_2$. So, Eqs. (7) and (8) can be further simplified as follows:

$$K_{dyn}(s) = \frac{F_T}{x_e}(s) = K_r + B_r s + A_p^2 K_1 \frac{I_i s^2 + R_i s}{I_i s^2 + R_i s + K_1},$$
(9)

$$\frac{P_1}{x_e}(s) = A_p K_1 \frac{I_i s^2 + R_i s}{I_i s^2 + R_i s + K_1}.$$
(10)

Eq. (9) and Eq. (10) can be simplified to a standard second-order system:

$$K_{d1}(s) = \frac{F_T}{x_e}(s) = K_r + B_r s + \lambda \frac{s^2 + 2\varsigma_2 \omega_{n2} s}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2},$$
(11)

$$\frac{P_1}{x_e}(s) = A_p K_1 \frac{I_i s^2 + R_i s}{I_i s^2 + R_i s + K_1} = A_p K_1 \frac{s^2 + 2\varsigma_2 \omega_{n2} s}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2}.$$
(12)

Thus, Eq. (11) can be transformed into the following standard form, where $\lambda = A_p^2 K_1$,

 $\tau = R_i / I_i$:

$$K_{d1}(s) = K_r + B_r s + \lambda \frac{(s+\tau)s}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2},$$
(13)

$$G(s) = \frac{P_1}{x_e}(s) = A_p K_1 \frac{(s+t)s}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2},$$
(14)

where $\varsigma_2 = \sqrt{R_i^2/4K_1I_i}, \, \omega_{n2} = \sqrt{K_1/I_i}.$

3.2. Model II combined inertial channel and orifice flow channel hydraulic mount

According to Fig. 1 and Fig. 2(b), the fluid equation of the hydraulic mount subjected to displacement excitation $x_e(t)$ for the parallel combination of inertial channel and orifice flow channel can be obtained as shown below:

$$P_{1} - P_{2} = I_{i}\dot{q}_{i} + R_{i}q_{i},$$

$$P_{1} - P_{2} = R_{0}q_{0},$$

$$(15)$$

$$P_{1} - P_{2} = R_{0}q_{0},$$

$$(16)$$

$$C_{1}\dot{P}_{1} = A_{p}\dot{x}_{e} - q_{i} - q_{0},$$

$$(17)$$

$$C_{2}\dot{P}_{2} = q_{i} + q_{0}.$$

$$(18)$$

In the equation, R_0 is the fluid resistance of the orifice flow channel. Since the inertia in the orifice flow channel is negligible, Eq. (16) only considers the viscous damping term resulting in the hydraulic mounts differential pressure change. Reference [33] shows that when the fluid is assumed to be in the form of laminar flow in the orifice flow channel, the fluid resistance of the orifice flow channel is solved for by the following expression:

$$R_0 = \frac{50.4vis}{\pi D_o^3},$$
(19)

where, D_o is the diameter of the orifice flow channel.

Thus, the hydraulic mounts parallel to the inertial channel and the orifice flow channel transfer the complex stiffness $K_{d2}(s)$ and the variation of the upper chamber pressure with the excitation amplitude G(s) defined as follows:

$$K_{d2}(s) = \frac{F_T}{x_e}(s)$$

$$= K_{e1} + R_{e2} + A_{e2}^2 K \qquad I_i R_0 s^2 + (K_2 I_i + R_i R_0) s + K_2 (R_i - R_0)$$
(20)

$$= K_r + B_r s + A_p^r K_1 \frac{1}{I_i R_0 s^2 + (K_2 I_i + R_i R_0 + K_1 I_i) s + (K_2 + K_1)(R_i + R_0)}{I_i R_0 s^2 + (K_2 I_i + R_i R_0) s + K_2 (R_i - R_0)}$$

$$G(s) = \frac{P_1}{x_e}(s) = A_p K_1 \frac{I_i R_0 s^2 + (K_2 I_i + R_i R_0 + K_1 I_i) s + (K_2 + K_1)(R_i + R_0)}{I_i R_0 s^2 + (K_2 I_i + R_i R_0 + K_1 I_i) s + (K_2 + K_1)(R_i + R_0)}.$$
(21)

Since $1/C_1 = K_1$, $1/C_2 = K_2$, and $K_1 \gg K_2$ can be simplified by Eqs. (20) and (21) as follows:

$$K_{dyn}(s) = \frac{F_T}{x_e}(s) = K_r + B_r s + A_p^2 K_1 \frac{I_i R_0 s^2 + R_i R_0 s}{I_i R_0 s^2 + (R_i R_0 + K_1 I_i) s + K_1 (R_i + R_0)'}$$
(22)

$$G(s) = \frac{P_1}{x_e}(s) = A_p K_1 \frac{I_i R_0 s^2 + R_i R_0 s}{I_i R_0 s^2 + (R_i R_0 + K_1 I_i) s + K_1 (R_i + R_0)}.$$
(23)

So, reduced to the standard second-order system form shown:

$$K_{dyn-2}(s) = \frac{F_T}{x_e}(s) = K_r + B_r s + \lambda \frac{s(s+\kappa)}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2},$$
(24)

$$G(s) = \frac{P_1}{x_e}(s) = A_p K_1 \left(\frac{s(s+\kappa)}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2} \right),$$
(25)

where the expressions for ω_{n2} , ς_2 , κ and F are shown:

$$\omega_{n2} = \sqrt{\frac{(R_i + R_0)K_1}{I_i R_0}}, \quad \varsigma_2 = \frac{(R_i R_0 + K_1 I_i)}{2} \sqrt{\frac{1}{K_1 I_i R_0 (R_i + R_0)}}, \quad \kappa = \frac{R_i}{I_i}, \quad \lambda = A_p^2 K_1.$$

3.3. Model III parallel combination of 3 flow channel hydraulic mount

To easily solve for the multi-inertia channel hydraulic mount transfer complex stiffness $K_{d3}(s)$ and the upper chamber pressure variation G(s) with excitation amplitude for the Z_5 and Z_6 configurations. The lumped parameter model of the multi-inertia channel hydraulic mounts with Z_5 and Z_6 configurations of Fig. 6(a) is equated to the equivalent mechanical model of Fig. 6(b). Because the cross-sectional area of the hydraulic mount inertia channel is much smaller than the equivalent cross-sectional area of the rubber main spring, when the fluid flows through the inertia channel, the fluid velocity is amplified to play the role of vibration attenuation, the literature [6] referred to it as a velocity amplification of large absorbing array.



Fig. 6. Multi-inertia channel hydraulic mount: a) lumped parameter model, b) equivalent mechanical model

According to Fig. 6(a), the mathematical model equation of the lumped parameters of the multi-inertia channel hydraulic mounts can be obtained as follows:

$C_1 \dot{P}_1 = A_p \dot{x}_e - q_{i1} - q_{i2} - q_{i3},$	(26)
$C_2 \dot{P}_2 = q_{i1} + q_{i2} + q_{i3},$	(27)
$P_1 - P_2 = I_{i1}\ddot{q}_{i1} + R_{i1}\dot{q}_{i1},$	(28)
$P_1 - P_2 = I_{i2} \dot{q}_{i2} + R_{i2} \dot{q}_{i2},$	(29)
$P_1 - P_2 = I_{i3} \dot{q}_{i3} + R_{i3} \dot{q}_{i3},$	(30)
$F_T = K_r x_e + B_e x_e + A_p P_1,$	(31)
$F_T = K_r x_e + B_e x_e + A_p P_1,$	(32)

where, q_{i1} , q_{i2} , q_{i3} are the flow rate through the three inertia channels of configuration Z_5 and Z_6 respectively, I_{i1} , I_{i2} , I_{i3} are the inertia flow through the three inertia channels respectively, R_{i1} , R_{i2} , R_{i3} are the liquid resistance flow through the three inertia channels respectively.

In Fig. 6(b), $k_1 = A_p^2/C_1$ is the equivalent linear stiffness of the upper chamber of the hydraulic mount, $k_2 = A_p^2/C_2$ is the equivalent linear stiffness of the lower chamber of the

hydraulic mount, b_{in} (n = 1, 2, 3) is the equivalent damping of the fluid in the inertia channel of the hydraulic mount, and m_{in} (n = 1, 2, 3) is the equivalent mass of the fluid in the inertia channel of the hydraulic mount [34]. It is known that the flow rate $q = A_i \dot{x}_i$, where A_i is the cross-sectional area of the inertial channel and \dot{x}_i is the average flow velocity of the fluid in the inertial channel. Therefore, Eqs. (32-35) can be written as shown in the following equation:

$$\frac{1}{C_1} \left(A_p x_e - A_{i1} x_{i1} - A_{i2} x_{i2} - A_{i3} x_{i3} \right) - \frac{1}{C_2} \left(A_{i1} x_{i1} + A_{i2} x_{i2} + A_{i3} x_{i3} \right)$$

$$= I_{i1} A_{i1} \dot{x}_{i1} + R_{i1} A_{i1} \dot{x}_{i1},$$
(33)

$$\frac{1}{C_1} \left(A_p x_e - A_{i1} x_{i1} - A_{i2} x_{i2} - A_{i3} x_{i3} \right) - \frac{1}{C_2} \left(A_{i1} x_{i1} + A_{i2} x_{i2} + A_{i3} x_{i3} \right) = I_{i2} A_{i2} \dot{x}_{i2} + R_{i2} A_{i2} \dot{x}_{i2},$$
(34)

$$\frac{1}{C_1} \left(A_p x_e - A_{i1} x_{i1} - A_{i2} x_{i2} - A_{i3} x_{i3} \right) - \frac{1}{C_2} \left(A_{i1} x_{i1} + A_{i2} x_{i2} + A_{i3} x_{i3} \right) = I_{i3} A_{i3} \ddot{x}_{i3} + R_{i3} A_{i3} \dot{x}_{i3},$$
(35)

$$F_T = K_r x_e + B_e x_e + A_p P_1 = K_r x_e + B_e x_e + \frac{1}{C_1} A_p \left(A_p x_e - A_{i1} x_{i1} - A_{i2} x_{i2} - A_{i3} x_{i3} \right),$$
(36)

where, A_{i1} , A_{i2} , A_{i3} are the cross-sectional areas of the three inertial channels in configurations Z_5 and Z_6 , and x_{i1} , x_{i2} , x_{i3} are the displacements of the fluid flowing through the three inertial channels in configurations Z_5 and Z_6 , respectively.

Combining the equivalent mechanical model of the multi-inertia channel hydraulic mount in Fig. 6(b) and the literature [2] the Eqs. (33-36) can be written as follows:

$$k_1(x_e - x_{i1} - x_{i2} - x_{i3}) - k_2(x_{i1} + x_{i2} + x_{i3}) = I_{i1}A_p^2 \ddot{x}_{i1} + R_{i1}A_p^2 \dot{x}_{i1},$$
(37)

$$k_1(x_e - x_{i1} - x_{i2} - x_{i3}) - k_2(x_{i1} + x_{i2} + x_{i3}) = I_{i2}A_p^2 \ddot{x}_{i2} + R_{i2}A_p^2 \dot{x}_{i2},$$
(38)

$$k_1(x_e - x_{i1} - x_{i2} - x_{i3}) - k_2(x_{i1} + x_{i2} + x_{i3}) = I_{i3}A_p^2 \dot{x}_{i3} + R_{i3}A_p^2 \dot{x}_{i3},$$
(39)

$$F_T = K_r x_e + B_e x_e + k_1 (x_e - x_{i1} - x_{i2} - x_{i3}).$$
(40)

Because $C_1 \ll C_2$, The equivalent masses of the three inertia channels of the hydraulic mounts of the Z_5 and Z_6 structures are $m_{i1} = I_{i1}A_p^2$, $m_{i2} = I_{i2}A_p^2$, $m_{i3} = I_{i3}A_p^2$, and the equivalent damping of the three inertial channels of the hydraulic mounts of the Z_5 and Z_6 are $b_{i1} = R_{i1}A_p^2$, $b_{i2} = R_{i2}A_p^2$, $b_{i3} = R_{i3}A_p^2$ respectively. Therefore, Eqs. (37)-(39) can be further simplified as shown below:

$$K_1(x_e - x_{i1} - x_{i2} - x_{i3}) = m_{i1}\ddot{x}_{i1} + b_{i1}\dot{x}_{i1},$$
(41)

$$K_1(x_e - x_{i1} - x_{i2} - x_{i3}) = m_{i2}\ddot{x}_{i2} + b_{i2}\dot{x}_{i2}, \tag{42}$$

$$K_1(x_e - x_{i1} - x_{i2} - x_{i3}) = m_{i3}\ddot{x}_{i3} + b_{i3}\dot{x}_{i3}.$$
(43)

Referring to the literature [32], the equivalent mass and equivalent damping of the inertial channel can be expressed as follows:

$$R_{eq} = \frac{\prod_{i=1}^{3} R_{ii}}{\sum_{i=1}^{3} (R_{ii})} = \frac{R_{i1} R_{i2} R_{i3}}{R_{i1} + R_{i2} + R_{i3}} \Leftrightarrow b_{eq} = A_p^2 R_{eq},$$
(44)

$$I_{eq} = \frac{\prod_{i=1}^{3} R_{ii}}{\sum_{i=1}^{3} (R_{ii})} = \frac{R_{i1} R_{i2} R_{i3}}{R_{i1} + R_{i2} + R_{i3}} \Leftrightarrow m_{eq} = A_p^2 I_{eq}.$$
(45)

Therefore, the Eqs. (40-43) can be reduced to the following expression:

$$K_1(x_e - x_{eq}) = m_{eq} \dot{x}_{eq} + b_{eq} \dot{x}_{eq}, \tag{46}$$

$$F_T = K_r x_e + B_e x_e + A_p P_1 = K_r x_e + B_e x_e + k_1 (x_e - x_{eq}).$$
(47)

As a result, the transfer function expressions for the transfer complex stiffness and the upper chamber pressure variation with excitation for the multi-inertia channel hydraulic mounts of Z_5 and Z_6 structures are as follows:

$$\frac{F_T}{x_e}(s) = K_r + B_r s + \frac{k_1 m_{eq} s^2 + k_1 b_{eq} s}{m_{eq} s^2 + b_{eq} s + k_1},$$
(48)

$$G(s) = \frac{P_1}{x_e}(s) = \frac{k_1}{A_p} \frac{m_{eq}s^2 + b_{eq}s}{m_{eq}s^2 + b_{eq}s + k_1}.$$
(49)

Therefore, the expressions for Eqs. (48) and (49) reduced to a standard second-order system are shown below:

$$K_{d3}(s) = \frac{F_T}{x_e}(s) = K_r + B_r s + \lambda \frac{s(s+\kappa)}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2},$$
(50)

$$G(s) = \frac{P_1}{x_e}(s) = \frac{R_1}{A_p} \frac{s(s+k)}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2},$$
(51)

where the expressions for κ , ω_{n2} and ς_2 are as follows:

$$\kappa = \frac{b_{eq}}{m_{eq}}, \quad \omega_{n2} = \sqrt{\frac{K_1}{m_{eq}}}, \quad \varsigma_2 = \frac{b_{eq}}{2}\sqrt{\frac{1}{K_1 m_{eq}}}, \quad \lambda = k_1.$$

The key parameters of the hydraulic mount can be obtained by combining model I, model II and mode III as shown in Table 2.

System parameters	Model I	Model II	Mode III
λ	$A_p^2 K_1$	$A_p^2 K_1$	<i>k</i> ₁
К	$\frac{R_i}{I_i}$	$\frac{R_i}{I_i}$	$rac{b_{eq}}{m_{eq}}$
ω_{n2}	$\sqrt{\frac{K_1}{I_i}}$	$\sqrt{\frac{(R_i + R_0)K_1}{I_i R_0}}$	$\sqrt{rac{K_1}{m_{eq}}}$
ς2	$\sqrt{\frac{R_i^2}{4K_1I_i}}$	$\frac{(R_i R_0 + K_1 I_i)}{2} \sqrt{\frac{1}{K_1 I_i R_0 (R_i + R_0)}}$	$\frac{b_{eq}}{2} \sqrt{\frac{1}{K_1 m_{eq}}}$

Table 2. Key parameters of the system for the analytical model.

4. Multi-fluid hydraulic mount time domain analytical formula

4.1. Time domain steady-state harmonic response

To analyze the hydraulic mounts in Models I, II, and III configurations for steady-state harmonic response, the hydraulic mounts transfer the force $F_T(t)$ and the pressure $P_1(t)$ in the upper chamber, the specific derivation process was referred to the reference [15, 16]. According to the convolution of the bishop function:

$$L[x(t) * y(t)] = X(s)Y(s),$$

where x(t) * y(t) is the mathematical expression of the convolution.

(52)

Therefore, the analytical equations of the transfer force $F_T(t)$ and the pressure $P_1(t)$ in the upper chamber under the steady-state harmonic response in the time domain for models I, II and III of the hydraulic mount configurations can be solved according to Eq. (52), which is expressed as the following equation:

$$F_T(t) = K_d(t) * x_e(t) = \int_0^t K_d(\tau) x_e(t-\tau) d\tau,$$
(53)

$$P_1(t) = G(t) * x_e(t) = \int_0^t G(\tau) x_e(t-\tau) d\tau.$$
(54)

4.1.1. Model I hydraulic mount

For the single inertial channel of model I, the third e hydraulic path of the right-hand side of Eq. (13) and Eq. (14) can be written as the following expressions:

$$K_{d1-h}(s) = \lambda \frac{(s+\tau)s}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2} = \lambda \left[1 - \frac{\omega_{n2}^2}{s^2 + 2\varsigma_2 \omega_{n2} + \omega_{n2}^2} \right],$$
(55)

$$G(s) = A_p K_1 \left[1 - \frac{\omega_{n2}^2}{s^2 + 2\varsigma_2 \omega_{n2} + \omega_{n2}^2} \right].$$
(56)

The inverse Laplace transform of Eqs. (55-56) is shown:

$$\begin{split} K_{d1-h}(t) &= L^{-1} \left\{ \lambda \left[1 - \frac{\omega_{n2}^2}{s^2 + 2\varsigma_2 \omega_{n2} + \omega_{n2}^2} \right] \right\} = \lambda \left[\delta(t) - \frac{\omega_{n2}}{\sqrt{1 - \varsigma_2^2}} e^{-\varsigma_2 \omega_{n2} t} \sin(\omega_d t) \right], \quad (57) \\ P_1(t) &= L^{-1} \left\{ A_p K_1 \left[1 - \frac{\omega_{n2}^2}{s^2 + 2\varsigma_2 \omega_{n2} + \omega_{n2}^2} \right] \right\} \\ &= A_p K_1 \left[\delta(t) - \frac{\omega_{n2}}{\sqrt{1 - \varsigma_2^2}} e^{-\varsigma_2 \omega_{n2} t} \sin(\omega_d t) \right], \end{split}$$

where $\omega_d = \omega_{n2}\sqrt{1-\varsigma_2^2}$ is the damped self-oscillation angular frequency of the system.

Therefore, combining Eq (53) and (54) give the following expression for the sine excitation $x_e(t) = A\sin\omega t$, the analytic expression for the model I hydraulic mounts transfer force $F_T(t)$ and the pressure $P_1(t)$ in the upper chamber:

$$F_{T}(t) = K_{r}A\sin\omega t + B_{r}A\omega\cos\omega t + A\lambda\sin\omega t$$

$$-A\lambda \frac{\omega_{n2}}{2\sqrt{1-\varsigma_{2}^{2}}} \left(\frac{e^{-at}b\sin\omega_{d}t - e^{-at}a\cos\omega_{d}t + b\sin\omega t + a\cos\omega t}{\omega^{2} + \omega_{n2}^{2} + 2\omega_{d}\omega} \right)$$

$$+ \frac{e^{-at}c\sin\omega_{d}t + e^{-at}a\cos\omega_{d}t - c\sin\omega t - a\cos\omega t}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} \right),$$

$$P_{1}(t) = AA_{p}K_{1}\sin\omega t - \frac{\omega_{n2}}{2\sqrt{1-\varsigma_{2}^{2}}} \left(\frac{e^{-at}b\sin\omega_{d}t - e^{-at}a\cos\omega_{d}t + b\sin\omega t + a\cos\omega t}{\omega^{2} + \omega_{n2}^{2} + 2\omega_{d}\omega} - \frac{-e^{-at}c\sin\omega_{d}t - e^{-at}a\cos\omega_{d}t + c\sin\omega t + a\cos\omega t}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} \right).$$
(60)

The expressions of a, b and c, in the form are as follows: $a = \omega_{n2}\varsigma_2$, $b = \omega + \omega_d$, $c = \omega - \omega_d$.

4.1.2. Model II hydraulic mount

Re-express the third term on the right-hand side of Eqs. (24) and (25) as Eqs. (61) and (61):

$$K_{d2-h}(s) = \lambda \left[1 - \frac{\omega_{n2}^2}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2} - \frac{s(2\varsigma_2 \omega_{n2} - \kappa)}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2} \right],\tag{61}$$

$$G(s) = \frac{P_1}{x_e}(s) = A_p K_1 \left(1 - \frac{\omega_{n2}^2}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2} - \frac{s(2\varsigma_2 \omega_{n2} - \kappa)}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2} \right).$$
(62)

The following expression for the inverse Laplace transforms of Eqs. (61-62):

$$K_{d2-h}(t) = \lambda \left[\delta(t) - \frac{\omega_{n2}}{\sqrt{1 - \varsigma_2^2}} e^{-\varsigma_2 \omega_{n2} t} \sin(\omega_d t) + \frac{2\varsigma_2 \omega_{n2} - \kappa}{\sqrt{1 - \varsigma_2^2}} e^{-\varsigma_2 \omega_{n2} t} \sin(\omega_d t - \theta) \right], \quad (63)$$

$$P_{1}(t) = A_{p}K_{1} \left[\delta(t) - \frac{\omega_{n2}}{\sqrt{1 - \varsigma_{2}^{2}}} e^{-\varsigma_{2}\omega_{n2}t} \sin(\omega_{d}t) + \frac{2\varsigma_{2}\omega_{n2} - \kappa}{\sqrt{1 - \varsigma_{2}^{2}}} e^{-\varsigma_{2}\omega_{n2}t} \sin(\omega_{d}t - \theta) \right], \quad (64)$$

where $\theta = \tan^{-1} \left(\sqrt{1 - \varsigma_2^2} / \varsigma_2 \right)$.

Combining Eqs. (53) and (54), the hydraulic mount with parallel combination of inertia channel and orifice flow channel in model II configuration can be obtained for the transfer force $F_T(t)$ and the pressure $P_1(t)$ in the upper chamber at a sinusoidal excitation $x_e(t) = A\sin\omega t$:

$$F_{T}(t) = K_{r}Asin\omega t + B_{r}A\omega cos\omega t + A\lambda sin\omega t$$

$$-A\lambda \frac{\omega_{n2}}{2\sqrt{1-\varsigma_{n2}^{2}}} \left(\frac{e^{-at}bsin\omega_{d}t - e^{-at}acos\omega_{d}t + bsin\omega t + acos\omega t}{\omega^{2} + \omega_{n2}^{2} + 2\omega_{d}\omega} + \frac{e^{-at}csin\omega_{d}t + e^{-at}acos\omega_{d}t - csin\omega t - acos\omega t}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} \right) + A\lambda \frac{(2\varsigma_{n2}\omega_{n2} - \kappa)}{2\sqrt{1-\varsigma_{n2}^{2}}}$$

$$\left(\frac{e^{-at}bsin(\omega_{d}t - \theta) - e^{-at}acos(\omega_{d}t - \theta) + bsin(\omega t - \theta) + acos(\omega t - \theta)}{\omega^{2} + \omega_{n2}^{2} + 2\omega_{d}\omega} + \frac{e^{-at}csin(\omega_{d}t - \theta) + e^{-at}acos(\omega_{d}t - \theta) - csin(\omega t - \theta) - acos(\omega t - \theta)}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} \right),$$

$$P_{1}(t) = AA_{p}K_{1} \left\{ sin\omega t - \frac{\omega_{n2}}{2\sqrt{1-\varsigma_{n2}^{2}}} \cdot \left(\frac{-e^{-at}acos\omega_{d}t + e^{-at}bsin\omega_{d}t + acos\omega t + bsin\omega t}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{-e^{-at}acos\omega_{d}t - e^{-at}csin(\omega_{d}t - \theta) + e^{-at}bsin(\omega_{d}t - \theta)}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} \right) + \frac{(2\varsigma_{n2}\omega_{n2} - \kappa)}{2\sqrt{1-\varsigma_{n2}^{2}}} \left(\frac{\left\{ -e^{-at}acos(\omega_{d}t - \theta) + e^{-at}bsin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{-e^{-at}acos(\omega_{d}t - \theta) - e^{-at}csin(\omega_{d}t - \theta) + e^{-at}bsin(\omega_{d}t - \theta)}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}acos(\omega_{d}t - \theta) - e^{-at}csin(\omega_{d}t - \theta) + e^{-at}bsin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}acos(\omega_{d}t - \theta) - e^{-at}csin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} \right) \right\}$$

4.1.3. Model III hydraulic mount

Re-express the third term on the right-hand side of Eqs. (50) and (51) as Eqs. (67) and (68):

$$K_{d3-h}(s) = \lambda \left[1 - \frac{\omega_{n2}^2}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2} - \frac{s(2\varsigma_2 \omega_{n2} - \kappa)}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2} \right], \tag{67}$$

$$G(s) = \frac{P_1}{x_e}(s) = \frac{k_1}{A_p} \left(1 - \frac{\omega_{n2}^2}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2} - \frac{s(2\varsigma_2 \omega_{n2} - \kappa)}{s^2 + 2\varsigma_2 \omega_{n2} s + \omega_{n2}^2} \right).$$
(68)

The following expression for the inverse Laplace transforms of Eqs. (67), (68):

$$K_{d2-h}(t) = \lambda \left[\delta(t) - \frac{\omega_{n2}}{\sqrt{1 - \varsigma_2^2}} e^{-\varsigma_2 \omega_{n2} t} \sin(\omega_d t) + \frac{2\varsigma_2 \omega_{n2} - \kappa}{\sqrt{1 - \varsigma_2^2}} e^{-\varsigma_2 \omega_{n2} t} \sin(\omega_d t - \theta) \right], \quad (69)$$

$$P_{1}(t) = \frac{k_{1}}{A_{p}} \bigg[\delta(t) - \frac{\omega_{n2}}{\sqrt{1 - \varsigma_{2}^{2}}} e^{-\varsigma_{2}\omega_{n2}t} \sin(\omega_{d}t) + \frac{2\varsigma_{2}\omega_{n2} - \kappa}{\sqrt{1 - \varsigma_{2}^{2}}} e^{-\varsigma_{2}\omega_{n2}t} \sin(\omega_{d}t - \theta) \bigg].$$
(70)

Combining Eqs. (53) and (54), the hydraulic mount for the parallel combination of inertial channels in the Model III configuration can be obtained for the transfer force $F_T(t)$ and the pressure $P_1(t)$ in the upper chamber at a sinusoidal excitation $x_e(t) = A \sin \omega t$:

$$F_{T}(t) = K_{r}A\sin\omega t + B_{r}A\omega\cos\omega t + A\lambda\sin\omega t$$

$$-A\lambda \frac{\omega_{n2}}{2\sqrt{1-\varsigma_{n2}^{2}}} \left(\frac{e^{-at}b\sin\omega_{d}t - e^{-at}a\cos\omega_{d}t + b\sin\omega t + a\cos\omega t}{\omega^{2} + \omega_{n2}^{2} + 2\omega_{d}\omega} + \frac{e^{-at}c\sin\omega_{d}t + e^{-at}a\cos\omega_{d}t - c\sin\omega t - a\cos\omega t}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} \right) + A\lambda \frac{(2\varsigma_{n2}\omega_{n2} - \kappa)}{2\sqrt{1-\varsigma_{n2}^{2}}}$$
(71)
$$\cdot \left(\frac{e^{-at}b\sin(\omega_{d}t - \theta) - e^{-at}a\cos(\omega_{d}t - \theta) + b\sin(\omega t - \theta) + a\cos(\omega t - \theta)}{\omega^{2} + \omega_{n2}^{2} + 2\omega_{d}\omega} + \frac{e^{-at}c\sin(\omega_{d}t - \theta) + e^{-at}a\cos(\omega_{d}t - \theta) - c\sin(\omega t - \theta) - a\cos(\omega t - \theta)}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} \right),$$

$$P_{1}(t) = \frac{Ak_{1}}{A_{p}} \left\{ \sin\omega t - \frac{\omega_{n2}}{2\sqrt{1-\varsigma_{n2}^{2}}} - \frac{(-e^{-at}a\cos\omega_{d}t + e^{-at}b\sin\omega_{d}t + a\cos\omega t + b\sin\omega t)}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{-e^{-at}a\cos\omega_{d}t - e^{-at}c\sin\omega_{d}t + a\cos\omega t + c\sin\omega t}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} \right)$$

$$+ \frac{(2\varsigma_{2}\omega_{n2} - \kappa)}{2\sqrt{1-\varsigma_{n2}^{2}}} \left\{ \frac{\left\{ -e^{-at}a\cos(\omega_{d}t - \theta) + e^{-at}b\sin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} + 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}a\cos(\omega_{d}t - \theta) - e^{-at}c\sin(\omega_{d}t - \theta) + e^{-at}b\sin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}a\cos(\omega_{d}t - \theta) - e^{-at}c\sin(\omega_{d}t - \theta) + e^{-at}b\sin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}a\cos(\omega_{d}t - \theta) - e^{-at}c\sin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}a\cos(\omega_{d}t - \theta) - e^{-at}c\sin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}a\cos(\omega_{d}t - \theta) - e^{-at}c\sin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}a\cos(\omega_{d}t - \theta) - e^{-at}c\sin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}a\cos(\omega_{d}t - \theta) - e^{-at}c\sin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}a\cos(\omega_{d}t - \theta) - e^{-at}c\sin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}a\cos(\omega_{d}t - \theta) - e^{-at}c\sin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}a\cos(\omega_{d}t - \theta) - e^{-at}c\sin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}a\cos(\omega_{d}t - \theta) - e^{-at}c\sin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}a\cos(\omega_{d}t - \theta) - e^{-at}c\sin(\omega_{d}t - \theta) \right\}}{\omega^{2} + \omega_{n2}^{2} - 2\omega_{d}\omega} - \frac{\left\{ -e^{-at}a\cos(\omega_{d}t - \theta) -$$

4.2. Transient response of ideal step excitation

Using Section 2, and the convolution method in Section 3.1, the analytical equations for the transfer force $F_T(t)$ and the pressure $P_1(t)$ in the upper chamber at step excitation $x_e(t) = A_1 u(t)$ are derived for the hydraulic mounts configurations in Models I, II, and III.

The analytical equations for the transfer force $F_T(t)$ and the pressure $P_1(t)$ in the upper chamber of the hydraulic mount of model I under the step excitation $x_e(t) = A_1u(t)$ are as follows:

$$F_{T}(t) = K_{r}A_{1}u(t) + B_{r}A_{1}\delta(t) + \int_{0}^{t} K_{d1-h}(\tau) x_{e}(t-\tau)d\tau = K_{r}A_{1}u(t) + B_{r}A_{1}\delta(t)$$

$$+ \frac{\lambda A_{1}}{\omega_{d}}e^{-at}(a\sin\omega_{d}t + \omega_{d}\cos\omega_{d}t),$$

$$P_{1}(t) = \frac{A_{p}K_{1}A_{1}}{\omega_{d}}e^{-at}\{a\sin\omega_{d}t - \omega_{d}\cos\omega_{d}t\}.$$
(73)

The analytical equations for the transfer force $F_T(t)$ and the pressure $P_1(t)$ in the upper chamber of the hydraulic mount of model II under the step excitation $x_e(t) = A_1 u(t)$ are as follows:

$$F_{T}(t) = K_{r}A_{1}u(t) + B_{r}A_{1}\delta(t) + \int_{0}^{t} k_{dyn-H}(\tau) x(t-\tau)d\tau$$

$$= K_{r}A_{1}u(t) + B_{r}A_{1}\delta(t) + \frac{\lambda A_{1}}{\omega_{d}}e^{-at}(asin\omega_{d}t + \omega_{d}cos\omega_{d}t)$$
(75)
$$-\lambda A_{1}\frac{2\varsigma_{2}\omega_{n2} - \kappa}{\omega_{d}\omega_{n2}} \{e^{-at}[asin(\omega_{d}t - \theta) + \omega_{d}cos(\omega_{d}t - \theta)] + asin\theta - \omega_{d}cos\theta\},$$

$$P_{1}(t) = \frac{A_{1}A_{p}K_{1}}{\omega_{d}}e^{-at}(asin\omega_{d}t + \omega_{d}cos\omega_{d}t)$$
(76)
$$+A_{1}K_{1}A_{p}\frac{2\varsigma_{2}\omega_{n2} - \kappa}{\omega_{d}\omega_{n2}} \{e^{-at}[asin(\omega_{d}t - \theta) + \omega_{d}cos(\omega_{d}t - \theta)] + asin\theta - \omega_{d}cos\theta\}.$$

The analytical equations for the transfer force $F_T(t)$ and the pressure $P_1(t)$ in the upper chamber of the hydraulic mount of model III under the step excitation $x_e(t) = A_1 u(t)$ are as follows:

$$F_{T}(t) = K_{r}A_{1}u(t) + B_{r}A_{1}\delta(t) + \int_{0}^{t} k_{dyn-H}(\tau) x(t-\tau)d\tau$$

$$= K_{r}A_{1}u(t) + B_{r}A_{1}\delta(t) + \frac{\lambda A_{1}}{\omega_{d}}e^{-at}(asin\omega_{d}t + \omega_{d}cos\omega_{d}t)$$

$$- \lambda A_{1}\frac{2\varsigma_{2}\omega_{n2} - \kappa}{\omega_{d}\omega_{n2}} \{e^{-at}[asin(\omega_{d}t - \theta) + \omega_{d}cos(\omega_{d}t - \theta)] + asin\theta - \omega_{d}cos\theta\},$$

$$P_{1}(t) = \frac{A_{1}K_{1}}{\omega_{d}A_{p}}e^{-at}(asin\omega_{d}t + \omega_{d}cos\omega_{d}t)$$

$$+ \frac{A_{1}K_{1}}{A_{p}}\frac{2\varsigma_{2}\omega_{n2} - \kappa}{\omega_{d}\omega_{n2}} \{e^{-at}[asin(\omega_{d}t - \theta) + \omega_{d}cos(\omega_{d}t - \theta)] + asin\theta - \omega_{d}cos\theta\}.$$

$$(78)$$

5. Effect of different parameters on hydraulic mount

5.1. Compliance of the upper chamber and the viscosity of the fluid

The upper chamber compliance C_1 of model I was varied in order to study the effect of the variation of C_1 on the hydraulic mount. The value of C_1 varies from $0.1C_1$ to $10C_1$, and the variation of the hydraulic mount transfer force is shown in Fig. 7(a). When changing the compliance of the hydraulic mount upper chamber of model I to $0.1 C_1$, the amplitude of the transfer force increases significantly and the period decreases, as shown in the solid blue line in

Fig. 7(a). Because according to Table 1 and Eq. (59), a decrease in C_1 leads to an increase in λ , ω_{n2} . Similarly, when C_1 increases from $1.0C_1$ to $10C_1$ leads to a decrease in λ and ω_{n2} , so the amplitude of the transfer force of the hydraulic mount of model 1 decreases significantly and the period increases for the same reason. When the viscosity vis is increased from 0.1vis to 10vis, it is clear from Fig. 7(b) that the transfer force of the hydraulic mount of model I gradually decreases, especially when increasing the fluid viscosity to 10vis is the most obvious. This is due to the increase in viscosity vis leads to an increase in κ and ς_2 obtained from Table 1, and the decrease in the transfer force $F_T(t)$ of the hydraulic mount of model I according to Eq. (59).



Fig. 7. Effect of upper chamber compliance and fluid viscosity on hydraulic mount transfer force: a) different C_1 values, b) different fluid *vis*

5.2. Effect of inertia channel and orifice flow channel parameters on hydraulic mount

The effect of inertia and damping of the Z_4 structural inertia channel fluid of model II on the transfer force of the hydraulic mount under ideal step excitation is shown in Fig. 8(a), (b). As can be seen from Eq. (4), the value of I_i is regulated by varying the length L and effective cross-sectional area A_i of the inertia channel, which has a significant effect on the peak and period of oscillation of the hydraulic mount under external excitation. Fig. 8(a) shows that as I_i increases, the period of oscillation $t_d = 2\pi/\omega_d$ and the peak value of the transfer force increase. This is due to the combination of Table 1 and Eq. (75) increasing I_i leads to a decrease in the rate of decay of the transfer force. As shown in Fig. 8(b), the step response of the hydraulic mount transfer force changes from an under-damped response to an over-damped response as the value of R_i changes. At the same time, the increase in R_i reduces the amplitude of the system oscillation, but the oscillation period T_d increases slightly. Fig. 8(c), (d) shows the step response of the hydraulic mount transfer force by varying the damping values of the short flow channel of the Z_3 and Z_4 structures for the Model II configuration. Increasing the damping values R_{01} and R_{02} of the orifice flow channels Z_3 and Z_4 , the transfer force $F_T(t)$ oscillation increases. Instead, as the damping values R_{01} and R_{02} decrease, the transfer force $F_T(t)$ decays more rapidly. However, comparing Fig. 8(c) with Fig. 8(d) shows that the larger the cross-sectional area of the orifice flow channel the more rapidly the transfer force of the system decays.

5.3. Effect of different structural configurations on transfer force

In order to compare the effects of different combinations of inertia channels and orifice flow channels on the performance of hydraulic mounts. Fig. 9 shows the variation of the transferred force for Z_2 in model 1 configuration, Z_4 in model 2 configuration and Z_6 in model 3 configuration for the same step excitation $x_e(t) = A_1 u(t)$ and sinusoidal excitation $x_e(t) = A_1 \sin(\omega t)$. It can be observed in Fig. 9 that the Z_4 structure in the model II configuration

has the fastest decay of transfer force, followed by the Z_6 structure in the model III configuration, and the Z_2 structure in the model I configuration has the slowest decay of transfer force. As a result, it can be concluded that the presence of the orifice flow channel increases the damping of the system and minimizes the overshoot of the transfer force step response. Meanwhile, comparing the Z_2 and Z_6 structures, it can be found that although the Z_6 structure decays faster than the Z_2 structure, the oscillation period T_d and the amount of overshoot of the Z_6 structure are significantly larger than those of the Z_2 structure.



Fig. 8. Effect of model II hydraulic mount transfer force under ideal step excitation: a) inertia channel inertia; b) inertia channel damping; c) Z₃ structural orifice flow channel damping;
 d) Z₄ structural orifice flow channel damping



Fig. 9. Combination of hydraulic mounts with different inertia channels to transfer forces: a) step excitation; b) half sine excitation

6. Time-domain response analysis

6.1. Steady-state harmonic excitation response analysis

The blue line in Fig. 10 compares the response of the hydraulic mount transfer force and upper

chamber pressure using the analytical solution of the steady-state response for Models I, II, and III solved in Section 3 with the red dashed line modeled using AMEsim. Fig. 10 show that the steady-state responses of the transfer force and upper chamber pressure solved using the analytical formula are in good agreement with the steady-state responses of the hydraulic mount solution built by AMEsim. For example, the errors between the RMS of the transfer force and the upper chamber pressure response solved by the two methods for the Z_1 the structure are 4.92 % and 9.81 %, respectively. The RMS of the transfer force and the pressure in the upper chamber of the Z_4 structure was 97.4724 N and 0.0178 MPa, respectively, when solved analytically, while the transfer force and the pressure in the upper chamber were 91.03 N and 0.01529 MPa when solved by AMEsim, with RMS errors of 6.61 % and 14.1 %. Similarly, the errors between the RMS of the transfer force and the upper chamber pressure response solved by the two methods for the Z_5 structure are 6.08 % and 10.81 %, respectively.



Fig. 10. Steady-state response of transfer force and upper chamber pressure for Z_1 , Z_4 and Z_5 structures at 25 Hz, amplitude of 0.3 mm: a) Z_1 structure; b) Z_4 structure, c) Z_5 structure

6.2. Ideal step excitation response analysis

The transfer forces for the structure of Model 1 configuration at amplitudes $A_1 = 1.0$ mm, 2.0 mm, 3.0 mm are normalized and then overlapped in the same plot as shown in Fig. 11(a). From Fig. 11(a), it is observed that the overshoot of the transfer force decreases with the increase of the excitation amplitude and the transfer force decays faster while the oscillation period is longer. This means that as the excitation amplitude increases, the natural frequency and damping ratio of the system changes. To further illustrate whether the hydraulic mounts in the inertial channel configuration are damping related when the excitation amplitude varies leading to changes in the transfer force. The transfer force step response solved using the model 1 configuration Z_1 structure in Section III is shown in Fig. 11(b) by changing the damping value of the inertia channel from $1.0R_i$ to an increased value of $3.0R_i$. Comparing Fig. 11(a), (b), it can be found that the damping of the inertial channel depends directly on the excitation amplitude. This conclusion is similar to that of reference studies on multi-inertia channel bushings [14].



Fig. 11. Transfer force response of Z_1 structure with different step excitation for model I configuration: a) Normalized response of transfer force with different excitation amplitude under AMEsim; b) Transfer force response with different damping values solved analytically

7. Conclusions

The main contribution of the article is to investigate the time domain response of hydraulic mounts for the combination of long and short inertia channels and orifice flow channels since there are few studies on the time domain of hydraulic mounts for the combination of long and short inertia channels and orifice flow channels. A combination of numerical analysis and hydraulic analysis software AMEsim was used to compare the steady-state response and step response of different combinations of short and long inertia channels and orifice flow channels. Firstly, the transfer functions of dynamic stiffness and an upper chamber pressure of the hydraulic mounts of six configurations are derived using the lumped model. Secondly, the convolution method is used to solve the time domain analytic equations for the steady-state and step responses of the hydraulic mounts conFig.d in Model I, Model II, and Model III. Study shows that (1) increased compliance and fluid viscosity in the upper chamber of the hydraulic mount can rapidly decrease the transfer force; (2) Reducing the inertia of the inertia channel of the hydraulic mount and increasing the damping of the inertia channel can make the fastest transfer force reach the minimum value while increasing the cross-sectional area of the orifice flow channel can improve the damping value of the hydraulic mount. (3) The transfer forces and upper chamber pressures of the hydraulic mounts of different configurations obtained by the analytical method are in good agreement with the steady-state response of the mounts built by the hydraulic software AMEsim; (4) The effective damping is dependent on the excitation amplitude for hydraulic mounts.

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Data availability

The datasets generated during and/or analyzed during the current study are available from the corresponding author on reasonable request.

Author contributions

Zhihong Lin – conceptualization, methodology, software, investigation, formal analysis, writing-original draft. Chunrong Wang – conceptualization, funding acquisition, writing-review and editing. Yunxiao Chen – software, validation. Mingzhong Wu – visualization, writing-review and editing. Feijie Zheng – software, validation. Jian Wei – visualization, writing-review and editing.

Conflict of interest

The authors declare that they have no conflict of interest.

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