

A simple harmonic quantum oscillator: fractionalization and solution

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Abstract. A quantum mechanical system that mimics the behavior of a classical harmonic oscillator in the quantum domain is called a simple harmonic quantum oscillator. The time-independent Schrödinger equation describes the quantum harmonic oscillator, and its eigenstates are quantized energy values that correspond to various energy levels. In this work, we first fractionalize the time-independent Schrödinger equation, and then we solve the generated problem with the use of the Adomian decomposition approach. It has been shown that fractional quantum harmonic oscillators can be handled effectively using the proposed approach, and their behavior can then be better understood. The effectiveness of the method is validated by a number of numerical comparisons.

Keywords: harmonic oscillator, Adomian decomposition, Hermite polynomial.

1. Introduction

In recent years, fractional calculus – a branch of classical calculus that deals with instructions on non-integer integration and differentiation – has become a fascinating area of study. The notion of fractional operators, which emerged nearly concurrently with their classical equivalents, has garnered considerable attention owing to their multifarious applicability across multiple fields of mathematics and science. It has been demonstrated that fractional differential equations are effective instruments for understanding and simulating intricate engineering, chemical, and physical processes [1-5].

The Schrödinger equation is a key tool in the field of quantum mechanics for explaining how quantum systems behave. This equation, which was created by Erwin Schrödinger in 1926, uses time as an independent variable to describe quantum systems. Time-independent Schrödinger equations (TISE) and time-dependent Schrödinger equations (TDSE) are the two categories into which Schrödinger wave equations fall [5, 6].

Numerous authors have offered both analytical and numerical solutions for TISE and TDSE. The Schrödinger issue has been solved using modern analytical techniques such the Elzaki decomposition methodology, homotopy perturbation, and Adomian decomposition methods. On the other hand, techniques such as the Numerov algorithm for harmonic and linear potentials have been used to produce approximative numerical solutions.

In this work, we study the semi-analytical solution of the Adomian decomposition approach for time-independent Schrödinger equations. The ADM provides an accurate solution and is well-known for its effectiveness in solving Sturm-Liouville issues, ordinary and partial differential equations, nonlinear and stochastic problems, and more. When compared to other approaches, the ADM solves the TISE more quickly and accurately while still offering a series solution [7-10].

The structure of this article is as follows: In part 2, we establish the notations and present the essential concepts, including the fractional operators of the Riemann-Liouville integral and derivative, and the Caputo derivative, which will be utilized throughout this essay. Section 3,

employs the ADM to solve the Schrödinger equation for the simple fractional harmonic oscillator, converting it into a Hermite differential equation. In Section 4, we depict some numerical comparisons between the ADM's solution and the solution reported in some references. Finally, Section 5 summarizes the conclusion of this work.

2. Preliminaries and notion

In this section, we go over the basic ideas of fractional calculus in detail, which forms the basis of our investigation. We also explore the properties of certain operators that are important to our work [11-14].

Definition 2.1. The Riemann-Liouville fractional integral of f of order γ , where $t > 0$, $n - 1 < \gamma \leq n$, $n \in \mathbb{N}$, of a function $f(t)$ is given as follows:

$$J^\gamma f(t) = \frac{1}{\Gamma(\gamma)} \int_0^t (t - \tau)^{\gamma-1} f(\tau) d\tau. \quad (1)$$

The following features of the Riemann-Liouville integral are worth noting:

$$J^0 f(t) = f(t), \quad (2)$$

$$J^\gamma t^m = \frac{\Gamma(m+1)}{\Gamma(m+\gamma+1)} t^{m+\gamma}, \quad m \geq -1, \quad (3)$$

$$J^\gamma J^\beta f(t) = J^{\beta+\gamma} f(t), \quad \gamma, \beta \geq 0, \quad (4)$$

$$J^\gamma J^\beta f(t) = J^{\gamma+\beta} f(t), \quad \gamma, \beta \geq 0. \quad (5)$$

Definition 2.2. Let $n - 1 < \gamma < n$, such that n is positive integer and $\gamma \in \mathbb{R}^+$, the Riemann Liouville derivative of fractional of order γ is given as follows:

$$D^\gamma f(t) = \frac{1}{\Gamma(n-\gamma)} \frac{d^n}{dt^n} \int_0^t \frac{f^{(n)}(\tau)}{(t-\tau)^{\gamma+1-n}} d\tau. \quad (6)$$

Definition 2.3. Suppose that $\gamma > 0$, $t > \gamma$, $t \in \mathbb{R}$. The Caputo fractional differential operator of order γ , is given as:

$$D^\gamma f(x) = \begin{cases} \frac{1}{\Gamma(n-\gamma)} \int_0^t \frac{f^{(n)}(\tau)}{(t-\tau)^{\gamma+1-n}} d\tau, & n-1 < \gamma < n \in \mathbb{N}, \\ \frac{d^n}{dt^n} f(t), & \gamma = n \in \mathbb{N}. \end{cases} \quad (7)$$

Remarks. The Caputo fractional derivative meets the following properties:

1. The power rule property, i.e.:

$$D^\gamma t^\varepsilon = \begin{cases} \frac{\Gamma(\varepsilon+1)}{\Gamma(\varepsilon-\gamma+1)} t^{\varepsilon-\gamma} = D^\gamma t^\varepsilon, & n-1 < \alpha < n, \quad \varepsilon > n-1, \quad \varepsilon \in \mathbb{R}, \\ 0, & n-1 < \alpha < n, \quad \varepsilon \leq n-1, \quad \varepsilon \in \mathbb{N}. \end{cases} \quad (8)$$

2. The constant property, i.e.:

$$D^\gamma c = 0. \quad (9)$$

3. The interpolation property, i.e.:

$$\lim_{\gamma \rightarrow n} D^\gamma f(t) = D^n f(t). \quad (10)$$

4. The linearity property, i.e.:

$$D^\gamma(\lambda f(t) + \psi g(t)) = \lambda D^\gamma f(t) + D^\gamma g(t)\psi. \quad (11)$$

5. The non-commutation property, i.e.:

$$D^\gamma D^\gamma f(t) = D^{\gamma+\gamma} f(t) \neq D^\gamma D^\gamma f(t). \quad (12)$$

3. Theory

One typical method for modeling systems with viscoelasticity in damping terms is to incorporate fractional derivatives to introduce memory effects in the system. The mathematical framework offered by fractional calculus makes it possible to include memory effects and produce a more realistic depiction of the behavior of the system.

The damping term in the differential equation can be written as a fractional derivative of the displacement or velocity variable when modeling viscoelasticity damping with fractional derivatives. The memory effects in the damping system are captured by this fractional derivative, enabling a more accurate depiction of the system's reaction to outside influences.

The core of this study is covered in this section, where we introduce our unique method for solving the Schrödinger equation for the basic fractional harmonic oscillator. In order to do this, we make use of the ADM, a reliable numerical technique that is well-known for its accuracy and efficiency in solving differential equations. Our objective is to transform the original Schrödinger equation into the Hermite differential equation, which is a more manageable form and has fractional order derivatives [15-17].

An important model for studying quantum systems with fractional characteristics is the simple fractional harmonic oscillator. Our goal is to achieve precise answers by utilizing the ADM, which will provide insight into the fascinating occurrences seen in the equation under consideration. The ADM is a perfect fit for this study because of its adaptability in handling nonlinearity and the complexities of fractional calculus [18-20].

In this context, the expression for the time-dependent Schrodinger's equation for a basic fractional harmonic oscillator is:

$$-\frac{\hbar^2}{2r} \frac{d^{2\alpha} \phi(x)}{dx^{2\alpha}} + \frac{1}{2} kx^2 \phi(x) = T \phi(x), \quad (13)$$

where h , r are constants, T is the energy. It should be observed here that the above model corresponds to inertia terms in the second-order mechanistic counterpart model when $\alpha = 1$. In the same regard, we should also note that since we are dealing with a physical problem, we will not be considering the positive exponent solution that diverges at. It was noted, meanwhile, that Eq. (13) can be changed to the following form:

$$\frac{d^{2\alpha} U}{d\gamma^2} - 2\gamma \frac{dU}{d\gamma} + 2 \left(\frac{T}{\hbar\omega} - \frac{1}{2} \right) U = 0, \quad (14)$$

where $\phi(\gamma) = U(\gamma)e^{-(\gamma^2/2)}$. Now define $\lambda = \left(\frac{T}{\hbar\omega} - \frac{1}{2} \right)$ and replace the dummy γ with x . This reduces Eq. (14) to the form shown below:

$$D^{2\alpha} U - 2xD^\alpha U + 2\lambda U = 0, \quad (15)$$

where $U(0) = c_0$, $D^\alpha U'(0) = c_1$.

The ADM is used to tackle this problem, by considering the following form:

$$D^{2\alpha}U = 2xD^\alpha U - 2\lambda U. \quad (16)$$

Operating J^α to both sides of Eq. (16) yields:

$$D^\alpha U - D^\alpha U'(0) = J^\alpha(2xD^\alpha U) - 2\lambda J^\alpha U,$$

or:

$$D^\alpha U = c_1 + 2J^\alpha(xD^\alpha U) - 2\lambda J^\alpha U. \quad (17)$$

Again, by taking J^α to both sides of Eq. (17), we obtain:

$$U - U(0) = J^\alpha c_1 + 2J^{2\alpha}(xD^\alpha U) - 2\lambda J^{2\alpha} U.$$

The result is as follows:

$$U = c_0 + c_1 \frac{x^\alpha}{\Gamma(\alpha + 1)} + 2J^{2\alpha}(xD^\alpha U) - 2\lambda J^{2\alpha} U, \quad (18)$$

where c_0 and c_1 are arbitrary constants. Following the ADM sense leads to the general solution for U , which would be as:

$$U = \sum_{n=0}^{\infty} U_n, \quad (19)$$

where:

$$U_0 = c_0 + c_1 \frac{x^\alpha}{\Gamma(\alpha + 1)}, \quad U_{n+1} = 2J^{2\alpha}(xD^\alpha U_n) - 2\lambda J^{2\alpha} U_n, \quad n \geq 0. \quad (20)$$

We will show some of the major calculations used above. The following steps show how we calculate U_1 up to U_5 . In this regard, we can find:

$$U_1 = 2J^{2\alpha}(xD^\alpha U_0) - 2\lambda J^{2\alpha} U_0,$$

which means:

$$U_1 = 2J^{2\alpha} \left(xD^\alpha \left(c_0 + c_1 \frac{x^\alpha}{\Gamma(\alpha + 1)} \right) \right) - 2\lambda J^{2\alpha} \left(c_0 + c_1 \frac{x^\alpha}{\Gamma(\alpha + 1)} \right),$$

or:

$$U_1 = \left(\frac{-2\lambda x^{2\alpha}}{\Gamma(2\alpha + 1)} \right) c_0 + \left(\frac{2x^{2\alpha+1}}{\Gamma(2\alpha + 2)} - \frac{2\lambda x^{3\alpha}}{\Gamma(3\alpha + 1)} \right) c_1.$$

The next term can be written as:

$$U_2 = 2J^{2\alpha}(xD^\alpha U_1) - 2\lambda J^{2\alpha}(U_1).$$

In other words, we have:

$$U_2 = -2^2 \lambda c_0 \left(\frac{\Gamma(\alpha + 2)x^{3\alpha+1}}{\Gamma(\alpha + 1)\Gamma(3\alpha + 2)} - \frac{x^{4\alpha}}{\Gamma(4\alpha + 1)} \right) + 2^2 c_1 \left(\frac{\Gamma(\alpha + 3)x^{3\alpha+2}}{\Gamma(\alpha + 2)\Gamma(3\alpha + 3)} - \lambda \frac{x^{4\alpha}}{\Gamma(4\alpha + 2)} \left(1 + \frac{\Gamma(2\alpha + 2)}{\Gamma(2\alpha + 1)} \right) + \frac{\lambda^2 x^{5\alpha}}{\Gamma(5\alpha + 1)} \right).$$

Once more, the next term can be written as:

$$U_3 = 2J^{2\alpha}(xD^\alpha U_2) - 2\lambda J^{2\alpha}(U_2).$$

This is equivalent to the following assertion:

$$U_3 = -2^3 c_0 \left(\frac{\lambda \Gamma(\alpha + 2)\Gamma(2\alpha + 3)x^{4\alpha+2}}{\Gamma(\alpha + 1)\Gamma(2\alpha + 2)\Gamma(4\alpha + 3)} + \lambda^2 x^{5\alpha+1} \left(\frac{\Gamma(3\alpha + 2)}{\Gamma(3\alpha + 1)\Gamma(5\alpha + 2)} + \frac{\Gamma(\alpha + 2)}{\Gamma(\alpha + 1)\Gamma(5\alpha + 2)} \right) - \frac{\lambda^3 x^{6\alpha}}{\Gamma(6\alpha + 1)} \right) + 2^3 c_1 \left(\frac{\Gamma(\alpha + 3)\Gamma(2\alpha + 4)x^{3\alpha+2}}{\Gamma(\alpha + 2)\Gamma(2\alpha + 3)\Gamma(4\alpha + 4)} - \lambda^2 x^{5\alpha+2} \left(\frac{\Gamma(3\alpha + 3)}{\Gamma(3\alpha + 2)\Gamma(5\alpha + 3)} \left(1 + \frac{\Gamma(2\alpha + 2)}{\Gamma(2\alpha + 1)} \right) + \frac{\Gamma(\alpha + 3)}{\Gamma(\alpha + 2)\Gamma(5\alpha + 3)} \right) + \frac{\lambda^2 x^{6\alpha+1}}{\Gamma(6\alpha + 2)} \left(\frac{\Gamma(4\alpha + 2)}{\Gamma(4\alpha + 1)} - \left(1 + \frac{\Gamma(2\alpha + 2)}{\Gamma(2\alpha + 1)} \right) \right) - \frac{\lambda^3 x^{7\alpha}}{\Gamma(7\alpha + 1)} \right).$$

Similarly, we can have:

$$U_4 = 2J^{2\alpha}(xD^\alpha U_3) - 2\lambda J^{2\alpha}(U_3),$$

which consequently gives:

$$U_4 = \frac{-2^4 \lambda c_0 \Gamma(\alpha + 2)\Gamma(2\alpha + 3)\Gamma(3\alpha + 4)x^{5\alpha+3}}{\Gamma(\alpha + 1)\Gamma(2\alpha + 2)\Gamma(3\alpha + 3)\Gamma(5\alpha + 4)} + 2^4 c_1 \frac{\Gamma(\alpha + 3)\Gamma(2\alpha + 4)\Gamma(3\alpha + 5)x^{5\alpha+4}}{\Gamma(\alpha + 2)\Gamma(2\alpha + 3)\Gamma(3\alpha + 4)\Gamma(5\alpha + 5)} + \frac{2^4 \lambda^2 c_0 x^{6\alpha+2}}{\Gamma(6\alpha + 3)} \left(\left(\frac{\Gamma(\alpha + 2)}{\Gamma(\alpha + 1)} \left(\frac{\Gamma(4\alpha + 3)}{\Gamma(4\alpha + 2)} + \frac{\Gamma(2\alpha + 3)}{\Gamma(2\alpha + 2)} \right) \right) + \frac{\Gamma(3\alpha + 2)\Gamma(4\alpha + 3)}{\Gamma(3\alpha + 1)\Gamma(4\alpha + 2)} \right) - \frac{2^4 \lambda c_1 x^{6\alpha+3}}{\Gamma(6\alpha + 4)} \left(\left(1 + \frac{\Gamma(2\alpha + 2)}{\Gamma(2\alpha + 1)} \right) \frac{\Gamma(3\alpha + 3)\Gamma(4\alpha + 4)}{\Gamma(4\alpha + 3)\Gamma(3\alpha + 2)} + \frac{\Gamma(\alpha + 3)\Gamma(4\alpha + 4)}{\Gamma(\alpha + 2)\Gamma(4\alpha + 3)} + \frac{\Gamma(\alpha + 3)\Gamma(2\alpha + 4)}{\Gamma(2\alpha + 3)\Gamma(\alpha + 2)} \right) - \frac{2^4 \lambda^3 c_0 x^{7\alpha+1}}{\Gamma(7\alpha + 2)} \left(\frac{\Gamma(5\alpha + 2)}{\Gamma(5\alpha + 1)} + \frac{\Gamma(3\alpha + 2)}{\Gamma(3\alpha + 1)} + \frac{\Gamma(\alpha + 2)}{\Gamma(\alpha + 1)} \right) + \frac{2^4 \lambda^2 c_1 x^{7\alpha+2}}{\Gamma(7\alpha + 3)} \left(\left(1 + \frac{\Gamma(2\alpha + 2)}{\Gamma(2\alpha + 1)} \right) \left(\frac{\Gamma(3\alpha + 3)}{\Gamma(3\alpha + 2)} + \frac{\Gamma(5\alpha + 3)}{\Gamma(5\alpha + 2)} \right) + \frac{\Gamma(\alpha + 3)}{\Gamma(\alpha + 2)} + \frac{\Gamma(4\alpha + 2)\Gamma(5\alpha + 3)}{\Gamma(4\alpha + 1)\Gamma(5\alpha + 2)} \right) + \frac{2^4 \lambda^4 c_0 x^{8\alpha}}{\Gamma(8\alpha + 1)} + \frac{2^4 \lambda^3 c_1 x^{8\alpha+1}}{\Gamma(8\alpha + 2)} \left(\left(1 + \frac{\Gamma(2\alpha + 2)}{\Gamma(2\alpha + 1)} \right) + \frac{\Gamma(4\alpha + 2)}{\Gamma(4\alpha + 1)} + \frac{\Gamma(6\alpha + 2)}{\Gamma(6\alpha + 1)} \right) + \frac{2^4 \lambda^4 c_1 x^{9\alpha}}{\Gamma(9\alpha + 1)} \tag{21}$$

If we continue in this manner, we got the desired solution reported in Eq. (19). To show the validity of our solution, we compare the ADM's solution with the solution reported in [20] in Fig. 1 and Fig. 2 by taking $\alpha = 1$ and $\lambda = \{0.5, 1.5\}$. Obviously, we can observe that the proposed solution coincides with the solution we compare, and this confirms the validity of the proposed solution.

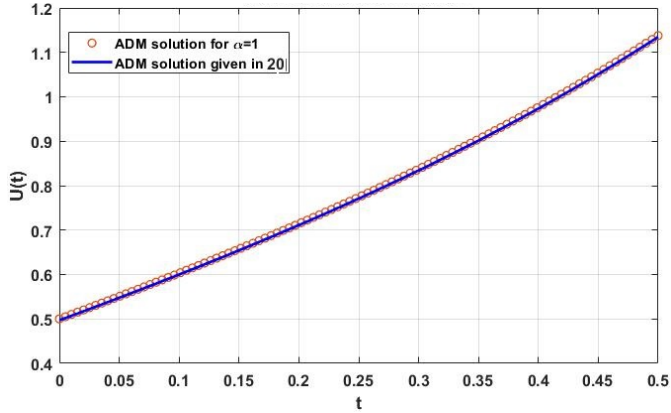


Fig. 1. ADM's solution vs. the solution in [20] according to $\lambda = 0.5$

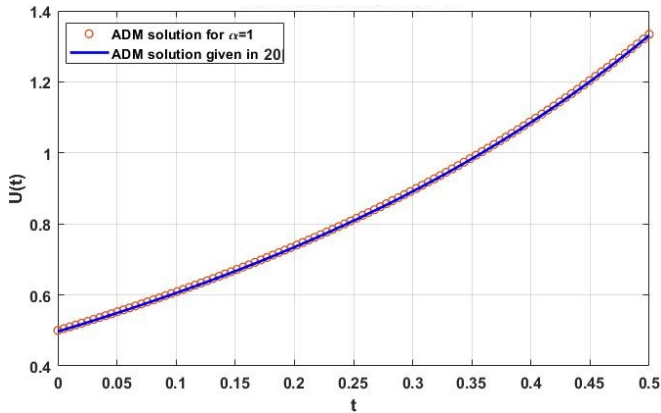


Fig. 2. ADM's solution vs. the solution in [20] according to $\lambda = 1.5$

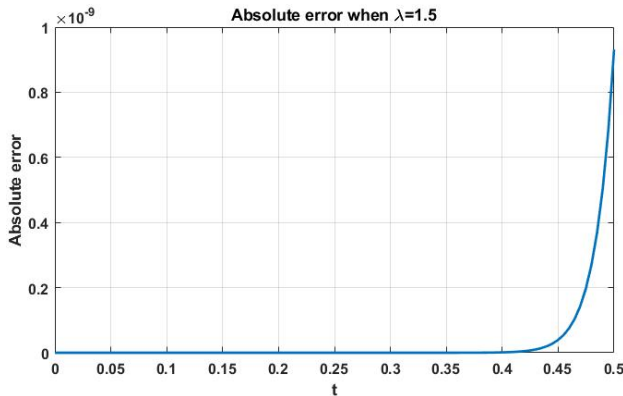


Fig. 3. Absolute error between the ADM's solution and the solution in [20]

For more description, we plot in Fig. 3 and Fig. 4, the absolute errors gained from the

performance comparison. These figures demonstrate the advantage of the suggested method of fractionalizing the model in question. However, the limitation of our proposed approach lies in the lack of capabilities to obtain a general form of U_n that can generate all other components of the desired solution, and this will be left to the future for more consideration.

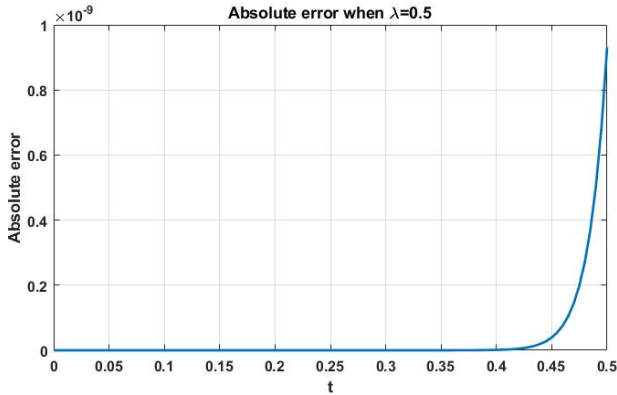


Fig. 4. Absolute error between the ADM's solution and the solution in [20] when $\lambda = 0.5$

4. Conclusions

In this study, a generalized form of the time-independent Schrödinger equation has been successfully proposed in its fractional simple harmonic quantum oscillator. The ADM has been used to obtain a semi-analytical solution for such an oscillator. This solution has proved its validity with an existing solution in [20] when a classical case of the generalized form is taken into consideration.

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Data availability

The datasets generated during and/or analyzed during the current study are available from the corresponding author on reasonable request.

Author contributions

Conceptualization, Iqbal M. Batiha; methodology, Iqbal H. Jebril and Abeer A. Al-Nana; validation, Shameseddin Alshorm; formal analysis, Iqbal M. Batiha.; investigation, Iqbal H. Jebril and Abeer A. Al-Nana; resources, Shameseddin Alshorm; data curation, Iqbal M. Batiha; writing-original draft, Abeer A. Al-Nana; visualization, Iqbal H. Jebril; supervision, Iqbal M. Batiha. All authors have read and agreed to the published version of the manuscript.

Conflict of interest

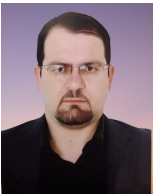
The authors declare that they have no conflict of interest.

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