

1D manipulator with vibration impact drive, based on which it is possible to create orthogonal manipulators and robots of any dimension

K. Ragulskis¹, L. Ragulskis²

¹Kaunas University of Technology, K. Donelaičio str. 73, LT-44249, Kaunas, Lithuania

²Department of Systems Analysis, Faculty of Informatics, Vytautas Magnus University, Vileikos str. 8, LT-44404, Kaunas, Lithuania

²Corresponding author

E-mail: ¹kazimieras3@hotmail.com, ²l.ragulskis@if.vdu.lt

Received 12 January 2024; accepted 8 March 2024; published online 17 March 2024
DOI <https://doi.org/10.21595/amr.2024.23930>



Copyright © 2024 K. Ragulskis, et al. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.

Abstract. Manipulator of the investigated type may move according to a straight line. It has an advantage in the fact that by choosing geometrical parameters of the manipulator it is possible to achieve its effective operation. This is presented by using analytical and graphical methods. The performed research shows that manipulators with vibration impact drives have some positive qualities. In their structure it is not necessary to include the self-stopping mechanism. In the conservative case of the system static position of equilibrium of the impact pair can be with negative, zero or positive tightening. In the case of zero tightening eigenfrequency and period of the system does not depend on the quantity of motion of impact excitation. In the case of harmonic forced excitation resonant motions take place in the vicinity of the eigenfrequency of the conservative system with zero tightening. Analytical – numerical calculations contribute to the creation of manipulators and robots with vibration impact drives.

Keywords: vibration impact drive, elements of robots, manipulators, graphical representation, nonlinear dynamics.

1. Introduction

Investigation of dynamics of manipulators and robots with vibration impact drive is important in engineering. Manipulator of the investigated type may move according to a straight line. It has an advantage in the fact that by choosing geometrical parameters of the manipulator it is possible to achieve its effective operation. This is presented by using analytical and graphical methods.

The model of the manipulator is described. First, the dynamics of the conservative system is investigated. Results for typical parameters of the manipulator were obtained. Then, the dynamics of the model with forced excitation for various parameters of the manipulator is investigated. The obtained graphical representations show the main dynamic regimes of the manipulator.

The investigation is performed by using analytical – numerical methods of analysis of nonlinear systems.

Resonances of vibrating systems are analyzed in [1]. Dynamics in vibrating systems for various excitations is investigated in [2]. Stabilization of nonlinear systems is described in [3]. Nonlinear effects in vibrating systems are analyzed in [4]. Investigation of periodic orbits is performed in [5]. Analysis of energy transfer is presented in [6]. Investigation of nonlinear interactions is described in [7]. Frequencies of vibrating systems are analyzed in [8]. Investigation of the pendulum mechanism is performed in [9]. Analysis of a piecewise linear system is presented in [10]. Resonant zones of nonlinear systems are described in [11]. Investigation of the Sommerfeld effect is performed in [12]. Analysis of isolated resonances is described in [13]. New types of robots are investigated in [14]. Analysis of the positioning system is performed in [15]. Investigation of the automatic positioning stage is presented in [16]. Analysis of the piezo-driven stage is performed in [17]. Investigation of the precision positioning stage is described in [18].

Analysis of essentially nonlinear dynamics is performed in [19]. Investigation of transmissions is described in [20]. Analysis of robots is performed in [21].

The purpose of the work is to achieve improved operation of the manipulator with vibration impact drive by the choice of structural, geometrical, dynamical parameters for effective operation of the manipulator by using analytical and numerical investigations.

In the conservative case of the system static position of equilibrium of the impact pair can be with negative, zero or positive tightening. Zero tightening is advantageous for practical applications and thus the investigations for zero tightening are performed.

Results of the performed investigation are used in the design of manipulators with vibration impact drive.

2. General dynamic model of the manipulator

The manipulator is shown in Fig. 1. In the figure XOY denotes the immovable coordinate system. By the number 1 the mass m_1 is denoted. By the number 2 immovable straight supports are denoted. By the letter C coefficient of stiffness of the spring is denoted. By the letter H coefficient of viscous damping of the damper of vibrations is denoted. By $F(\omega t)$ the exciting force is denoted. By the number 0 the case with the mass m_0 is denoted. Further H_{02} denotes the coefficient of viscous friction between the case and the immovable straight supports. The quantities x_0 and x_1 denote the displacements shown in the figure. Further the upper dot denotes differentiation with respect to the time variable. The force of resistance of external media to the motion of the case is denoted as $-A - B\dot{x}_0$, where A and B are constants and \dot{x}_0 is the velocity of motion of the case.

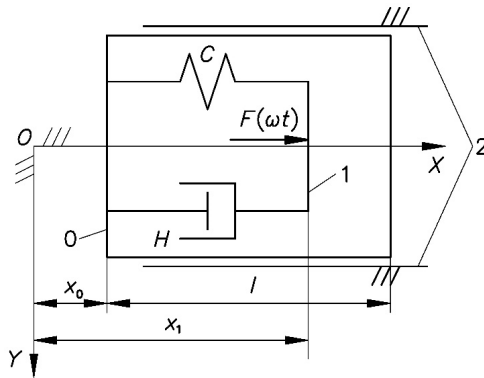


Fig. 1. Model of the system: XOY – the immovable coordinate system; 1 – mass; 2 – immovable straight supports; C – coefficient of stiffness of the spring; H – coefficient of viscous damping of the damper of vibrations; $F(\omega t)$ – the exciting force

In the conservative case of the system static position of equilibrium of the impact pair can be with negative, zero or positive tightening. Zero tightening is advantageous for practical applications and thus the investigations for zero tightening are performed. The distance l is the position of equilibrium of the spring.

When the following condition is satisfied:

$$x_1 - x_0 < l, \quad (1)$$

dynamics of the system is described by the following equations:

$$m_0\ddot{x}_0 + C(x_0 - x_1 + l) + H(\dot{x}_0 - \dot{x}_1) + H_{02}\dot{x}_0 = -A - B\dot{x}_0, \quad (2)$$

$$m_1\ddot{x}_1 + C(x_1 - x_0 - l) + H(\dot{x}_1 - \dot{x}_0) = F\sin\omega t, \quad (3)$$

where F is the amplitude of harmonic excitation and ω is the frequency of harmonic excitation.

When the following condition is satisfied:

$$x_1 - x_0 = l, \quad (4)$$

impact of the mass m_1 with the mass m_0 takes place in a moment of time, where the change of velocities of those masses takes place according to the theorem of quantity of motion and losses of velocity of impact, that is:

$$m_0 \dot{x}_0^+ + m_1 \dot{x}_1^+ = m_0 \dot{x}_0^- + m_1 \dot{x}_1^-, \quad (5)$$

$$R = -\frac{\dot{x}_1^+ - \dot{x}_0^+}{\dot{x}_1^- - \dot{x}_0^-}, \quad (6)$$

where R is the coefficient of restitution of the impact. From those equations \dot{x}_1^+ and \dot{x}_0^+ are found.

The following notations are introduced:

$$\mu = \frac{m_0}{m_1}, \quad p^2 = \frac{C}{m_1}, \quad h = \frac{H}{m_1}, \quad h_{02} = \frac{H_{02}}{m_1}, \quad a = \frac{A}{m_1}, \quad b = \frac{B}{m_1}, \quad f = \frac{F}{m_1}. \quad (7)$$

When the following condition is satisfied:

$$x_1 - x_0 < l, \quad (8)$$

dynamics of the system is described by the following equations:

$$\mu \ddot{x}_0 + p^2(x_0 - x_1 + l) + h(\dot{x}_0 - \dot{x}_1) + h_{02}\dot{x}_0 = -a - b\dot{x}_0, \quad (9)$$

$$\ddot{x}_1 + p^2(x_1 - x_0 - l) + h(\dot{x}_1 - \dot{x}_0) = f \sin \omega t. \quad (10)$$

When the following condition is satisfied:

$$x_1 - x_0 = l, \quad (11)$$

impact takes place, that is:

$$\dot{x}_0^+ = \frac{1}{1 + \mu} [(\mu - R)\dot{x}_0^- + (1 + R)\dot{x}_1^-], \quad (12)$$

$$\dot{x}_1^+ = \frac{1}{1 + \mu} [\mu(1 + R)\dot{x}_0^- + (1 - \mu R)\dot{x}_1^-]. \quad (13)$$

3. Simplified dynamic model with one degree of freedom

First simplified dynamic model of the manipulator with one degree of freedom is investigated.

When the following condition is satisfied:

$$x_1 < l, \quad (14)$$

dynamics of the system is described by the following equation:

$$\ddot{x}_1 + p^2(x_1 - l) + h\dot{x}_1 = f \sin \omega t. \quad (15)$$

When the following condition is satisfied:

$$x_1 = l, \quad (16)$$

impact takes place, that is:

$$\dot{x}_1^+ = -R\dot{x}_1^- \quad (17)$$

3.1. Conservative system

For the conservative system it is assumed that:

$$h = f = 0, \quad R = 1. \quad (18)$$

When the following condition is satisfied:

$$x_1 < l, \quad (19)$$

dynamics of the system is described by the following equation:

$$\ddot{x}_1 + p^2(x_1 - l) = 0. \quad (20)$$

When the following condition is satisfied:

$$x_1 = l, \quad (21)$$

impact takes place, that is:

$$\dot{x}_1^+ = -\dot{x}_1^-. \quad (22)$$

Further investigation of the conservative system is performed.

The following parameters of the system are assumed:

$$l = 1, \quad p = 1. \quad (23)$$

The following initial conditions are assumed:

$$x_1(0) = l, \quad \dot{x}_1(0) = 1. \quad (24)$$

Results for the case of the conservative system are presented in Fig. 2.

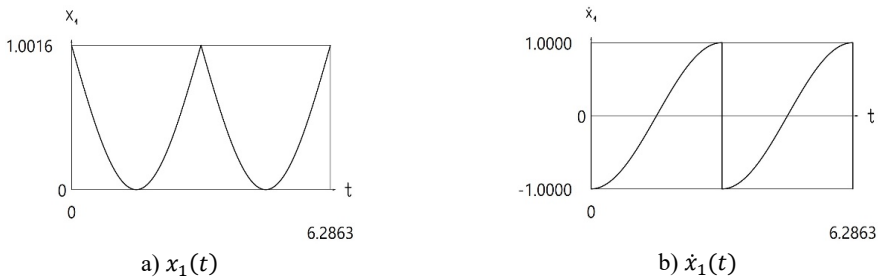


Fig. 2. Dynamics of the one-dimensional model for the case of the conservative system

3.2. Investigation of the dynamics of the system

The following parameters of the system are assumed:

$$\omega = 1, \quad l = 1, \quad h = 0.1, \quad p = 1, \quad f = 6, \quad R = 0.7. \quad (25)$$

The following initial conditions are assumed:

$$x_1(0) = l, \quad \dot{x}_1(0) = 0. \quad (26)$$

Results of investigation of dynamics are presented in Fig. 3.

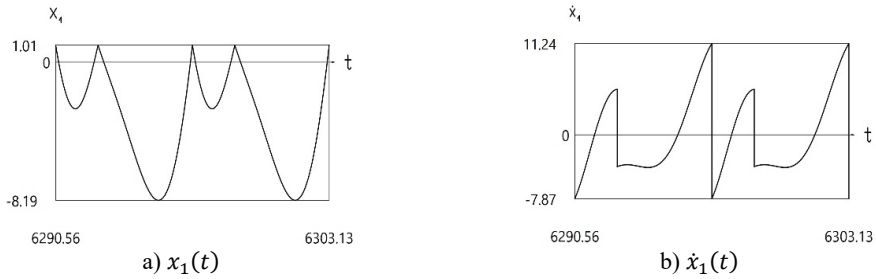


Fig. 3. Dynamics of the one-dimensional model

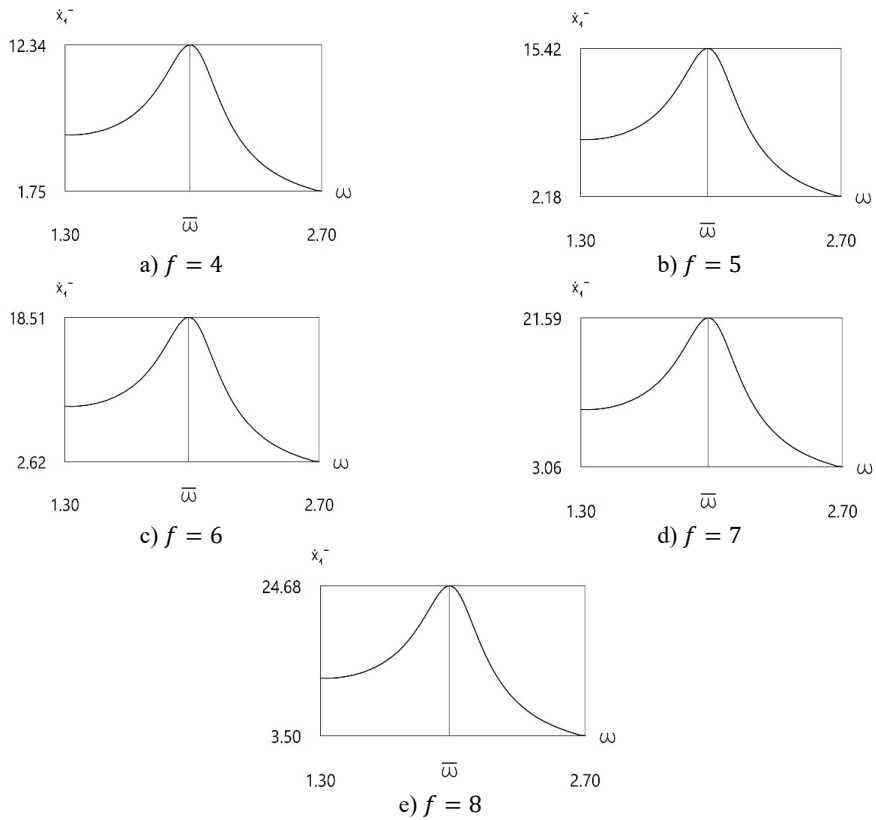


Fig. 4. Velocity before impact as function of frequency of excitation

Further investigation of the velocity before impact as function of frequency of excitation is performed.

The following parameters of the system are assumed:

$$l = 1, \quad h = 0.1, \quad p = 1, \quad R = 0.7. \quad (27)$$

The following initial conditions are assumed:

$$x_1(0) = l, \quad \dot{x}_1(0) = 0. \quad (28)$$

Results for different values of f are presented in Fig. 4.

Results for all values of f are presented in Fig. 5.

From the presented results in the vicinity of the first resonance the first resonant frequency is observed.

Further the full dynamic model of the manipulator is investigated.

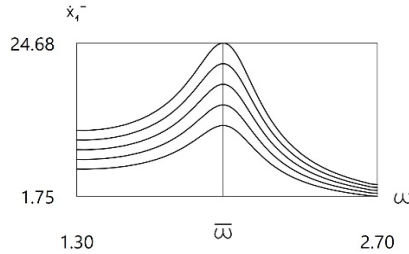


Fig. 5. Velocity before impact as function of frequency of excitation for all values of f

4. Conservative system

For the conservative system it is assumed that:

$$h = h_{02} = a = b = f = 0, \quad R = 1. \quad (29)$$

When the following condition is satisfied:

$$x_1 - x_0 < l, \quad (30)$$

dynamics of the system is described by the following equations:

$$\mu \ddot{x}_0 + p^2(x_0 - x_1 + l) = 0, \quad (31)$$

$$\ddot{x}_1 + p^2(x_1 - x_0 - l) = 0, \quad (32)$$

$$\frac{\mu \ddot{x}_0}{\ddot{x}_1} = -1. \quad (33)$$

When the following condition is satisfied:

$$x_1 - x_0 = l, \quad (34)$$

impact takes place, that is:

$$\dot{x}_0^+ = \frac{1}{1 + \mu} [(\mu - 1)\dot{x}_0^- + 2\dot{x}_1^-], \quad (35)$$

$$\dot{x}_1^+ = \frac{1}{1 + \mu} [2\mu\dot{x}_0^- + (1 - \mu)\dot{x}_1^-]. \quad (36)$$

Investigation of dynamics is performed according to the Eqs. (29-36). Initial conditions for x_0 , \dot{x}_0^+ , x_1 and \dot{x}_1^+ are to be known. It is convenient to use the point of impact, which is described by the Eqs. (34-36). In those equations for example by freely assuming x_1 , \dot{x}_0^- , \dot{x}_1^- the quantities x_1 , \dot{x}_0^+ , \dot{x}_1^+ are found. Further according to the Eqs. (30-33) solution till the next impact according to the Eqs. (34-36) is performed and so on.

Further graphical relationships of the variables are obtained, and their analysis is performed as well as of the quantity of motion $S = \mu\dot{x}_0^- + \dot{x}_1^- = \mu\dot{x}_0^+ + \dot{x}_1^+$ and of the spectral relationships. From the Eq. (33) it follows that for the conservative system the value of S is constant.

Further investigation of the conservative system is performed.

The following parameters of the system are assumed:

$$l = 1, \quad \mu = 5, \quad p = 1. \quad (37)$$

The following initial conditions are assumed:

$$x_0(0) = 0, \quad x_1(0) = l, \quad \dot{x}_0(0) = -0.4, \quad \dot{x}_1(0) = 1. \quad (38)$$

Results for the case of the conservative system are presented in Fig. 6.

From the presented results the period of vibrations of the investigated manipulator is determined.

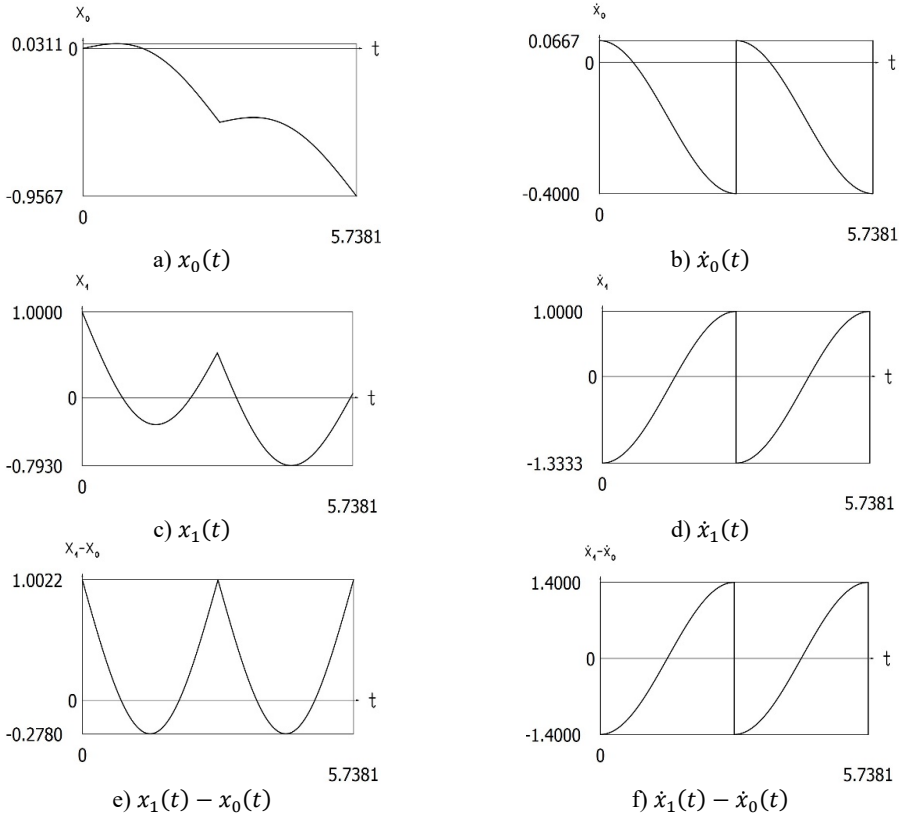


Fig. 6. Dynamics of the manipulator for the case of the conservative system

5. Investigation of the dynamics of the system with external harmonic excitation in the vicinities of resonances

The following parameters of the system are assumed:

$$l = 1, \quad \mu = 5, \quad h = 0.1, \quad h_{02} = 0.1, \quad b = 0.1, \quad p = 1, \quad a = 0.5, \quad f = 6, \quad R = 0.7. \quad (39)$$

The following initial conditions are assumed:

$$x_0(0) = 0, \quad x_1(0) = l, \quad \dot{x}_0(0) = 0, \quad \dot{x}_1(0) = 0. \quad (40)$$

The eigen period was determined as:

$$\bar{T} = 2.868. \quad (41)$$

The eigenfrequency is:

$$\bar{\omega} = \frac{2\pi}{\bar{T}}. \quad (42)$$

It is assumed that:

$$\omega = \bar{\omega} + \delta\omega. \quad (43)$$

Results for different values of $\delta\omega$ are presented in Fig. 7, Fig. 8, and Fig. 9.

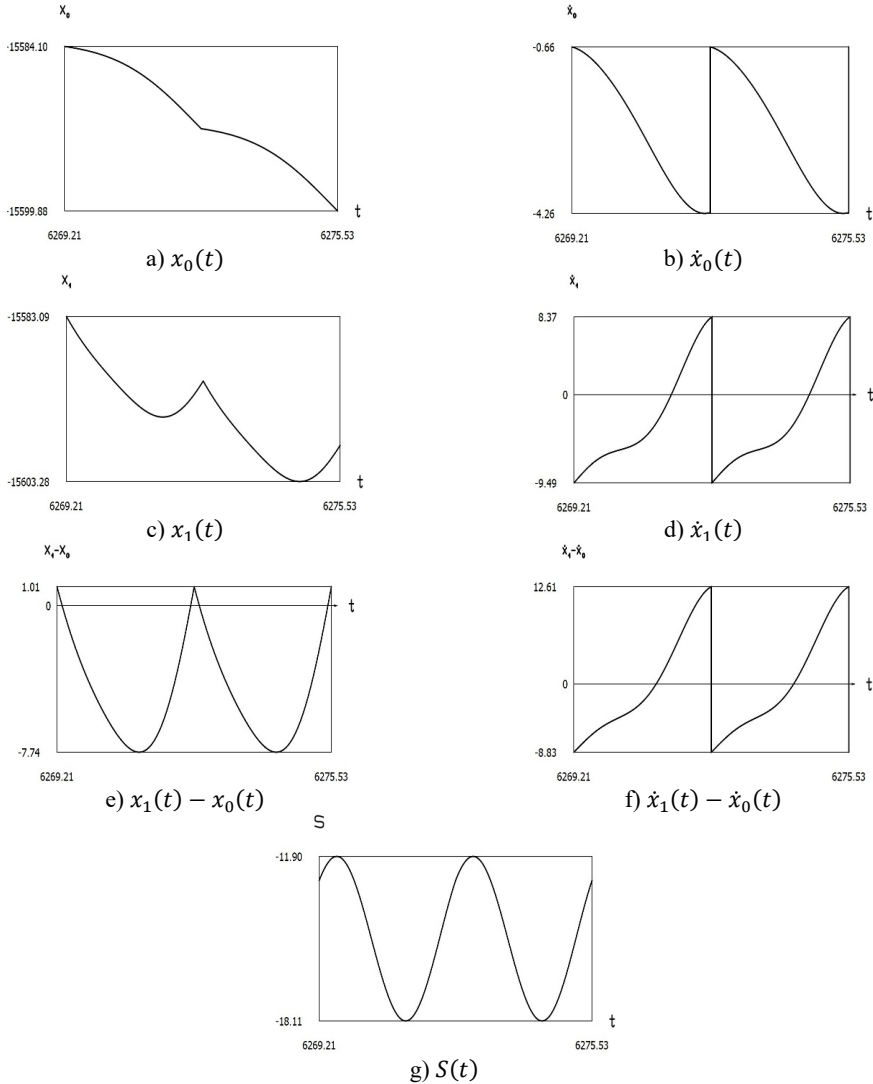


Fig. 7. Dynamics of the manipulator when $\delta\omega = -0.2$

From the presented results below the first resonance, at the first resonance and after the first resonance the effectiveness of operation of the manipulator with vibration impact drive at the first

resonant frequency is observed.

Further investigation of the quantity of motion during impact as function of frequency of excitation is performed.

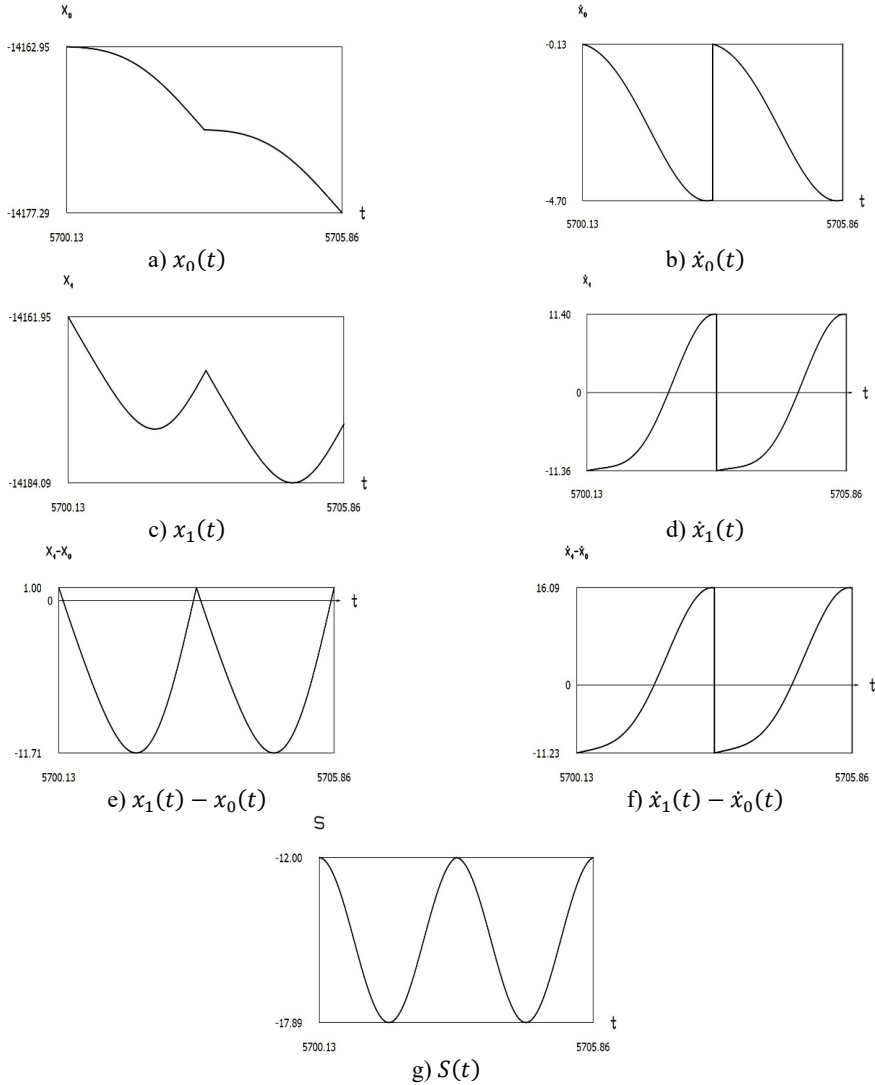


Fig. 8. Dynamics of the manipulator when $\delta\omega = 0$

The following parameters of the system are assumed:

$$l = 1, \quad \mu = 5, \quad h = 0.1, \quad h_{02} = 0.1, \quad b = 0.1, \quad p = 1, \quad a = 0.5, \quad R = 0.7. \quad (44)$$

The following initial conditions are assumed:

$$x_0(0) = 0, \quad x_1(0) = l, \quad \dot{x}_0(0) = 0, \quad \dot{x}_1(0) = 0. \quad (45)$$

Results for different values of f are presented in Fig. 10.

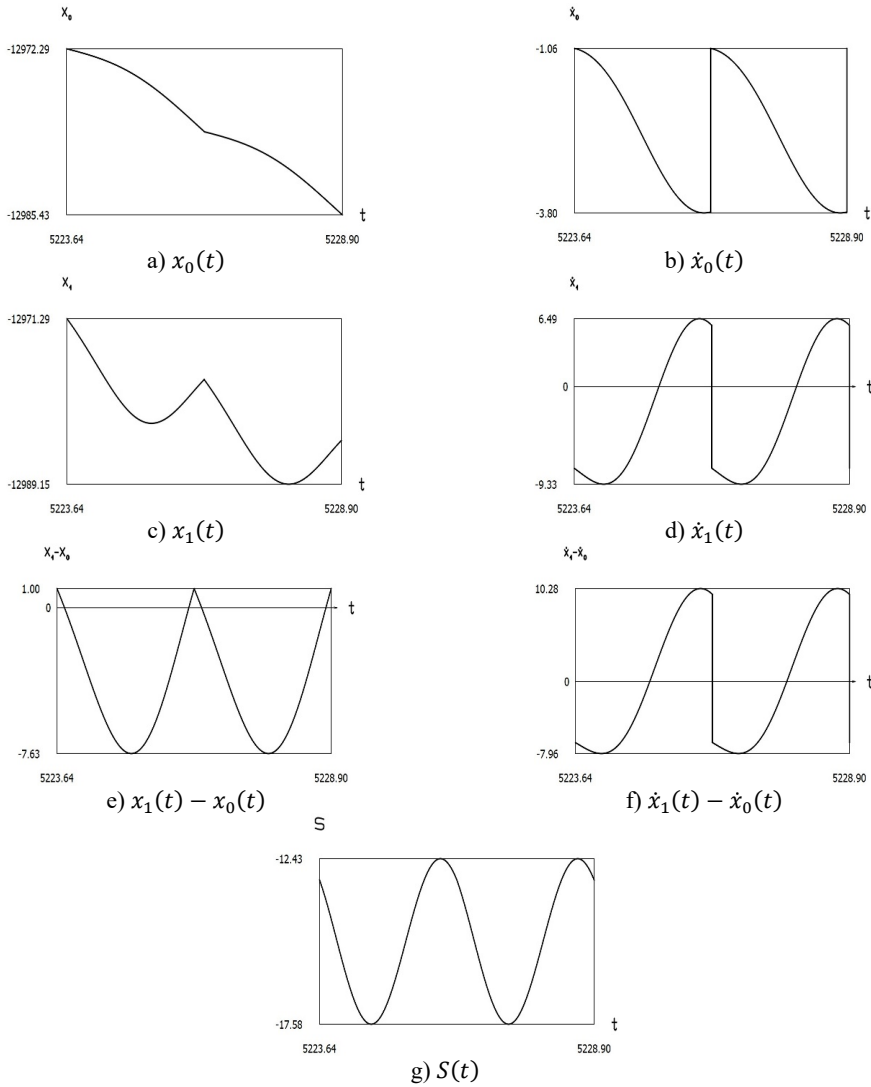


Fig. 9. Dynamics of the manipulator when $\delta\omega = 0.2$

From the presented results in the vicinity of the first resonance, the effectiveness of operation of the manipulator with vibration impact drive at the first resonant frequency is observed.

6. Investigation of the harmonics of the conservative system

The following parameters of the system are assumed:

$$l = 1, \quad \mu = 5, \quad p = 1. \quad (46)$$

The following initial conditions are assumed:

$$x_0(0) = 0, \quad x_1(0) = l, \quad \dot{x}_0(0) = 0, \quad \dot{x}_1(0) = S. \quad (47)$$

Constant part and the first three harmonics of S are presented in Fig. 11.

From the graphical representation the first three eigenfrequencies of the investigated

manipulator with vibration impact drive are seen.

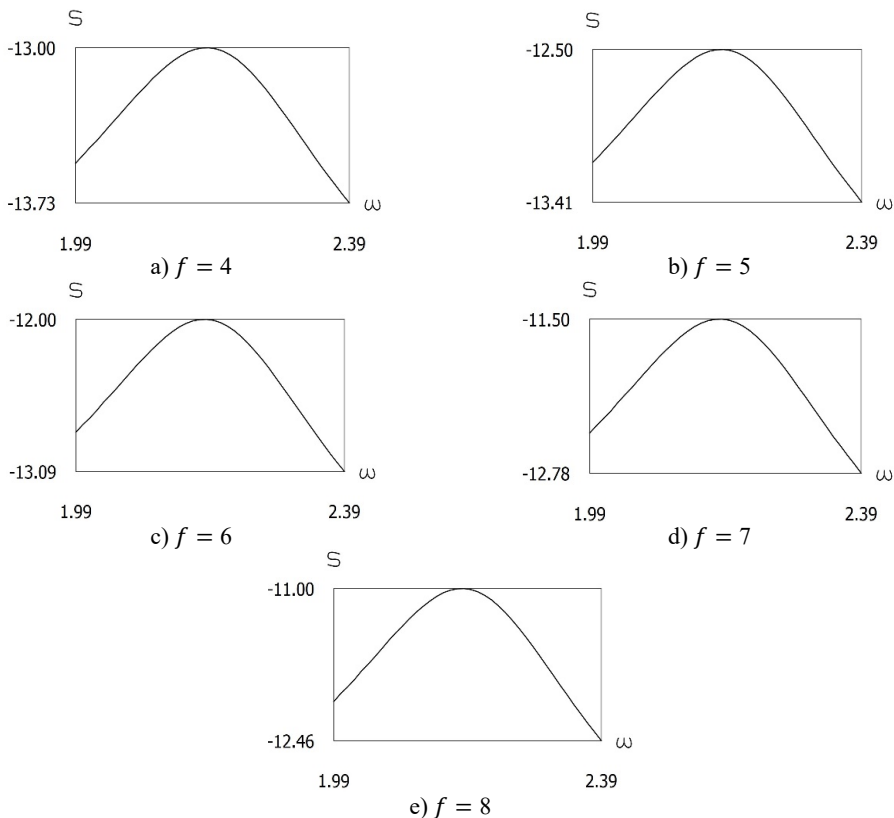


Fig. 10. Quantity of motion during impact as function of frequency of excitation

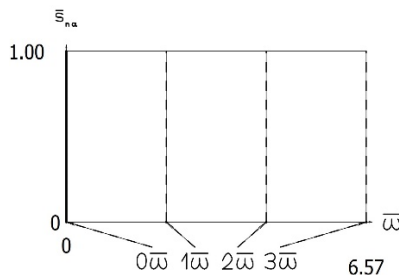


Fig. 11. Constant part and the first three harmonics of S

7. Conclusions

The proposed 1D manipulator with vibration impact drive is simple because it does not include the self-stopping mechanism and its control is performed by the external force.

In the conservative case of the system static position of equilibrium of the impact pair can be with negative, zero or positive tightening. Zero tightening is advantageous for practical applications and thus the investigations for zero tightening are performed.

In the case of zero tightening eigenfrequency and period of the system does not depend on the quantity of motion of impact excitation. In this case the static position of the system in the position of equilibrium is the point of contact of both masses.

In the case of harmonic forced excitation resonant motions take place in the vicinity of the

eigenfrequency of the conservative system with zero tightening.

This enables us to create advanced manipulators and complicated robots from manipulators of this type.

Acknowledgements

The authors have not disclosed any funding.

Data availability

The datasets generated during and/or analyzed during the current study are available from the corresponding author on reasonable request.

Author contributions

Kazimieras Ragulskis: conceptualization, formal analysis, investigation, methodology, supervision, validation, visualization, writing – original draft preparation, writing – review and editing. Liutauras Ragulskis: visualization.

Conflict of interest

The authors declare that they have no conflict of interest.

References

- [1] W. V. Wedig, “New resonances and velocity jumps in nonlinear road-vehicle dynamics,” *Procedia IUTAM*, Vol. 19, pp. 209–218, Jan. 2016, <https://doi.org/10.1016/j.piutam.2016.03.027>
- [2] T. Li, E. Gourc, S. Seguy, and A. Berlioz, “Dynamics of two vibro-impact nonlinear energy sinks in parallel under periodic and transient excitations,” *International Journal of Non-Linear Mechanics*, Vol. 90, pp. 100–110, Apr. 2017, <https://doi.org/10.1016/j.ijnonlinmec.2017.01.010>
- [3] V. A. Zaitsev, “Global asymptotic stabilization of periodic nonlinear systems with stable free dynamics,” *Systems and Control Letters*, Vol. 91, pp. 7–13, May 2016, <https://doi.org/10.1016/j.sysconle.2016.01.004>
- [4] H. Dankowicz and E. Fotsch, “On the analysis of chatter in mechanical systems with impacts,” *Procedia IUTAM*, Vol. 20, pp. 18–25, Jan. 2017, <https://doi.org/10.1016/j.piutam.2017.03.004>
- [5] S. Spedicato and G. Notarstefano, “An optimal control approach to the design of periodic orbits for mechanical systems with impacts,” *Nonlinear Analysis: Hybrid Systems*, Vol. 23, pp. 111–121, Feb. 2017, <https://doi.org/10.1016/j.nahs.2016.08.009>
- [6] W. Li, N. E. Wierschem, X. Li, and T. Yang, “On the energy transfer mechanism of the single-sided vibro-impact nonlinear energy sink,” *Journal of Sound and Vibration*, Vol. 437, pp. 166–179, Dec. 2018, <https://doi.org/10.1016/j.jsv.2018.08.057>
- [7] J. S. Marshall, “Modeling and sensitivity analysis of particle impact with a wall with integrated damping mechanisms,” *Powder Technology*, Vol. 339, pp. 17–24, Nov. 2018, <https://doi.org/10.1016/j.powtec.2018.07.097>
- [8] E. Salahshoor, S. Ebrahimi, and Y. Zhang, “Frequency analysis of a typical planar flexible multibody system with joint clearances,” *Mechanism and Machine Theory*, Vol. 126, pp. 429–456, Aug. 2018, <https://doi.org/10.1016/j.mechmachtheory.2018.04.027>
- [9] U. Starossek, “Forced response of low-frequency pendulum mechanism,” *Mechanism and Machine Theory*, Vol. 99, pp. 207–216, May 2016, <https://doi.org/10.1016/j.mechmachtheory.2016.01.004>
- [10] S. Wang, L. Hua, C. Yang, Y.O. Zhang, and X. Tan, “Nonlinear vibrations of a piecewise-linear quarter-car truck model by incremental harmonic balance method,” *Nonlinear Dynamics*, Vol. 92, No. 4, pp. 1719–1732, Feb. 2018, <https://doi.org/10.1007/s11071-018-4157-6>
- [11] P. Alevras, S. Theodossiades, and H. Rahnejat, “On the dynamics of a nonlinear energy harvester with multiple resonant zones,” *Nonlinear Dynamics*, Vol. 92, No. 3, pp. 1271–1286, Feb. 2018, <https://doi.org/10.1007/s11071-018-4124-2>

- [12] A. Sinha, S. K. Bharti, A. K. Samantaray, G. Chakraborty, and R. Bhattacharyya, “Sommerfeld effect in an oscillator with a reciprocating mass,” *Nonlinear Dynamics*, Vol. 93, No. 3, pp. 1719–1739, Apr. 2018, <https://doi.org/10.1007/s11071-018-4287-x>
- [13] G. Habib, G. I. Cirillo, and G. Kerschen, “Isolated resonances and nonlinear damping,” *Nonlinear Dynamics*, Vol. 93, No. 3, pp. 979–994, Apr. 2018, <https://doi.org/10.1007/s11071-018-4240-z>
- [14] V. Glazunov, *New mechanisms in contemporary robot engineering*. Moscow: Tehnosphere, 2018.
- [15] H. Shinno, H. Yoshioka, and H. Sawano, “A newly developed long range positioning table system with a sub-nanometer resolution,” *CIRP Annals*, Vol. 60, No. 1, pp. 403–406, Jan. 2011, <https://doi.org/10.1016/j.cirp.2011.03.027>
- [16] W. Kokuyama, T. Shimoda, and H. Nozato, “Primary accelerometer calibration with two-axis automatic positioning stage,” *Measurement*, Vol. 204, p. 112044, Nov. 2022, <https://doi.org/10.1016/j.measurement.2022.112044>
- [17] K. Cai, Y. Tian, F. Wang, D. Zhang, and B. Shirinzadeh, “Development of a piezo-driven 3-DOF stage with T-shape flexible hinge mechanism,” *Robotics and Computer-Integrated Manufacturing*, Vol. 37, pp. 125–138, Feb. 2016, <https://doi.org/10.1016/j.rcim.2015.07.004>
- [18] C. Lin, M. Jiang, and S. Zheng, “Establishment and verification of the analytical model for the critical parameters in the kinematics model of the precision positioning stage,” *Sensors and Actuators A: Physical*, Vol. 320, p. 112572, Apr. 2021, <https://doi.org/10.1016/j.sna.2021.112572>
- [19] V. Ragulskienė, *Vibro-Shock Systems (Theory and applications)*. Vilnius: Mintis, 1974.
- [20] R. Kurila and V. Ragulskienė, *Two-Dimensional Vibro-Transmissions*. Vilnius: Mokslas, 1986.
- [21] Ragulskis et al., *Vibromotors for Precision Microrobots*. New York: Hemisphere, 1987.



Kazimieras Ragulskis is a member of Academies of Sciences of the USSR (later of the Russian Academy of Sciences) and Lithuania. His research area is precise vibromechanics and vibroengineering, the basis of it are vibrations and waves in nonlinear dynamical systems.



Liutauras Ragulskis is working as a research associate at Vytautas Magnus University. His research interests are numerical calculations of vibrating systems.