# Liquid induced vibrations of partially liquid-filled elastic cylindrical-conical shells 

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#### Abstract

The main objective of this study to develop efficient numerical techniques based on coupled finite and boundary element methods to estimate natural vibration frequencies in compound reservoirs with liquid. Elastic structures composed of cylindrical and conical shells connected by rings have been examined. The space between the shells is filled with an ideal, incompressible fluid. The advantage of the proposed approach lies in the ability to investigate both free and forced vibrations of empty and liquid-filled shells structures within the framework of the single computer technology. The benchmark test calculations demonstrated the high accuracy and efficiency of the proposed approach. The novelty and practical value of the obtained results lie in the ability to investigate fuel tanks of complex shapes under various conditions.


Keywords: coaxial elastic shells, boundary and finite element methods.

## Nomenclature

| n | External unit normal to wetted surfaces |
| :--- | :--- |
| $S_{0}$ | Liquid free surface |
| $S_{1}$ | Walls and bottom wetted surfaces |
| $R_{1}$ | Radius of circular cylindrical shell |
| $L$ | The cylindrical shell length |
| $R_{2}$ | Larger radius of the conical shell |
| $R_{3}$ | Smaller radius of the conical shell |
| $\rho_{l}$ | The liquid density |
| $\rho_{s}$ | The structure material density |
| $E$ | Young's module |
| $v$ | Poisson's ratio |

## 1. Introduction

Shells and shell-like structures play a significant role in modern mechanical engineering. These components are commonly used as fuel containers in launch vehicles, oil storage tanks, containers for storing various substances, including potentially dangerous ones. Nowadays, the cylindrical [1], conical, toroidal and coaxial cylindrical shells [2] are commonly employed as fuel tanks. The most precise strength assessment of liquid-filled tanks under significant external forces often relies on comprehensive, full-scale field experiments. However, these experiments are both expensive and pose risks. So, virtual experimentation is gaining prominence in scientific research as a safer and more cost-effective alternative. Various effective methods have been elaborated for computer modeling dynamic processes, including finite (FEM) [3] and boundary (BEM) [4] element methods, finite volume methods [5], and series expansions [6]. The initial step in dynamic
modeling involves determining the frequencies and modes of natural vibrations. This aspect has been explored in [2], [7], and [8]. Papers [9], [10] are dedicated to developing analytical methods for analyzing vibrations in fluid-filled shells with baffles. The investigation [2] focused on examining the natural frequencies of toroidal and coaxial cylindrical shells, while [8] and [11] delved into the mitigation of liquid sloshing using rigid baffles. Furthermore, research involving the theoretical and experimental stability aspects of cylindrical shells and beams has been accomplished in [1], [12]. The design and manufacture of each new powerful launch vehicles necessitate the development of new fuel tanks. Compound shells are becoming more common in these tanks. To mitigate undesirable resonance, it is necessary to study the entire frequency spectrum of the tank's vibrations. However, the vibrations of compound shells have not received sufficient attention in the literature [13]. Therefore, the vibrations study of compound coaxial cylindrical and cylindrical-conical shells with liquids is a topical problem.

## 2. Methodology

### 2.1. Problem statement

Elastic compound shells of revolution, partially filled with liquids, are considered, Fig. 1(a-b). The wetted surface of the shell is denoted as $S_{1}$, while $S_{0}$ represents the liquid free surface. It is assumed that the liquid is an ideal and incompressible. Reservoirs consist of coaxial cylindrical and conical shells connected by rings. The liquid is located between the shells. The surfaces of the shells and bottoms are wetted, and the free surface at rest is at a height of $H$, forming a ring, which is described in polar coordinate system as $\left\{z=H, R_{1}<r<R_{0}\right\}$. Here for the shell structure, depicted in Fig. 1(a), $R_{0}=R_{3}+\left(R_{2}-R_{3}\right)(L-H) / L$, whereas $R_{0}=R_{3}+\left(R_{2}-R_{3}\right) H / L$ for the shell structure, depicted in Fig. 1(b). The wetted surface $S_{1}$ consists of cylindrical and conical parts as well as the bottom. The liquid within the shells undergoes vibrations due to applied loads. The resulting flow from these loads is assumed to be irrotational. If $\mathbf{V}\left(V_{x}, V_{y}, V_{z}\right)$ represents the fluid velocity vector, the incompressibility condition of the continuum can be derived from the equation $\operatorname{div} \mathbf{V}=0$. So, the velocity potential $\Phi$ can be introduced as $\mathbf{V}=\operatorname{grad} \Phi$, and it leads to the Laplace equation. The motion governing equations of liquid-filled elastic shells are given by [13]:

$$
\begin{equation*}
\mathbf{L U}+\mathbf{M} \ddot{\mathbf{U}}=\mathbf{P}_{\mathbf{1}}+\mathbf{Q}, \tag{1}
\end{equation*}
$$

where $\mathbf{L}$ and $\mathbf{M}$ denote the operators corresponding to elastic and mass forces acting on the shell structure, $\mathbf{U}=\left(U_{1}, U_{2}, U_{3}\right)$ is the displacement vector, $\mathbf{Q}$ is the vector of external surface loads, and $\mathbf{P}_{\mathbf{1}}$ signifies hydrodynamic pressure, $\mathbf{P}_{\mathbf{1}}=p \mathbf{n}$.


Fig. 1. Cylindrical - conical compound shells and their drafts
The pressure value $p$, according to the Bernoulli integral, can be represented as $p=-\rho_{l} \frac{\partial \Phi}{\partial t} p_{0}$, where $p_{0}$ is atmospheric pressure. The region occupied by the liquid is denoted by $\Omega$, and $\mathbf{P}$ is for
points within $\Omega$. Boundary conditions for the Laplace equation within $\Omega$ are outlined as follows. On surfaces $S_{1}$, no-slip conditions are enforced, and on $S_{0}$ the dynamic condition is set:
$\nabla^{2} \Phi=0, \quad \mathbf{P} \in \Omega, \quad \frac{\partial \Phi}{\partial \mathbf{n}}=(\dot{\mathbf{U}}, \mathbf{n}), \quad \mathbf{P} \in S_{1}, \quad \frac{\partial \Phi}{\partial \mathrm{t}}=0, \quad \mathbf{P} \in S_{0}$.
So, it is necessary to determine the unknown functions $\Phi$ and $\mathbf{U}$ through Eq. (1-2) including boundary and initial data. The unknown vector-function is represented as following series:
$\mathbf{U}(x, y, z, t)=\sum_{k=1}^{n} \mathbf{u}_{k}(x, y, z) c_{k}(t)$,
where functions $\mathbf{u}_{k}(x, y, z)$ are natural vibration modes of the empty shell structure, and $c_{k}(t)$ are time-dependent unknown factors. To determine the vibration modes $\mathbf{u}_{k}(x, y, z)$, the FEM is used [4]. To determine potential $\Phi$, boundary value problem Eq. (2) is solved. The following series is incorporated $\Phi(r, \theta, z, t)=\sum_{k=1}^{n} \varphi_{k}(r, \theta, z) \dot{c}_{k}(t)$. Here the functions $\varphi_{k}(x, y, z)$ are solutions of boundary value problem Eq. (2) with $\frac{\partial \varphi_{k}}{\partial \mathbf{n}}=\left(\mathbf{u}_{k}, \mathbf{n}\right)$.

For shells of revolution, all unknown functions for each $l$ can be expressed in the form [4]:

$$
\begin{equation*}
\mathbf{u}_{k}(r, \theta, z)=\mathbf{u}_{k}(r, z) \cos (l \theta), \quad \varphi_{k}(r, \theta, z)=\varphi_{k}(r, z) \cos (l \theta) . \tag{4}
\end{equation*}
$$

Such presentations simplify the dynamical analysis by reducing the problem to a series of uncoupled single-degree-of-freedom systems, depending on $l$ only.

## 3. Results and discussion

### 3.1. Validation

The validation study involves truncated isotropic conical tanks. Let $R_{1}$ and $R_{2}$ are the cone radii at its larger and smaller ends, $\alpha$ is the semi-vertex angle, and $H$ is the cone height. In numerical simulations, the shell thickness is $h / R_{1}=0.01$. The semi-vertex angle is $\alpha=45^{\circ}$, $H / R_{2}=1, E=2.11 \cdot 10^{6} \mathrm{MPa}, v=0.3, \rho_{s}=8000 \mathrm{~kg} / \mathrm{m}^{3}, \rho_{l}=1000 \mathrm{~kg} / \mathrm{m}^{3}$, and $R_{1}=1 \mathrm{~m},[16]$. Shell edge fixations are clamped-clamped (CC), clamped-simply supported (C-SS), simply supported -clamped (SS-C), and clamped -free (C-F) ones. Convergence was achieved when the number of finite elements on the tank walls equals 30 . Table 1 provides comparison of obtained results with data of [16]. The specified above boundary conditions are considered, with wave numbers $l=0,1, \ldots, 7$.

Table 1. Dimensionless frequency parameter $\Omega_{1}$ for empty conical tank

| $l$ | Boundary conditions |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | SS-C |  | C-SS |  | C-C |  |
|  | present FEM | Shu [16] | present FEM | Shu [16] | present FEM | Shu [16] |
| 0 | 0,8701 | 0,8700 | 0,7157 | 0,7151 | 0.8730 | 0.8732 |
| 1 | 0,8119 | 0,8118 | 0,7097 | 0,7090 | 0,8120 | 0,8120 |
| 2 | 0,6614 | 0,6613 | 0,6479 | 0,6475 | 0,6694 | 0,6696 |
| 3 | 0.5247 | 0.5245 | 0,5203 | 0,5201 | 0,5427 | 0,5428 |
| 4 | 0,4321 | 0,4319 | 0,4166 | 0,4161 | 0,4563 | 0,4566 |
| 5 | 0,3828 | 0,3826 | 0.3596 | 0.3592 | 0,4087 | 0,4089 |
| 6 | 0,3739 | 0,3737 | 0,3458 | 0,3450 | 0,3960 | 0,3964 |
| 7 | 0,3984 | 0,3981 | 0,3651 | 0,3648 | 0,4141 | 0,4143 |

In this comparison, dimensionless frequencies are obtained for truncated isotropic elastic
conical tanks for $k=1, \alpha=45^{\circ}$, using the frequency parameter $\Omega_{1}$ as in [16]. The results obtained by the proposed reduced BEM are in good agreement with data of [16].

### 3.2. Vibrations of elastic cylindrical-conical shells.

The vibration frequencies and modes of the empty and fluid-filled structures, depicted in Fig. 1, are estimated. The next parameters are chosen: $L=1 \mathrm{~m}, R_{1}=0.5 \mathrm{~m}, R_{2}=0.5 \mathrm{~m}$, $R_{3}=0.25 \mathrm{~m}$, the filling level $H=1.25 \mathrm{~m}$. The material and mechanical properties are $\rho_{s}=7800 \mathrm{~kg} / \mathrm{m}^{3}, E=2.110^{6} \mathrm{MPa}, v=0.3$, the thickness of the shell structure is 3 mm . The structure has been rigidly bottom fixed. The frequencies and modes of the empty structure are calculated first; the FEM [6] is used. Table 2 shows the lowest vibration frequencies of unfilled structures, Fig. 1(a)-(b).

Table 2. Vibration frequencies of unfilled tanks, Hz

| No. | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Shell structure, Fig. 1(a) |  |  |  |  |  |  |  |  |
| Values | 127,59 | 128,07 | 132,6 | 133,6 | 157,79 | 158,55 | 167,61 | 169,86 |
| Shell structure, Fig. 1(b) |  |  |  |  |  |  |  |  |
| Values | 121,57 | 122,04 | 124,8 | 125,56 | 154,81 | 155,42 | 157,98 | 158,18 |

The frequencies of the structure depicted in Fig. 1(b) slightly differ from the frequencies of the structure shown in Fig. 1(a). The corresponding vibration modes $\mathbf{u}_{k l}$ are shown in Fig. 2.

b) Modes of shell structure Fig. 1(b)

Fig. 2. Vibration modes of the unfilled structures
Note that, as in [16], the lowest frequencies correspond to 6th and 7th wavenumbers. It should also be noted the presence of practically multiple frequencies. Next, the base functions $\varphi_{k}(r, \theta, H)$ and the corresponding frequencies are obtained. These functions are vibration modes of the structure walls, considering the liquid added masses. The frequencies are provided in Table 3.

Table 3. Vibration frequencies of elastic walls considering liquid added masses ( Hz )

| No | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Shell structure, Fig. 1(a) |  |  |  |  |  |  |  |  |
| Values | 45,327 | 45,597 | 47,141 | 47,602 | 51,264 | 51,433 | 52,787 | 54,053 |
| Shell structure, Fig. 1(b) |  |  |  |  |  |  |  |  |
| Values | 43,542 | 43,619 | 48,367 | 48,384 | 50,004 | 50,291 | 57,81 | 57,896 |

Significant decreasing in frequencies compared to the empty shells has been observed.

Multiple frequencies are also available. Fig. 3 depicts some characteristic vibration modes.
The lowest frequency corresponds to the eighth wave number. This result cannot be predicted theoretically, and only computer experiments allow to understand the nature of vibration patterns.

b) Modes of shell structure Fig. 1(b)

Fig. 3. Vibration modes $\varphi_{k}(r, \theta, H)$ of elastic walls considering liquid added masses

## 4. Conclusions

The approach has been developed based on the application of reduced finite and boundary element methods as well as mode superposition methods for analysing free vibrations of the coaxial elastic shell structures partially filled with liquid. The developed approach allows for the investigation of free and forced vibrations of unfilled shells, as well as for the problem of vibrations of liquid-filled shells. The further intention is to analyse the forced vibrations.

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## Data availability

The datasets generated during and/or analyzed during the current study are available from the corresponding author on reasonable request.

## Conflict of interest

The authors declare that they have no conflict of interest.

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